

## RESEARCH ARTICLE

# Modelling the effects of solar activity on the ionospheric $F_2$ critical frequency over Wakkanai

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**Abstract:** For many decades ionospheric researchers investigated the variations in the ionosphere due to solar activity. The suggested relevant models are based on single-station data, considering regional and global geographic conditions. The present study investigated the impact of the solar cycles 21<sup>st</sup> (1976 to 1986) and 23<sup>rd</sup> (1996 to 2008) on the ionospheric  $F_2$  layer's critical frequency ( $f_0F_2$ ) at mid-latitude over the Wakkanai region (45.39° N, 141.68° E), Japan. The statistical analyses showed that monthly median  $f_0F_2$  has a significant non-linear association with high sunspot numbers (SSN) over Wakkanai, which represents a saturation effect depending on the time of the day in different months and on the magnitude of the solar cycle. Polynomial empirical models of  $f_0F_2$  based on parameters such as SSN and geomagnetic index  $A_p$  were examined. Considering the rate of change in solar activity factor much improved the accuracy of our empirical model and also reduces the hysteresis effect. The most appropriate empirical model for single-station diurnal models of  $f_0F_2$  was developed using Fourier series. Diurnal models incorporate Japanese standard time, months and sunspot numbers. The computed  $f_0F_2$  models were compared with the IRI-2012 model's predicted  $f_0F_2$  values, which demonstrated the better accuracy of the Fourier model compared to the global IRI model. The models obtained in this study are useful for researchers and organisations working in the field of sunspot performance relating to the dynamics of the ionosphere.

**Keywords:**  $f_0F_2$ , geomagnetic  $A_p$  index, ionosphere, mid-latitude, sunspot.

## INTRODUCTION

The atmosphere of the earth consists of the layers of gases surrounding the earth. Each layer differs in the composition of gas species, altitudes, temperatures and pressures. The Earth's atmosphere extends from 50 km to 1000 km where the solar radiations (especially; EUV and X-rays) highly influence it and create the region of charged particles (positive and electrons). The region where the concentration of charged particles is high enough to reflect the trajectory of radio waves sent from the ground is called the ionospheric layer. The ionosphere lies between 140 km to 600 km in altitude and under certain solar-terrestrial conditions; ionosphere is sub-divided into D, E,  $F_1$ ,  $F_2$  layers.  $F_2$  layer is the most important region of the ionosphere from the point of view of radio communication and navigation. However, the ionosphere is not a stable ionised medium. The variation of electron density in the  $F_2$  layer possibly enhances or reduces the transmission of a signal propagating at a high frequency between the transmitter and the receiver or it may suddenly interrupt the communication. This disturbance highly depends on the level of solar activity (Sizun, 2003; Jixiang, 2006).

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### Variation in ionosphere under solar activity

Much effort has been made to study the variations in Earth's ionosphere with respect to different locations on Earth and at different times. Thus, by using ionospheric characteristic data, a large number of stations have developed regional and global models of the ionosphere. These empirical models are established by conducting statistical studies of ionospheric data over a prolonged period. Bilitza (2002) has presented a review of the many available empirical models. Among those models, the international reference ionosphere (IRI) model is broadly used. From different ionospheric parameters, the critical frequency of the  $F_2$  layer ( $f_0F_2$ ) is one of the most vital parameters that represents the concentration of electron density in the  $F_2$  layer. Critical frequency is defined as the maximum frequency, which is transmitted vertically to the sky and is refracted back to the ground. The  $f_0F_2$  describes the  $F_2$  region of the ionosphere and at night time,  $f_0F_2$  completely explores the F region.

At middle and low latitudes, the solar extreme ultraviolet radiation (EUV) is the primary source of ionisation in the F region (Lakshmi *et al.*, 1988). Thus, ionospheric empirical models depend on the high energy radiation coming from the Sun. However, the intensity of EUV and other high energy radiation depends on the magnitude of the solar cycle. As a result of the lack of long-term direct measurements of solar EUV and other high energy radiation, researchers have to rely on other solar activity proxies. X-rays, EUV and UV can be precisely estimated by different solar indices such as sunspot numbers, solar radio flux at 10.7 cm wavelength (F10.7), Mg II core-to-wing index and He 1083 (Maltby & Albregtsen, 1979; Sams III *et al.*, 1992; Acton, 1996; Floyd *et al.*, 2005; Ramesh & Rohini, 2008; Lukianova & Mursula, 2011; Yan *et al.*, 2011; Usoskin, 2017). The most interesting solar photospheric feature that can be simply observed through a solar telescope in white light are the sunspots, which are the seat of solar activity.

Sunspots initiate as smaller, darker regions on the solar disk which may develop into larger and much darker spots, surrounded by light (dark) regions called penumbrae (Antia *et al.*, 2003). The sunspot number is one of the most commonly used solar proxies because the sunspot number is well correlated to high energy radiation and with other solar proxies. It is convenient to use it because of its long series of observations and its characteristics of reliability and predictability. The solar activity reliance of ionospheric characteristics has been found in early regular ionospheric observations (Richards, 2001; Sethi *et al.*, 2002; Atac *et al.*, 2009;

Maruyama, 2010; Adeniyi & Ikubanni, 2013; Mielich & Bremer, 2013; Elias *et al.*, 2014; Ikubanni & Adeniyi, 2017). Earlier researchers suggested the idea about the linear relationship between sunspot numbers and  $f_0F_2$  for every hour and month, which has been frequently used for global, regional as well as single station ionospheric models (De Franceschi & De Santis, 1994; Bilitza, 2001; Holt *et al.*, 2002), while after further studying proved that nonlinear relation exists (Pancheva & Mukhtarov, 1998; Ozguc *et al.*, 2007; Dabas *et al.*, 2008).

