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A practical indicator for surface ocean heat and freshwater buoyancy fluxes and its application to the NCEP reanalysis data

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ABSTRACT

The buoyancy flux at the air/sea interface plays a key role in water mass transformation and mixing as it modifies surface water density and in turn drives overturning and enhances stratification. It is the interplay of these two independent heat and freshwater buoyancy flux components that is of central importance when analysing mechanisms of the ocean/atmosphere interaction. Here, a diagnostic quantity (Θ_B) is presented that allows to capture the relative contribution of both components on the buoyancy flux in one single quantity. Using NCEP reanalysis of heat and freshwater fluxes (1948–2009) demonstrates that Θ_B is a convenient tool to analyse both the temporal and spatial variability of their corresponding buoyancy fluxes. For the global ocean the areal extent of buoyancy gain and loss regions changed by 10%, with the largest extent of buoyancy gain during the 1970–1990 period. In the subpolar North Atlantic, and likewise in the South Pacific, decadal variability in freshwater flux is pronounced and, for the latter region, takes control over the total buoyancy flux since the 1980s. Some of the areal extent time series show a significant correlation with large-scale climate indices.

1. Introduction

The ocean's surface layer is stirred by the wind and undergoes a cycle of convection and restratification in response to changes in surface water density. The surface water density is modified by the surface density flux (Q_{ρ}) that is composed of a heat flux (F_T) scaled by the thermal expansion coefficient (α) , and a freshwater flux (F_S) scaled by the haline contraction coefficient (β) . Multiplying the sum of the scaled fluxes with the sea surface density (ρ_S) results in the surface ocean density flux (e.g. Schmitt et al., 1989; Zhang and Talley, 1998)

$$Q_{\rho} = \rho_{\rm S} (\alpha F_{\rm T} + \beta F_{\rm S}),$$

where

$$F_{\rm T} = -Q_{\rm net}/\rho_S c_{\rm p}$$

and

$$F_{\rm S} = (E - P) \cdot S/(1 - S/1000).$$

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In the above equations, $Q_{\rm net}$ is the net heat flux (negative if ocean loses heat), E-P the local evaporation minus precipitation balance (assuming no influence from run-off and ice sheet melting), $c_{\rm p}$ the specific heat of seawater and S the sea surface salinity.

Considering the gravitational acceleration g and a reference density ρ_0 , the surface density flux becomes a buoyancy flux $\mathcal{B} = -g \, Q_\rho/\rho_0$ in units $W \, kg^{-1}$. The minus sign indicates that a water parcel becomes less buoyant when it gains density. Therefore, \mathcal{B} drives overturning ($\mathcal{B} < 0$) or stratification ($\mathcal{B} > 0$) of surface waters and is thus ultimately related to the formation and destruction of water masses at the ocean surface (e.g. Walin, 1982; Tziperman, 1986). The distribution of \mathcal{B} has been studied in theoretical studies (Garrett et al., 1995; Nurser et al., 1999), for the North Atlantic basin (e.g. Schmitt et al., 1989; Speer, 1997), the Southern Ocean (e.g. Marshall, 1997; Speer et al., 2000; Badin and Williams, 2010), shelf seas (Badin et al., 2010, and references therein) and on global scales (e.g. Speer et al., 1995; Zhang and Talley, 1998).

 \mathcal{B} is a quantity that is typically calculated from atmospheric and oceanic parameters at their interface, the sea surface. As such, atmospheric and ocean variability both reflect on \mathcal{B} in a composite way, making \mathcal{B} an important parameter to analyse variability of the atmosphere/ocean system. Changes in atmospheric forcing over the oceans, for example, as a response to

338 Tellus 63A (2011), 2

climate variability, can influence the air/sea exchange and thereby leave an imprint in the ocean. Likewise, changes in ocean circulation, for example, the intensity of the northward residual circulation in the Southern Ocean (Badin and Williams, 2010), may influence the air/sea temperature difference, and thereby the heat flux and buoyancy flux accordingly.