### IRI model and its uses

International reference ionosphere is a purely empirical standard model of the ionosphere, based on the large collection of ground-based and space observations. It provides monthly average values of ionospheric parameters (such as electron density, ion and electron temperatures, composition of gas species etc.) under the quiet geomagnetic condition (Rawer *et al.*, 1978). In the studies, the  $f_0F_2$  data of IRI-2012 ([https://omniweb.gsfc.nasa.gov/vitmo/iri2012\\_vitmo.html](https://omniweb.gsfc.nasa.gov/vitmo/iri2012_vitmo.html)) model was considered for comparison with the spectral model. The critical frequency  $F_2$  layer directly represented the electron density of the  $F_2$ . The highest electron density ( $N_mF_2$ ) reaches the  $F_2$  peak height ( $h_mF_2$ ) from where it divides into lower and upper parts. In IRI both parts are normalised to the  $F_2$  peak density and height. Electron density below the  $F_2$  peak is normalised to the E and  $F_2$  peaks. The region between  $F_1$  and  $F_2$  height is of special interest because of its effect on HF radio wave propagation. Electron density ( $N_e$ ) in that region is described by the IRI function as shown in equation (1),

$$N_e(h) = N_mF_2 \times \frac{\exp(-Z^{B_1})}{\cosh(Z)}, \quad Z = \frac{h_mF_2 - h}{B_0} \quad \dots(1)$$

where the shape of the electron density profile is determined by the lower thickness parameter ( $B_0$ ) and the shape parameter ( $B_1$ ). The  $B_0$  is defined as ( $B_0 = h_mF_2 - hx$ ) in which the  $hx$  is the height where the density profile dropped down to ( $0.24 \times N_mF_2$ ). The parameters  $B_0$  and  $B_1$  are upgraded every few years by IRI for better description of the global and temporal variation of the ionosphere (Bilitza *et al.*, 2014).

In this study, we have illustrated the effects of sunspots on  $f_0F_2$  over Wakkanaï. Similar to other models, we have constructed a local-station model of  $f_0F_2$  that usually gives an average value of  $f_0F_2$  incorporating the effect of temporal, seasonal effect, solar cycle and the latitude of the station. It has been observed that specific

models of a particular region are very useful and widely referred to ionospheric services (Pancheva & Mukhtarov, 1998). The diurnal  $f_0F_2$  models for each month for solar maximum and solar minimum years are successfully presented in this study.

## METHODOLOGY

From 1948 onwards the ionosonde/digisonde station in Wakkanai has consistently measured ionospheric parameters. These measurements are recorded by the World Data Center (WDC) for Ionosphere and Space Weather in the Japanese city of Tokyo. The center is under the management of the National Institute of Information and Communications Technology (NICT), Japan. The geographic location of Wakkanai is  $45.39^\circ\text{N}$  and  $141.68^\circ\text{E}$ , in the mid-latitude region of Japan. In this research, the hourly data of  $f_0F_2$  in Japanese standard time (JST) and sunspot number were used for the years of 21<sup>st</sup> and 23<sup>rd</sup> solar cycles, instead the years of 22<sup>nd</sup> (1986 – 1996) solar cycle due to the huge deficiency in observational data of  $f_0F_2$ . To study the average behaviour of the ionosphere the monthly median hourly values of  $f_0F_2$  were extracted from the entire database, while the missing  $f_0F_2$  values were filled with IRI-2007 model values obtained from the Virtual Ionosphere Thermosphere Mesosphere Observatory (VITMO) (<http://omniweb.gsfc.nasa.gov/vitmo/>). Considering only the effect of seasons,  $f_0F_2$  is independent of solar zenith angle. Generally,  $f_0F_2$  peaks in winter and decreases in summer, and this effect is known as winter anomaly. In the summer season, the enhancement in night-time  $f_0F_2$  called midlatitude summer nighttime anomaly (MSNA) has also been observed over Wakkanai  $45^\circ\text{N}$  (Mushtaq *et al.*, 2015). The preference of monthly median values over monthly mean values is because of the 27 days solar rotation. The solar rotation has conspicuous variation in EUV radiation, which contributes 85 – 95 % in electron density in the ionosphere (Ruiping *et al.*, 2012).

Regression method was used to investigate the dependence of sunspot numbers on the ionospheric  $F_2$  layer (i.e.,  $f_0F_2$ ) over Wakkanai. Regression analysis makes it possible to find the line which best describes the association between two variables. This line is usually referred to as the ‘line of best fit’ (Jaisingh, 2006). Spectral models explain the complex data as a hypothesis; it predicts the dynamics of  $f_0F_2$  in the frequency domain by observing long-term data records. The diurnal cycles have a significant impact on many phenomena in the Earth’s atmosphere, mainly on the ionosphere. To model the complex nature of the ionosphere the spectral approach has been implemented.

For solar proxy, the monthly mean sunspot numbers are used for the said solar cycles given by the Royal Observatory of Belgium located in Uccle, Belgium, which has the geographic position  $50.798^\circ\text{N}$  and  $4.358^\circ\text{E}$ . The relative sunspot number is calculated on the basis of all observations available from different observatories. However, for representing global geomagnetic activity, the monthly mean geomagnetic index  $A_p$  is utilised for 1976 – 1986 and 1996 – 2008. Prior to January 1997, data were taken from the Institute für Geophysik, Gottingen, Germany. Since 1997 up to now, documentation, calculation, and distribution of indices have been moved to Geo Forschungs Zentrum, Helmholtz Center, Postdam, Germany. In this study, data extraction, calculations and plotting were all done in Python 2.7, MATLAB 2015, Easyfit.

## RESULTS AND DISCUSSION

### Empirical models of $f_0F_2$ and sunspot number

It is significant to find the empirical models between monthly median hourly  $f_0F_2$  with any best fitted solar proxies. The key parameter used is the sunspot number for ionospheric long-term predictions. Sunspots are one of the best ways to indicate the emission of high energy radiation from the solar surface and its atmosphere (i.e., EUV and X-rays), which are the main sources of ionisation in the ionosphere. We fitted four regression models and to authenticate the correlation of solar activity on  $f_0F_2$ , F-test has been performed at 99 % confidence interval which verified that regression analyses between SSN and  $f_0F_2$  are statistically significant for all months. The first regression model is the linear line of best fit. The linear relationships between  $f_0F_2$  and solar proxies have been used in many previous models which are in good agreement with our results (Gulyaeva, 1999; Holt *et al.*, 2002; Yanhong *et al.*, 2002; Yadav *et al.*, 2011). A good linear relation depends on the following situations:

i) It has been found that at moderate sunspot numbers ( $60 < \text{SSN} < 160$ ) the linear relation better explained the behaviour of  $f_0F_2$ . However, no linear correlation is found at low and high SSN (Ikubanni *et al.*, 2013).

ii) In different locations on the Earth, the nature of ionosphere is varying. At equator and low-latitudes the effect of saturation is so strong that no linear relation is suitable for explaining  $f_0F_2$ . Mid-latitudes show less variability in  $f_0F_2$  at high SSN, therefore the linear relation can also be considered (Liu *et al.*, 2006). Equation (2) describes the linear correlation between the sunspot number and  $f_0F_2$  as,

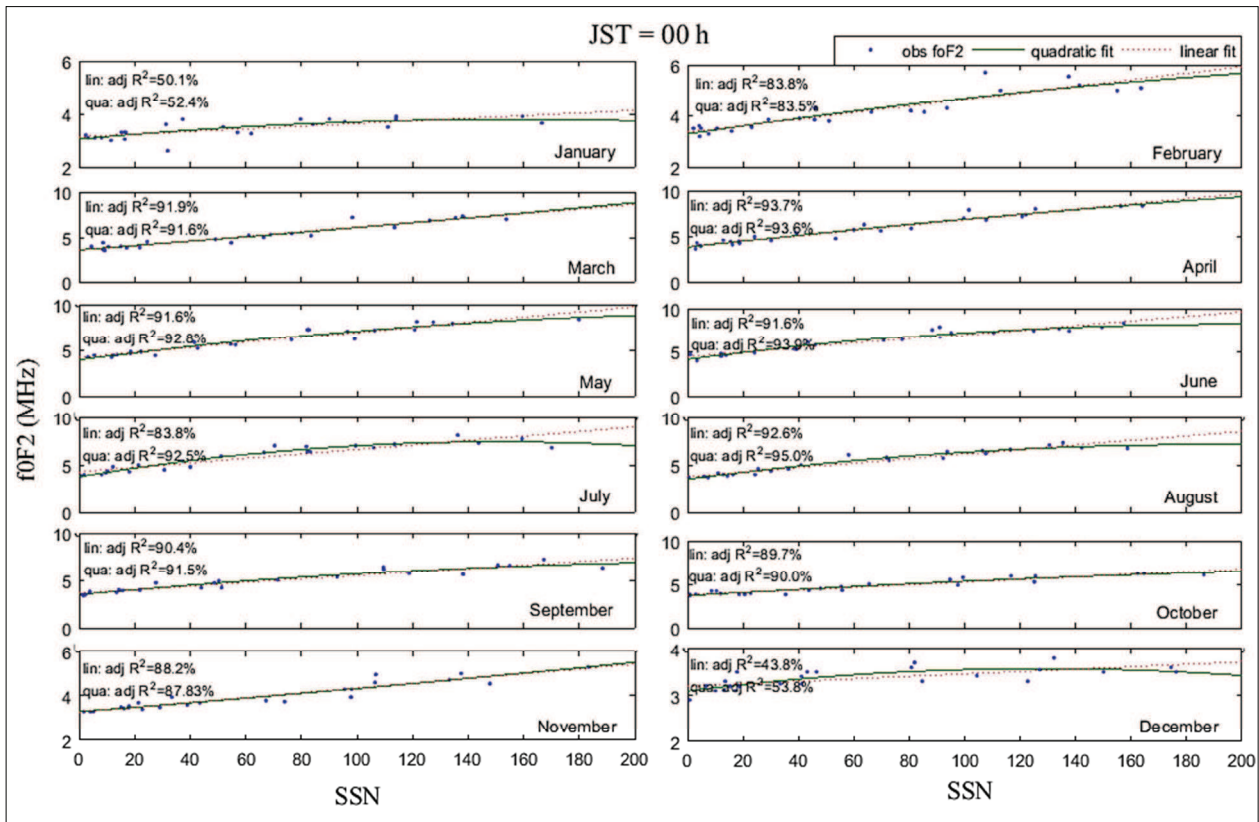
$$foF2(t)_m = a_{0m}(t) + a_{1m}(t) \cdot SSN_m \quad \dots(2)$$

where  $\alpha_0$  and  $\alpha_1$  are the coefficients, while  $t$  and  $m$  show the Japanese standard time in hours and different months, respectively and SSN is the monthly mean sunspot number for the current month. We analysed the relationship at midnight ( $t = 00$  h) and noontime ( $t = 12$  h) in JST, whilst  $m$  shows the whole months of a year, and the linear line of best fit is represented as dotted lines as shown in Figures 1 and 2.

At midnight, the linear regression fit seems good in many months, especially around the equinoctial time and the highest adjusted coefficient of determination ( $adj\_R^2$ ) occurred in April ( $adj\_R^2 = 93.7\%$ ), while the lowest  $adj\_R^2$  was found in December (i.e.,  $43.8\%$ ). Higher values of  $adj\_R^2$  illustrate that the regression model explains the majority of the variations in  $f_0F_2$  while low values of  $adj\_R^2$  show the opposite.

At noontime, linear fits are still a good indication of the relationship between sunspot numbers and  $f_0F_2$ , and the maximum and minimum  $adj\_R^2$  are  $94.8\%$  and  $81.2\%$  in the months of March and July, respectively. It has been observed that  $f_0F_2$  shows a good linear relationship for low and moderate numbers of sunspots while it shows almost constant or decreasing values at extremely high solar epochs, which is known as the saturation effect. Saturation depends on local hours, months and on the magnitude of the solar cycle. In Figures 1 and 2 it is observed that, in winter, (including the months December and January) the  $f_0F_2$  values show generally constant values of high sunspot numbers, which our first regression model does not demonstrate. Therefore, a better model is expected by using the second regression model, i.e., quadratic relationship between sunspot numbers and  $f_0F_2$  as,

$$foF2(t)_m = b_{0m}(t) + b_{1m}(t) \cdot SSN_m + b_{2m}(t) \cdot SSN_m^2 \quad \dots(3)$$



**Figure 1:** The relationship between monthly median  $f_0F_2$  and monthly mean SSN at JST 0 h for the whole interval of 1976 – 1986 and 1996 – 2008 over Wakkanai. Observed values are plotted as dots, dotted lines represent the linear regression fit while solid lines show the quadratic regression fit.