Although \mathcal{B} , over a larger part of the ocean area, is dominated by its thermal component (\mathcal{B}_T), it is the interplay of \mathcal{B}_T with the haline component (\mathcal{B}_S) that ultimately determines the impact of atmospheric and ocean changes on the ocean ventilation and subsequently the interior ocean structure. (Note: buoyancy flux is not the single driver for interior ocean water mass renewal, and processes, such as mixing, act as transmitters of surface information to propagate into the ocean's interior.)

Most studies on surface ocean buoyancy forcing put a focus on \mathcal{B}_T . One possible reason is that seasonal buoyancy flux variability is often controlled by heat flux variability, including the changes in the thermal expansion coefficient (Large and Nurser, 2001). However, it is well established from observations (e.g. Fairbanks, 1989) and model experiments (e.g. Vellinga and Wood, 2002) that \mathcal{B}_S is crucial for long-term variability of the atmosphere/ocean system and as such it is mandatory to consider both, heat and freshwater forcing.

Methods to conduct a joint analysis of spatially distributed time series of heat and freshwater buoyancy fluxes are rather limited: Either one has to apply a temporal averaging or an areal averaging to analyse the spatial distribution and temporal evolution of the two flux components, respectively. A different approach is analysing the flux ratio $(\mathcal{B}_T/\mathcal{B}_S;$ Zhang and Talley, 1998; Schmitt et al., 1989). The advantage of a ratio is that it combines the two flux components into a single number which opens the way for a joint analysis of the two components. But a ratio has obvious disadvantages for description and analysis: (1) a ratio is ambiguous, in case both contributors are either negative or positive the ratio is always positive—although the total flux is fundamentally different $(\mathcal{B} < 0 \text{ versus } \mathcal{B} > 0)$. (2) The ratio goes to infinity in case the flux in the denominator goes to zero, while the total flux can still be well defined. (3) The flux magnitude of either one or the other components is scaled differently as the numerator is the main contributor for ratios between -1 and 1, and the denominator is the main contributor for ratios between $\pm \infty$ and ± 1 .

Here we introduce a diagnostic that considers the relative contribution of the heat and freshwater buoyancy flux components but require neither temporal nor spatial averaging and also overcomes the disadvantages of a ratio. The diagnostic uses calculating the phase angle of the flux components as $atan2(\mathcal{B}_T, \mathcal{B}_S)$. As Fig. 1 illustrates, the atan2 is a two-argument function, available in most mathematical computation packages, that computes the angle between the positive x-axis of a plane (here \mathcal{B}_S , Fig. 1) and the point (x, y) (here $(\mathcal{B}_S, \mathcal{B}_T)$ on it. [Note that for most software packages the input arguments are defined as such that the y-axis is the first input argument [atan2(y, x)].]

For practical reasons it is convenient to associate a negative angle with a negative total buoyancy flux (which is similar to a buoyancy loss) and therefore the phase angle from the *atan*2

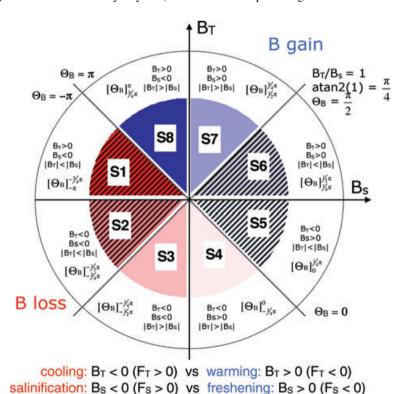


Fig. 1. Schematic of buoyancy flux angle, Θ_B , relative to the buoyancy flux components (\mathcal{B}_T and \mathcal{B}_S). The numbers indicate the eight segments S1–S8 as referenced in Table 1 and in the text. Redish segments indicate total buoyancy loss, \mathcal{B}_S dominated segments are hatched. In each segment, the respective \mathcal{B}_T , \mathcal{B}_S composition and Θ_B ranges are given.