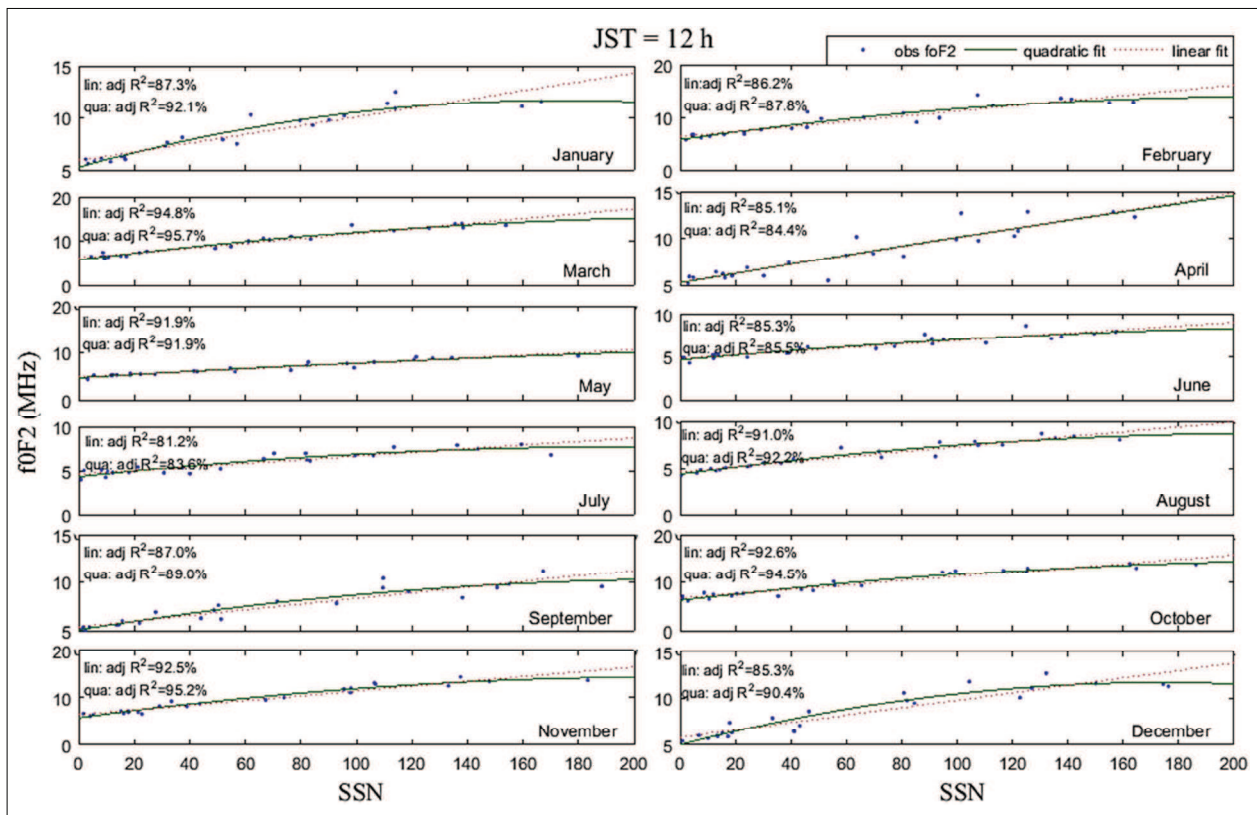
$b_0$ ,  $b_1$  and  $b_2$  are the coefficients of the quadratic regression equation at specified JST time ( $t$ ) and month ( $m$ ) represented by a solid line in Figures 1 and 2. In this case the significance of coefficient  $b_2$  is;

- i)  $b_2 < 0$ , indicates that  $f_0F_2$  declines with increasing solar activity, which represents the saturation effect.
- ii)  $b_2 > 0$ , implies that  $f_0F_2$  rises at excessively high solar activity epochs.

The saturation at noon and midnight in different months is clearly shown by the second regression model in Figures 1 and 2. The maximum difference in  $adj\_R^2$  (i.e., 5 and 10 %) is found in December at noon and midnight, respectively. The clear saturation in the noontime monthly median values of  $f_0F_2$  during the winter season (including months of December and January) at Wakkanai (45.39 °N) is due to the enhancement of the recombination process. Molecular nitrogen ( $N_2$ ) plays a

key role in the recombination process (Rishbeth, 1998). Hedin *et al.* (1979) demonstrated that the density ratio of atomic oxygen and molecular nitrogen ( $O/N_2$ ) decreases significantly with the increase in solar activity for winter at 45° latitude. This is due to the increase in the scale height of molecular species, especially  $N_2$ , rather than the scale height of atomic oxygen at high solar epochs (Schunk & Walker, 1973). This anomalous role of  $N_2$  raises the recombination process, emerging as a saturation effect.

Polynomial regression models of higher orders have been tested, and a good relation between  $f_0F_2$  and SSN at low and moderate sunspot numbers was found. However, randomly steep downward or upward trends at larger sunspot numbers contradicts the actual physical phenomena of the saturation effect. Along with this, the adjusted  $R^2$  also showed no improvement with higher orders.



**Figure 2:** The relationship between monthly median  $f_0F_2$  and monthly mean SSN at JST 12 h for the whole interval of 1976 – 1986 and 1996 – 2008 over Wakkanai. Observed values are plotted as dots, dotted lines represent the linear regression fit while solid lines show the quadratic regression fit.

It has been shown that for lower levels of solar activity,  $f_0F_2$  may vary during ascending and descending parts of the 11-year solar cycles; an effect which is known as the hysteresis effect. This effect is attributed to the activities of geomagnetic storms throughout the solar cycle (Rao & Rao, 1969; Apostolov & Alberca, 1995; Mikhailov & Mikhailov, 1995). The enhanced geomagnetic activity during the declining period of the solar cycle can produce stronger storm effects on the  $F_2$  layer than during the increasing period of the solar cycle. Kane (1992) argued that two key parameters i.e., solar and the geomagnetic index, should be taken into account for long-term prediction models. Therefore, Xu *et al.* (2008) used a multiple regression model by taking two parameters i.e., the sunspot number as solar activity index and geomagnetic index  $A_p$  for geomagnetic perturbations of the dependence on the monthly median  $f_0F_2$  over Chongqing, China.