Segment	Buoyancy loss				Buoyancy gain			
	S1	S2	S 3	S4	S5	S6	S 7	S8
	$-\pi$	$-3\pi/4$	$-\pi/2$	$-\pi/4$	0	$\pi/4$	$\pi/2$	$3\pi/4$
Θ_B	to	to	to	to	to	to	to	to
	$-3\pi/4$	$-\pi/2$	$-\pi/4$	0	$\pi/4$	$\pi/2$	$3\pi/4$	π
\mathscr{B}_{T}	>0	<0	<0	<0	<0	>0	>0	>0
\mathscr{B}_{S}	<0	<0	<0	>0	>0	>0	>0	<0
$ \mathcal{B}_{T} > \mathcal{B}_{S} $	No	No	Yes	Yes	No	No	Yes	Yes
Global area (%)	6	7	25	11	6	10	20	15

Table 1. Eight segments of Θ_B and associated sign and ratios of the thermal (\mathcal{B}_T) and haline (\mathcal{B}_S) buoyancy flux components (compare also Fig. 1). Last row give the percentage of area occupation in the long-term mean as derived from the NCEP reanalysis data.

function must be shifted by $+\pi/4$ (see Fig. 1). Because of the shift, absolute jumps greater than π should be shifted to their 2π complements. In the following we shall refer to the shifted phase angle as 'buoyancy flux angle'

$$\Theta_B = atan2(\mathscr{B}_{\mathrm{T}}, \mathscr{B}_{\mathrm{S}}) + \frac{\pi}{4}.$$

 Θ_B divides in eight segments between $-\pi$ and π that represent all eight possible combinations of \mathcal{B}_T and \mathcal{B}_S (Fig. 1 and Table 1). Segments S1 to S4 represent all \mathcal{B}_T and \mathcal{B}_S combinations that in total lead to a buoyancy loss ($\mathcal{B} < 0$) of surface waters, while segment S5–S8 represent the combinations that lead in the sum to a buoyancy gain ($\mathcal{B} > 0$). The segments S1, S2 and S5, S6 are those where the haline buoyancy flux dominates the total buoyancy flux ($|\mathcal{B}_S| > |\mathcal{B}_T|$), and consequently the remaining segments indicate the dominance of the thermal flux. As Θ_B is an angular measure it is not meaningful to average Θ_B values, instead the \mathcal{B}_T and \mathcal{B}_S input data should be averaged (if required).

The following section discusses the global distribution of Θ_B . It will be demonstrated how Θ_B can be used to analyse and interpret temporal variability of \mathcal{B}_T and \mathcal{B}_S in a joint way. Conclusions are drawn in the last section.

2. Global distribution and variability of the buoyancy flux angle Θ_B

We make use of 62 yr (January 1948–December 2009) of monthly NCEP reanalysed atmospheric data (Kalnay et al., 1996) combined with observations time series of sea surface temperature (SST) for the same period (Reynolds et al., 2002; Smith et al., 2008) to calculate Θ_B . For sea surface salinity, a climatology (Levitus et al., 1994) was used as no time series with global coverage is available.

From the long-term mean of the monthly heat and freshwater buoyancy fluxes (Figs. 2a and b) we calculate the global distribution of Θ_B (Fig. 3). Interannual variability is calculated from annual averages of monthly \mathcal{B}_T and \mathcal{B}_S data. Because of the

non-linearity of α and β on temperature and salinity, respectively, it is important to calculate the buoyancy fluxes prior to doing the averaging.

2.1. Long-term mean Θ_R distribution

At first glance, the climate pattern of Θ_B (Fig. 3) shows many similarities with the distributions of \mathcal{B}_T and \mathcal{B}_S (Figs. 2a and b). One obvious difference is that regions with strong \mathcal{B}_T , like the western boundary current regions (see Fig. 2a), are not peculiar in Θ_B . By definition, Θ_B is uneffected by the magnitude of \mathcal{B}_T , and the dominance of \mathcal{B}_T over \mathcal{B}_S in the western boundary current regions continues eastward over large parts of the adjacent subtropical gyres. Likewise, the pattern of the ratio $\mathcal{B}_T/\mathcal{B}_S$ (Fig. 2c) show similarities with Θ_B (Fig. 3). However, one can identify regions were ambiguity appears and very large or even vanishing flux ratios appear. The interpretation of the time series will demonstrate the difficulties in operating with a non-symmetric scale.