Our analysis of  $f_0F_2$  over Wakkanai (45.39 °N) supports the model by Kane (1992) that uses both solar and geomagnetic index parameters which improved the regression model. Our third regression model is represented as dashed-square lines in Figure 3 and expressed as,

$$\begin{aligned} f_0F_2(t)_m = & c_{0m}(t) + c_{1m}(t) \cdot SSN_m + c_{2m}(t) \cdot SSN_m^2 \\ & + c_{3m}(t) \cdot A_{pm} \cdot SSN_m + c_{4m}(t) \cdot A_{pm} \\ & + c_{5m}(t) \cdot A_{p^2} \end{aligned} \quad \dots(4)$$

where  $c_0$  to  $c_5$  are the coefficients of the third regression model in a different JST time ( $t$ ) and month ( $m$ ) while  $A_p$  and SSN are the monthly mean geomagnetic index and sunspot number at the present month, respectively. The coefficients  $c_1$  and  $c_2$  are the quantitative expressions of the current month's sunspot number whilst  $c_4$  and  $c_5$  represent the current month's  $A_p$  index values. However,  $c_3$  depicts the combination of solar and geomagnetic disturbances.

The standard deviations of the change in  $f_0F_2$  calculated from regression models and observed values for the years 1976 – 1986 and 1996 – 2008 are illustrated in Figure 3. The diurnal variations of these standard deviations are determined for all months. It is seen that the trend is slightly nearer to zero when using both factors (i.e., solar and geomagnetic index) for the empirical model represented as dashed-square lines. It concludes that the third regression model better explains the observed values rather than the second regression

model. The root mean square deviation (RMSD) is often used to measure the difference between two time series trends; RMSD is calculated from the trend of the standard deviations measured from equations (3) and (4). In this scenario, the higher values of RMSD in a specific month indicate that the calculated  $f_0F_2$  values from equation (4) are much closer to the observed values than the  $f_0F_2$  calculated from equation (3) for that particular month. The maximum RMSD was reported in the months of January and July (i.e., 0.049 and 0.0488) and the maximum differences appear at 11:00 and 09:00 hours JST, respectively. Therefore, the third regression model is preferable over the second regression model. However, it is insignificant that in the months of April, May and September, where the RMSD values are the least i.e., 0.0089, 0.017 and 0.0122, respectively.

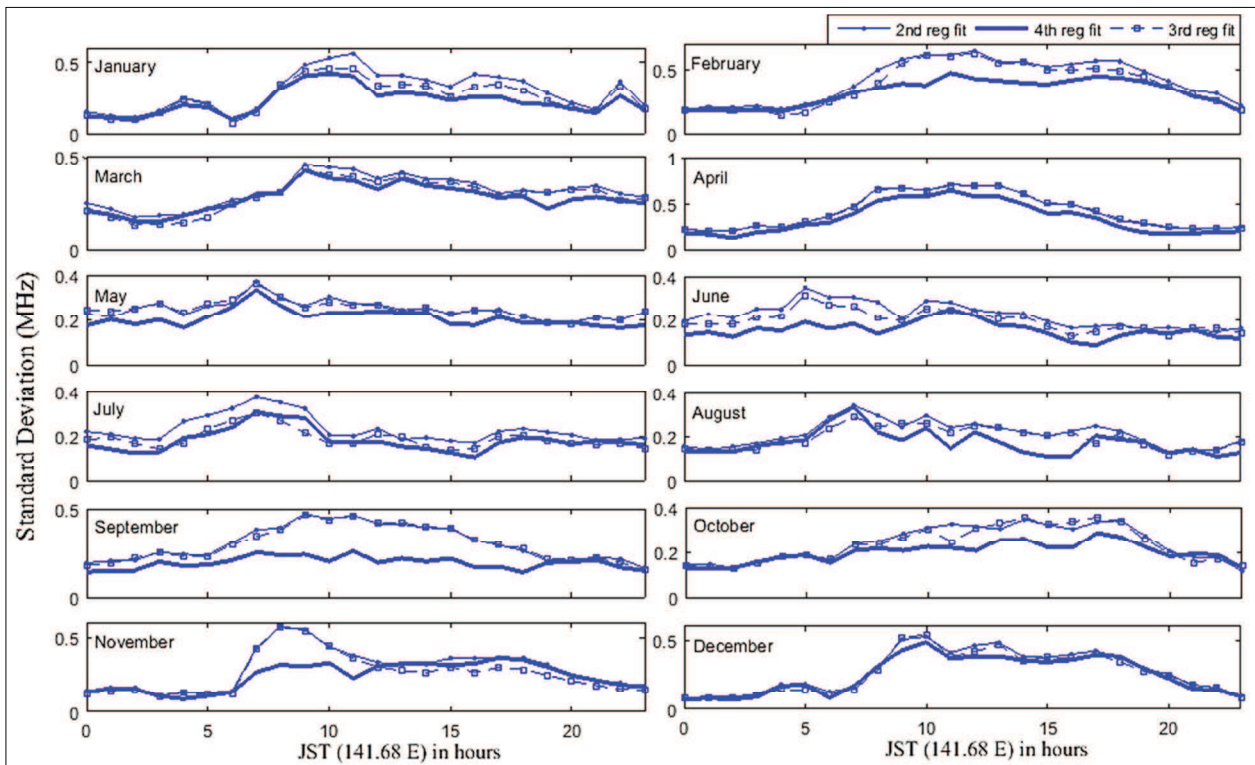
The hysteresis effect has a great impact on the mid-latitude compared to the low and high-latitudes, which cannot be neglected for empirical and spectral models of  $f_0F_2$ , especially at Wakkanai (45° N). Rao and Rao (1969) observed that noontime  $f_0F_2$  values are higher in the decreasing part of the solar cycle while it is low in the rising phase of the solar cycle. This is the consequence of some weak solar activities which can occur with or without sunspot formation because the required magnetic field threshold to form sunspot is 1500 Gauss (Penn & Livingston, 2006). However, the weak magnetic field can produce events such as solar flares, coronal mass ejection, solar winds and energetic protons. Those events occur near or close to the time of formation of sunspots (Legrand & Simon, 1989). High-speed solar wind streams from low solar-latitude coronal holes enhanced much later than sunspots maxima in the falling part of a solar cycle whose effect on ionosphere is even more non-linear (Simon & Legrand, 1986). Therefore in a systematic way, one can consider the rate of change in sunspot numbers for a few-month's time interval from a given month. This approach can potentially cover all the solar activities that occur before and after the formation of sunspots as shown previously by Pancheva and Mukhtarov (1998). Pancheva and Mukhtarov (1998) found that the rate of change in solar activity is low in the descending part of the solar cycle than in the ascending part. The influence of the tendency for a change of solar activity is essential for the years when solar activity varies very rapidly (i.e.; ascending part of the solar cycle). The study (Pancheva & Mukhtarov, 1998) used the Kr parameter and tried to get more accuracy in empirical models, although it is inconvenient because the Kr parameter includes the factor of solar activity after five months from the given month, which may be unknown. Then a new finding was published by Liu *et al.* (2004) in which they only replaced