Overall we find in the NCEP reanalysis data an equal partition of the surface ocean area into buoyancy loss (S1–S4) and buoyancy gain (S5–S8) regions (see Table 1, last row). Furthermore, 71% of the surface ocean is predominately controlled by the thermally driven buoyancy loss (S3, S4) and buoyancy gain (S7, S8), confirming the common global mean picture of two-thirds thermal and one-third haline surface buoyancy forcing (Speer et al., 1995, see Table 1).

Considering individual segments, 25% of the global ocean area is occupied by S3, indicating a net heat loss is enforced by net evaporation (compare also Figs. 2a and b). Large parts of the subtropical gyres are associated with S3 where the world oceans' Central Waters are being formed (Sverdrup et al., 1942). The second largest portion is S7, covering 20% of the global ocean area in the NCEP reanalysis data. This segment indicates a buoyancy gain of surface waters which stems from warming as well as freshening effects. S7 occupies a major part of the

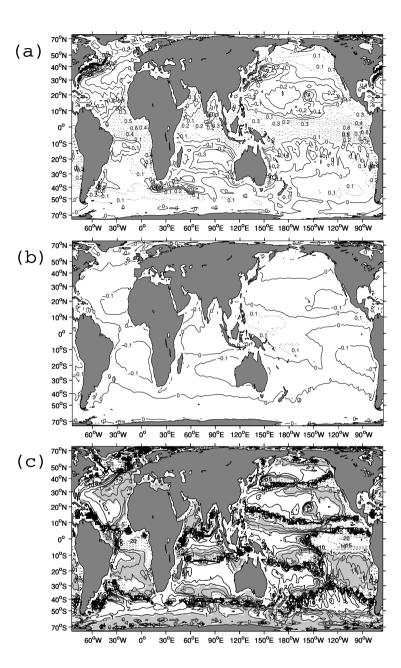


Fig. 2. Long-term mean distributions of (a) thermal buoyancy flux (\mathcal{B}_T) and (b) haline buoyancy flux (\mathcal{B}_S) . Note that the same contour interval was used in both panels (0.1 W kg^{-1}) , and dashed lines indicate buoyancy gain. (c) Buoyancy flux ratio $(\mathcal{B}_T/\mathcal{B}_S)$ where grey shading indicates freshwater flux dominates buoyancy flux, dashed lines indicate negative ratios.

tropical upwelling regions, parts of the North Pacific subpolar gyre and most areas in the transition region between the southern hemisphere subtropical gyres and the Antarctic Circumpolar Current.

S8 (net buoyancy gain) and S4 (net buoyancy loss) occupy about 15 and 11%, respectively, of the ocean's area. In both segments the haline buoyancy forcing opposes the thermal buoyancy forcing, whilst the latter has the larger amplitude.

S8 areas indicate that a net buoyancy gain $(\mathcal{B} > 0)$ occurs from the warming of surface waters $(\mathcal{B}_T > 0)$ but also that these regions lose buoyancy through net evaporation $(\mathcal{B}_S < 0)$. This is a typical situation for the eastern boundary upwelling regions

and for the equatorial 'cold tongue' regions in the eastern tropical Pacific and Atlantic Ocean.

S4 areas are of particular interest as here the net heat loss drives overturning of surface waters but a net freshwater gain is imprinted as a low salinity signature onto the overturning waters. This segment is found in the temperate regions of the subpolar North Atlantic and in the Mode and Intermediate Water formation areas at the poleward rim of the subtropical gyres (McCartney, 1977; Hanawa and Talley, 2001). Other regions with S4 characteristics are transition zones between the tropical and subtropical Pacific and Indian Oceans at about 10°N/S (e.g. Tomczak and Godfrey, 1994).

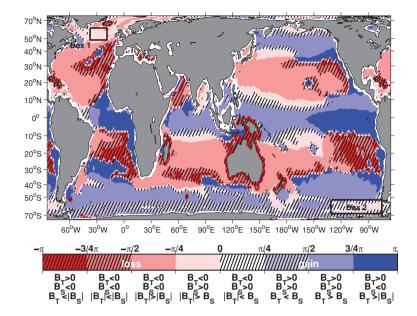


Fig. 3. Long-term mean distribution of buoyancy flux angle Θ_B . Boxes used for time series analysis are indicated by black frames.

In the area south of about 60° S, haline buoyancy fluxes dominate the buoyancy forcing of surface waters, as indicated by S5 (with opposing heat loss) and S6 (with accompanying heat gain). The net heat gain in S6 can be related to the northward advection of cold waters with the residual circulation (e.g. Speer et al., 2000; Badin and Williams, 2010) or by other processes, such as upwelling of cold interior waters.

Finally, the $\mathcal{B} < 0$ areas dominated by haline driven buoyancy loss (S1, S2) are the eastern parts of the subtropical gyres, outside of the coastal upwelling regions. In these areas intense evaporation ($\mathcal{B}_{\rm S} < 0$) controls the net buoyancy loss and thus sinking of surface waters. Depending on the local advection and mixed layer entrainment, the thermal buoyancy flux support (S2) or oppose (S1) the haline buoyancy loss. These areas are associated with the formation of saline Subtropical Underwater and eastern Subtropical Mode Waters (Mamayev, 1975; Emery and Meincke, 1986; Hanawa and Talley, 2001).

2.2. Analysis of surface ocean buoyancy flux time series by using Θ_B

Time series of buoyancy flux at the sea surface and the relative importance of its two components \mathcal{B}_T and \mathcal{B}_S can be systematically and simultaneously analysed by using Θ_B . Here we use time series of areal occupation of the eight segments, thereby eliminating the need for temporal or spatial averaging. This is in contrast to constructing time series of the individual fluxes or of their flux ratio which always require areal averaging. Three example areas are selected, the subpolar North Atlantic (Box 1), the South Pacific (Box 2) and the global ocean south of 70° N (see Fig. 3). The Θ_B analysis is presented and also the areal averaged \mathcal{B}_T , \mathcal{B}_S , and their ratio ($\mathcal{B}_T/\mathcal{B}_S$) for reference purposes.

The discussion of the time series includes the correlation with large-scale global climate indices as the North Atlantic Oscillation (NAO; Hurrell and Deser, 2009, December through March index based on the difference of normalized sea level pressure between Lisbon, Portugal and Stykkisholmur/Reykjavik, Iceland), the Southern Annular Mode (SAM; Marshall, 2003, observation based index), the Pacific Decadal Oscillation (PDO; Zhang et al., 1997, leading PC of monthly SST anomalies in the North Pacific Ocean, poleward of 20°N), and the East Atlantic pattern (EAP; Josey and Marsh, 2005, Rotated Principal Component Analysis applied to NCEP/NCAR reanalysis data).

Variability of heat and freshwater forcing in Box 1, in the subpolar North Atlantic, has been extensively discussed in the literature before (Reverdin et al., 2002; Josey and Marsh, 2005). Josey and Marsh (2005) analysed evaporation and precipitation anomalies in Box 1 for the period 1958 to 2002 (also using the NCEP reanalysis). They found not only a significant increase in net precipitation but a highly negative correlation with sea surface salinities. The EAP was found to be well correlated with the increase in precipitation while the NAO seems to play only a minor role (see also Reverdin et al., 2002). The subpolar North Atlantic is a major water mass formation area with deep reaching overturning of surface waters during winter, and therefore the freshening is not confined to the surface waters but also propagates into the interior ocean (Josey and Marsh, 2005; Curry and Mauritzen, 2005).

Time series of \mathcal{B}_T , \mathcal{B}_S and \mathcal{B} averaged over Box 1 (Fig. 4, upper left) show that in the annual mean net buoyancy loss ($\mathcal{B} < 0$) is dominated by \mathcal{B}_T (cooling). \mathcal{B}_S is one order of magnitude smaller and mostly opposite to \mathcal{B}_T (except for 1969, 1970). \mathcal{B}_T shows interannual variability with no obvious trend while \mathcal{B}_S may contain more of a decadal fluctuation with an overall