Kr with F107p (i.e., the prior month's F107 values) as a solar activity parameter over the Wuhan region (30.6 °N, 114.4 °E) that gave an even better result. In the present study, SSNp was used (i.e., the monthly mean sunspot number of the previous month). Therefore, the fourth regression model can be given as,

$$foF2(t)_m = d_{0m}(t) + d_{1m}(t).SSN_m + d_{2m}(t).SSN_m^2 + d_{3m}(t).SSNp_m + d_{4m}(t).SSNp_m^2 \dots(5)$$

The trend of standard deviations by evaluating equation (5) is represented as a solid line in Figure 3. For all months, it is seen that the fourth regression model has a much lower standard deviation that made it easier to explain

the observed  $f_0F_2$  values over Wakkanai. The highest drop in the standard deviation occurs in November from about 0.57 MHz to 0.31 MHz at 08:00 hour JST. For all months, RMSD is computed from the standard deviation trends measured from equations (3) and (5). The RMSD has high values, which is an indication that the best relationship exists between the calculated  $f_0F_2$  from equation (5) with observed  $f_0F_2$ , especially in February and September with maximum RMSD value of 0.114 and 0.136, respectively. Introducing SSNp improves the depiction of solar cycle dependence of  $f_0F_2$ . In other words, it is observed that the rate of change in solar activity has a significant effect on  $f_0F_2$  over Wakkanai and its response depends on the month and partially on the diurnal cycle as seen in Figure 3. The results demonstrate that the fourth regression model accurately explains the complex behaviour of the ionospheric  $f_0F_2$  parameter.



**Figure 3:** Diurnal variation of the standard deviation of  $f_0F_2$  computed from equations (3), (4) and (5) from the observed  $f_0F_2$  for the whole interval 1976 – 1986 and 1996 – 2008 over Wakkanai. Line dots represent the second regression fit, dashed-square lines show the third regression fit, solid lines are for the fourth regression fit.

**Spectral analysis**

We have constructed a diurnal  $f_0F_2$  model for Wakkanai 45.39 °N, based on Fourier series. Four key parameters

were used for the model. These are  $t$ ,  $m$ , SSN and SSNp, which represent the diurnal, seasonal and solar cycle variations and the effect of prior solar activity on  $f_0F_2$ , respectively, where  $t$  is the Japanese standard time,  $m$  is

the month, SSN and SSNp are the monthly mean sunspot numbers' value in the specified and previous month, respectively.

**Fourier model and its authentication**

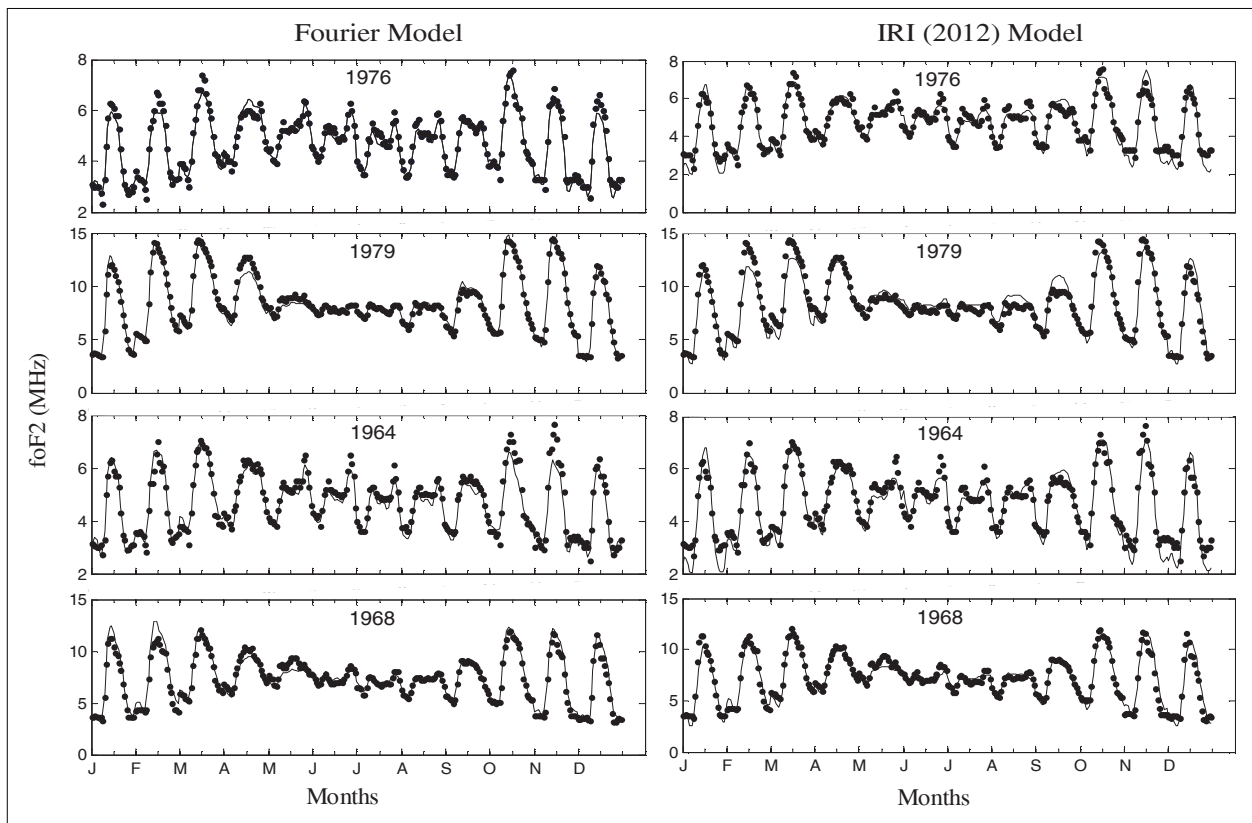
Fourier series represents periodic functions in terms of cosines and sines. There are diurnal periodicities present in the ionospheric parameters. Therefore, the Fourier series is generally preferred when constructing spectral models. The diurnal variations can be expressed by a Fourier series with periods of 24 hours and higher harmonics. A Fourier series for  $f_0F_2$  can be generally expressed as,