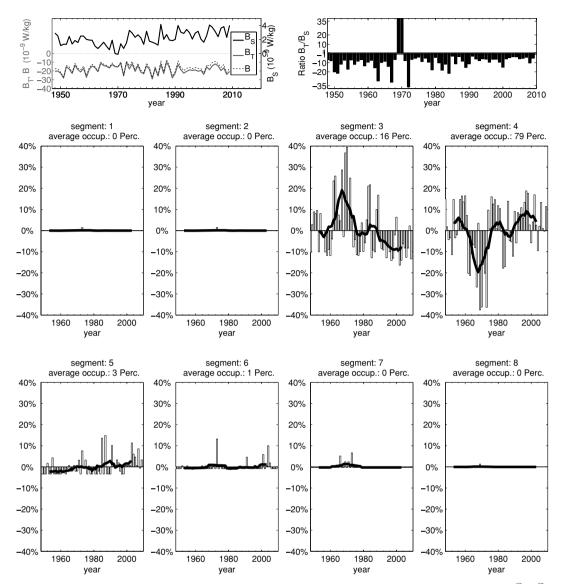


Fig. 4. North Atlantic subpolar gyre Box 1 (see Fig. 3): Time series (1948–2009) of (upper left) area averaged buoyancy fluxes \mathcal{B}_T , \mathcal{B}_S and \mathcal{B} (note difference in scales); (upper right) buoyancy flux ratio $\mathcal{B}_T/\mathcal{B}_S$; (lower eight panels) area occupation of the eight Θ_B segments S1–S8. The black line indicates the 11-yr running mean. The time averaged area occupation of each segment for the box is given in the title.

trend towards increasing buoyancy gain (freshwater gain), in accordance with earlier investigations (Reverdin et al., 2002; Josey and Marsh, 2005). The buoyancy flux ratio (Fig. 4, upper right) clearly shows for all years that the dominance of \mathcal{B}_T as $|\mathcal{B}_T/\mathcal{B}_S|$ is always greater 1. For the years 1969 and 1970, \mathcal{B}_S changed sign as evaporation dominated the freshwater flux in the areal average.

The time series of the eight Θ_B segments (Fig. 4, lower panels) also shows, as expected, the dominance of \mathcal{B}_T buoyancy loss as well as an opposing \mathcal{B}_S buoyancy gain (S4) in Box 1. In addition the Θ_B analysis shows that on average 16% of the area in Box 1 is under the influence of S3 (\mathcal{B}_S and \mathcal{B}_T are both negative) and that the variability of areal occupation of S3 and S4 are mirror

images of each other. What seem to be exceptional years (1969 and 1970) from the areal average \mathcal{B}_T , \mathcal{B}_S and their ratio show up as decadal variability with long-term trends in S3 and S4 in the Θ_B analysis. The positive (negative) correlation between the time series of segment S3 (S4) and EAP amounts to 0.3, significant at a 95% level. S4 is also positively correlated with the PDO by 0.3 (95% significance level). No significant correlation (on 95% level) between any segment and the NAO was found.

The second area (Box 2 in Fig. 3) examined in greater detail is the formation region of Antarctic Intermediate Water (AAIW) in the eastern South Pacific from 55°S to 70°S (McCartney, 1977, 1982; Hanawa and Talley, 2001). The areal average \mathcal{B}_S and \mathcal{B}_T buoyancy fluxes (Fig. 5, upper left) are of

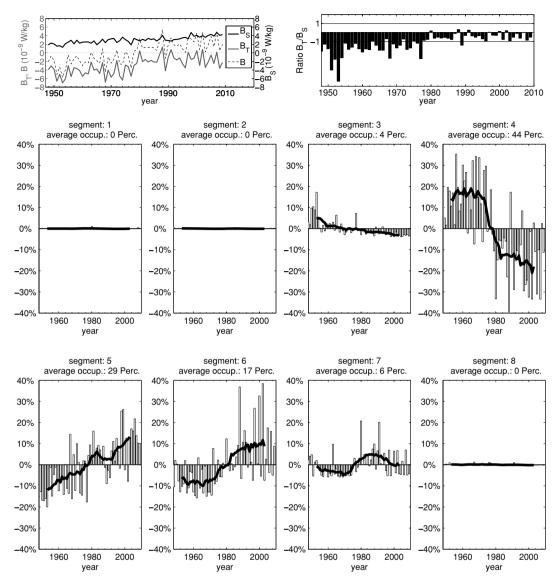


Fig. 5. As Fig. 4, but for South Pacific Box 2 (see Fig. 3).

comparable magnitude, except for some short periods of opposite sign $(\mathcal{B}_S > 0, \mathcal{B}_T < 0)$. \mathcal{B}_S nearly doubled in magnitude with a rather linear increase over the entire time series. Simultaneously the thermally driven buoyancy loss vanishes more and more, leading to a transition from a \mathcal{B}_T -dominated net buoyancy loss to a \mathcal{B}_S -dominated net buoyancy gain region. This is also reflected in the buoyancy flux ratio (Fig. 5, upper right) that switched from a \mathcal{B}_T dominance $(\mathcal{B}_T/\mathcal{B}_S < -1)$ for the 1950 to 1980s to a \mathcal{B}_S dominance $(\mathcal{B}_T/\mathcal{B}_S > -1)$ in the later decades. The positive correlation between anomalous annual $\mathcal{B}_S(\mathcal{B}_T)$ and SAM (PDO) amounts to 0.46 (0.39), both at the 99% significance level.

The Θ_B analysis shows that for Box 2 the area occupation in S4 decreases from more than 60% to about 20% while S5 (average 29%) and S6 (average 17%) occupation increases in

parallel (Fig. 5, lower panels). Based on repeated hydrographic observations, a warmer and more saline AAIW composition was found comparing the 1960s with recent years (Wong et al., 1999; Shaffer et al., 2000). The less dense water above the AAIW were found to be warmer and fresher (Boening et al., 2008).

Finally, the global buoyancy flux variability is shown in Fig. 6. Compared to the magnitude in \mathcal{B} variability in Boxes 1, 2, the magnitude of the global flux variability is small. The dominant source for \mathcal{B} variability are the fluctuations in \mathcal{B}_T whilst the variability of \mathcal{B}_S is about 10 times smaller (Fig. 6, upper left). During the first two decades (1950–1970) and the last two decades (mid 1990–2009), the \mathcal{B}_T buoyancy flux was anomalous high compared to the 1970–1990 decades.

The variability in global ocean areal occupation for the \mathcal{B}_{T} -dominated segments is 10%, which is surprisingly large. The

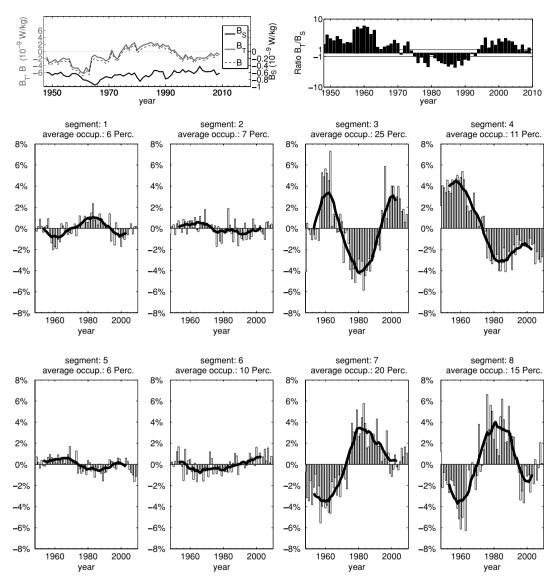


Fig. 6. As Fig. 4, but for the global ocean south of 70°N.

 \mathcal{B}_T and \mathcal{B}_S combinations in S3/S8 and in S4/S7 compensate each other, explaining why the magnitude of the area average flux is small (Fig. 6, upper left). For the transfer of air/sea flux variability into the ocean interior, not only the flux magnitude is important but the areal extent were net buoyancy loss occurs. The area occupation of Θ_B in S4 (S7) is positively (negatively) correlated with the PDO by 0.55 (both at the 99% significance level). Corresponding fluctuations in S3 and S8 are negatively and positively correlated, respectively, with the NAO at 0.31 and 0.55 (both at the 95% significance level).