$$foF2(t)_m = x_{0,m} + \sum_{n=1}^N \left( x_{n,m} \cos\left(\frac{2\pi n t}{T}\right) + y_{n,m} \sin\left(\frac{2\pi n t}{T}\right) \right) \dots(6)$$

where the harmonic number is  $n$  ( $=0, 1, 2, 3, \dots, N$ ), and in our models, we take 4 harmonic numbers to ensure model accuracy and  $T$  is equal to 24 hours. The coefficients

$x_{n,m}$  and  $y_{n,m}$  are functions of SSN and SSNp. These coefficients are computed from equation (6) by using the curve fitting tool in MATLAB software, of the specified months for solar minimum and solar maximum years (i.e., 1976 and 1979, respectively); however, the  $f_0F_2$  values for every month are evaluated from equation (5).

The predicted  $f_0F_2$  values from the Fourier model [i.e., equation (6)] and IRI model are represented as solid lines, whilst the observed  $f_0F_2$  values are symbolised with dots as shown in Figure 4. Left panels are for the Fourier model while right panels are for the IRI-2012 model; both are compared with the observed  $f_0F_2$  values. The solar minimum year (1976) and solar maximum year (1979) were chosen for constructing the  $f_0F_2$  models. For verification of the models, we used the years 1964 and 1968, which are the solar minima and solar maxima years of solar cycle 20, respectively. As  $f_0F_2$  data were missing for the years of solar cycle 22 (1986–1996) and 24 (2008 – present), we were unable to work on it. The years used for verification of our models were not included initially in the dataset for regression models. Therefore, the

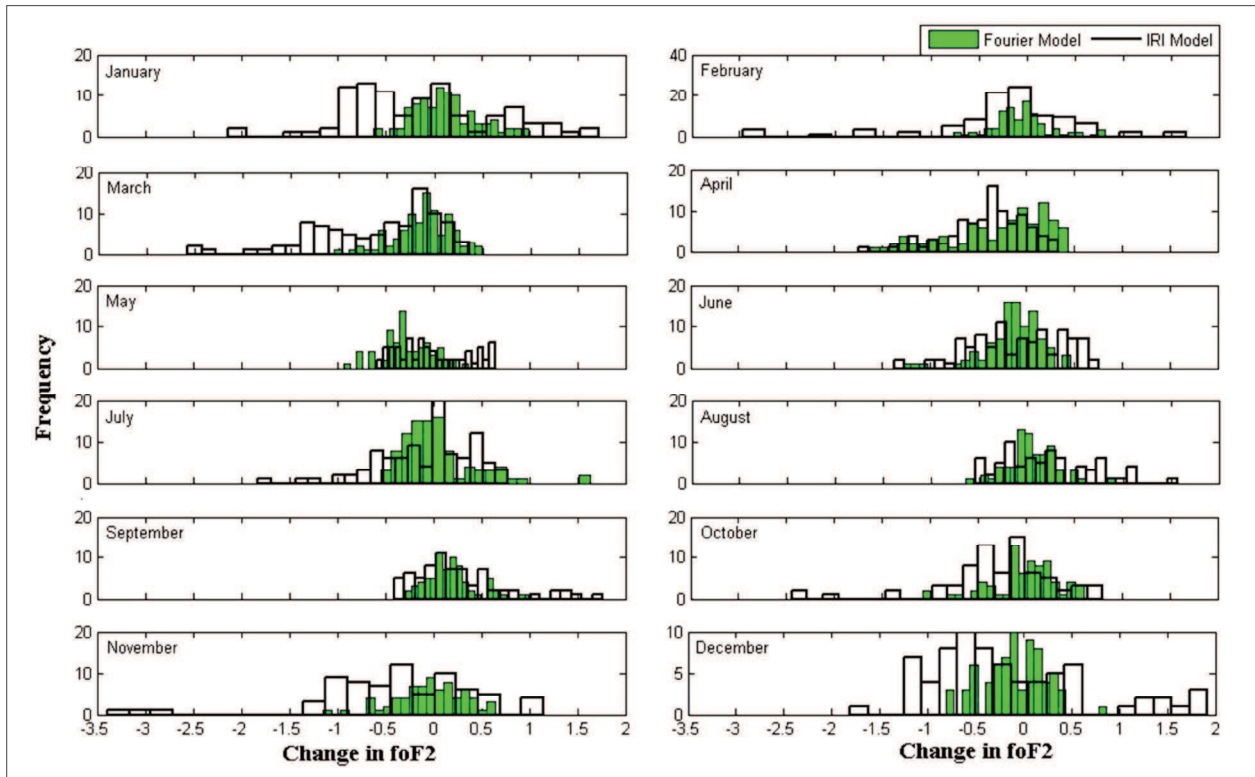


**Figure 4:** The computed and predicted diurnal variation of the monthly median  $f_0F_2$  represented as solid lines with observed  $f_0F_2$  shown with dots for solar minimum years (1976 and 1964) and for solar maximum years (1979 and 1968).