A systematic shift during the 1960s/1970s as seen in S4/S7, has been identified in many atmospheric variables, such as rainfall, winds pattern, etc. (see Baines and Folland, 2007, for a summary). Baines and Folland (2007) identified the ocean/atmosphere system inherent internal variability to be the

dominante cause for the shift. But other factors, such as green-house gas emissions, have been identified as well (Meehl et al., 2009). However, it remains unclear how much inconsistencies in the NCEP data sets contribute to generating such variability. These inconsistencies may be caused by numerous reasons, including the general lack of data in certain regions, the sudden improvement of the observatory network (e.g. satellite remote sensing), or shortcomings in parametrizing the atmospheric boundary layer.

3. Conclusion

Buoyancy fluxes (\mathcal{B}) at the ocean surface represent a combination of the haline and thermal forcing and therefore a joint interpretation of the haline (\mathcal{B}_S) and thermal (\mathcal{B}_T) buoyancy

flux components is desirable when analysing \mathscr{B} variability. In the past a simple ratio of $\mathscr{B}_T/\mathscr{B}_S$ has been presented as a diagnostic tool (e.g. Schmitt et al., 1989; Speer et al., 1995; Zhang and Talley, 1998), but a ratio has disadvantages, such as ambigutey, infinity problems, and asymmetric scaling. As such a ratio is not well suited for quantitative diagnostics.

Here we present a diagnostic tool, called the buoyancy flux angle (Θ_B) , that combines \mathcal{B}_T and \mathcal{B}_S in a single quantity. Θ_B is divided into eight characteristic segments. Considering the areal occupation of each of the eight Θ_B segments, the temporal variability in \mathcal{B}_T and \mathcal{B}_S is analysed in a joint quantity.

NCEP reanalysis time series (1948–2009) of atmosphere and ocean surface parameters are used to illustrate the Θ_B analysis and to compare with results obtained by looking at \mathcal{B}_T and \mathcal{B}_S individually. Two important water mass formation regions of the global ocean are chosen for demonstration purposes: the subpolar North Atlantic (Box 1) and the eastern South Pacific (Box 2).

In Box 1, a pronounced freshening of surface waters has been reported and related to an increase in net precipitation (Reverdin et al., 2002; Josey and Marsh, 2005). The Θ_B analysis confirmed the freshening via a positive trend in area coverage of S4, which indicates that \mathcal{B}_T drives \mathcal{B} but is opposed by \mathcal{B}_S . Being positive (net freshwater gain) during most of the time, in 1969/1970 the areal averaged \mathcal{B}_S turned negative (evaporation exceeded net precipitation). By looking at the Θ_B time series (S3), we see that these 'exceptional' years were part of a multidecadal variability which is positively correlated with the Pacific Decadal Oscillation at 0.3. No significant correlation with the North Atlantic Oscialltion could be found for the S3 and S4 time series.

Box 2 is a key water mass formation area for Antarctic Intermediate Water. The area experienced a transition from being dominated by a \mathcal{B}_T driven net surface buoyancy loss to a \mathcal{B}_S driven net buoyancy gain. Θ_B analysis showed that the decrease in area with a thermal driven buoyancy loss (S4) was occupied in similar partitions from haline dominated buoyancy gain areas (S5 and S6). The area changes are correlated at -0.4 (S4) and +0.3 (S5 and S6) with the Southern Annular Mode.

For the global ocean area south of 70° N we found that the heat and freshwater buoyancy flux variability is small in magnitude, as expected when using a data set that has been forced to be nearly closed in global heat and freshwater fluxes. Nevertheless, Θ_B analysis revealed pronounced changes in the surface ocean area occupation, in particular of thermal buoyancy flux dominated loss and gain areas. These changes are likely associated with corresponding changes in the ventilating water masses.

Especially for ocean model simulations, the heat and freshwater forcing is crucial for experimental results. Θ_B offers a simple way for analysing forcing data sets and allows to easier

identify inconsistencies in the forcing fields. Θ_B analysis can be particularly useful when investigating the output of coupled ocean/atmosphere models as all input data needed to calculate Θ_B is available.

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