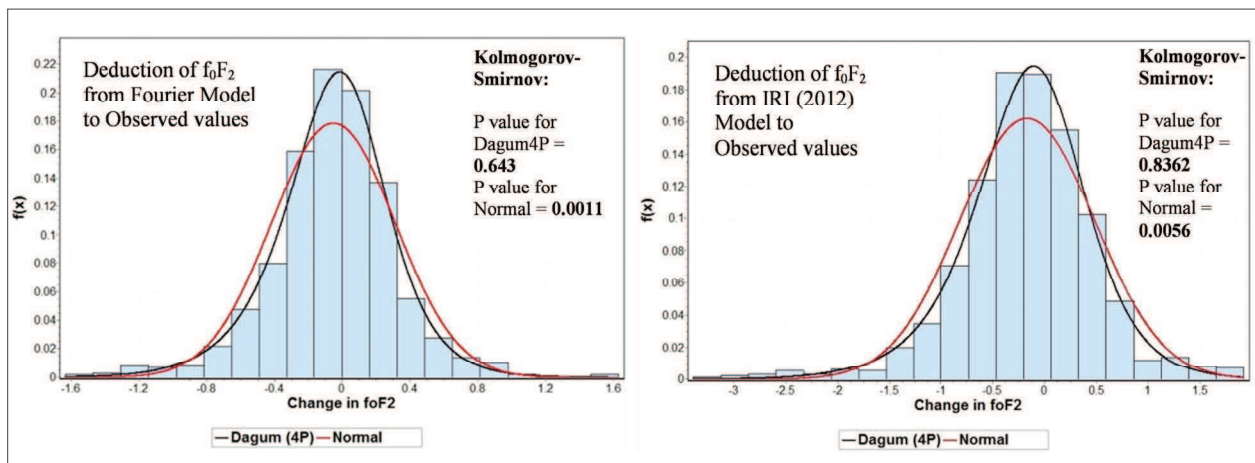


correctness of our work can be demonstrated by showing the results of the model  $f_0F_2$  values that almost do not change in the years which are not included originally in

the dataset for analysis. Hence our Fourier models can be employed for the long-term prediction of the ionospheric  $f_0F_2$  parameter over Wakkanai (45° N).



**Figure 5:** The deviation of computed and predicted (Fourier and IRI-2012)  $f_0F_2$  from observed  $f_0F_2$  for the years 1976, 1979, 1996 and 2000. Green colour bars are for the Fourier model and non-colour bars are for the IRI-2012 model.



**Figure 6:** The probability density function plot of the change in  $f_0F_2$  of the solar minimum and maximum years of the cycle 21 and 23. The Dagum 4P (in black colour) and normal (in red colour) distributions are fitted.

The deviation of models (Fourier and IRI-2012) from observed  $f_0F_2$  for the solar minimum (1976 and 1996) and maximum (1979 and 2000) years is shown in Figure 5. The missing data of  $f_0F_2$  of the year 2000 for the months of May and August to December are not included in the graph. Figure 5 demonstrates that:

- i) Firstly, the deviation of Fourier model  $f_0F_2$  mostly lays closer to zero as shown in the green colour bars, while divergence is much larger for IRI model  $f_0F_2$  as displayed in none colour bars. This illustrates that Fourier model better explains the real behaviour of F2 layer critical frequency than IRI-2012 model over Wakkanai.
- ii) Secondly, the estimation of  $f_0F_2$  by Fourier model in the months of April, June and October have more negative values that illustrate our model significantly and underestimate the observed  $f_0F_2$  values while in July and September, model overestimates the  $f_0F_2$  values. Although for all months, the scatter is much higher in the IRI-2012 model as compared to the observed  $f_0F_2$ .

The probability density functions (PDF) of the change in  $f_0F_2$  values are plotted in Figure 6, where the frequency on the y-axis shows the probability that x-axis can occupy between two maxima. It explained that high frequency peak near zero means observed  $f_0F_2$  values are close to the calculated (or IRI)  $f_0F_2$  values and low frequency shows opposite behaviour. The peak frequency near to zero is accurately represented by Dagum type II (4P) distribution that is also compared with the normal distribution. The Dagum 4P distribution is a heavy-tail distribution that contains four parameters; two parameters describe the shape, one is for scale and last is the location parameter. The Dagum 4P is frequently used in economic problems in which null and negative data have assets, and it is highly applicable in the cases of extreme data values that are not supposed to be neglected. In our case, comparison of models gives high frequencies with long-tail that are fully explained by the Dagum 4P distribution besides normal distribution that partially explained peak frequency with overestimating frequency at the second-standard deviation of the mean (Kleiber, 2007; Benjamin et al., 2013). The p-value test has preferred the Dagum 4P distribution and aborted the normal distribution fitting. PDF reveals that the IRI model is scattered more and under estimate  $f_0F_2$  values from observed values with respect to the Fourier model.

As a worldwide empirical model, the IRI model is outstanding at many spatial locations in the Earth's ionosphere. However, the behaviour of the ionosphere is

extremely complex and it changes in different latitudes. We worked on single-station empirical models so the effects of solar activity on  $f_0F_2$  over Wakkanai 45 °N are highly considered that enabled us to precisely measure the saturation and hysteresis effects and include them in models. A similar methodology has been applied by Liu et al. (2004) over Wuhan 30.7 °N and by Xu et al. (2008) over Chongqing 29.56 °N, China. Therefore, the Fourier model provides much higher accuracy as compared to the IRI model over Wakkanai 45 °N.

## CONCLUSION

The non-linear relationship between  $f_0F_2$  and SSN depends on the time of a day, season-months and the magnitude of solar activity. However, a quadratic regression fit is still better for all months to explain the small variations and especially the saturation effects.

To reduce the hysteresis effect, integrated factors of geomagnetic index  $A_p$  and SSN are used, which improved the empirical relationship. However, the relation is still lacking, especially in the months of April, May and September, while introducing prior month sunspot numbers (i.e., SSNp) with current month SSN significantly improved the relationship between solar cycle variation and its influence on  $f_0F_2$  over Wakkanai. The tendency of change in solar activity factor makes the empirical model more appropriate than other regression models and it is suitable to use equation (5) for the spectral models of  $f_0F_2$ .

Spectral models are constructed using Fourier series for the solar minimum and solar maximum years (1976 and 1979). A good agreement has been found between the computed  $f_0F_2$  from our models and the observed  $f_0F_2$ . We applied our models to the years 1964 and 1968 showing the solar minimum and solar maximum years. Comparable predictions of  $f_0F_2$  from the IRI-2012 model were also analysed. From Figures 5 and 6, it can be seen that the accuracy of our Fourier model is better than the IRI-2012 model, while the IRI model largely underestimated the  $f_0F_2$  values from observed values. The implementation of the Dagum 4P distribution is appropriate to model high frequencies and extreme values. The Dagum distribution is an appealing choice for modelling extreme values when using maximum likelihood estimators.

The Fourier models can be used as a relevant tool to study the effect of sunspot numbers on the F layer of ionosphere during solar minimum and solar maximum periods and it can be also used as a reference for predictions of the  $f_0F_2$  parameter.

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