

# Observational quantification of a total aerosol indirect effect in the Arctic

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## ABSTRACT

We use 6 yr of multisensor radiometric data (1998–2003) from the U.S. Department of Energy Atmospheric Radiation Measurement (ARM) program to provide an observational quantification of the short-wave aerosol first indirect effect in the Arctic. Combined with the previously determined long-wave indirect effect, the total (short-wave and long-wave) first indirect effect in the high Arctic is found to yield a transition from surface warming of  $+3 \text{ W m}^{-2}$  during March to a cooling of  $-11 \text{ W m}^{-2}$  during May, therefore altering the seasonal cycle of energy input to the Arctic Earth–atmosphere system. These data also reveal evidence of a first indirect effect that affects optically thinner clouds during summer, which may represent an additional negative climate feedback that responds to a warming Arctic Ocean with retreating sea ice.

## 1. Introduction

Aerosol–cloud interactions are a major source of uncertainty in current global climate change simulations that critically need direct observations to resolve (Penner et al., 2006). Northern high latitudes are showing some of the most dramatic manifestations of anthropogenic climate change (Stroeve et al., 2007) and recent work showed that springtime pollution aerosol act to enhance the downwelling long-wave radiation from Arctic clouds by altering their microphysics via a mechanism known as the first indirect effect ( $\text{IDE}_1$ ), which yields enhanced surface long-wave fluxes comparable to those caused by  $\text{CO}_2$  abundance increases (Lubin and Vogelmann, 2006; Garrett and Zhao, 2006). However, assessing the total  $\text{IDE}_1$  depends on the concurrent impact at short wavelengths, as the Arctic insolation increases from zero during winter to a maximum at the vernal equinox. Here we determine the short-wave component of the  $\text{IDE}_1$  in the Arctic utilizing the unique multiyear and multisensor data from the U.S. Department of Energy Atmospheric Radiation Measurement (ARM) program's North Slope of Alaska (NSA;  $71^\circ 19' 37.73''\text{N}$ ,  $156^\circ 36' 56.70''\text{W}$ ) site at Barrow (Stamnes et al., 1999). The result can serve as a benchmark climatological result against which climate model parametrizations can be tested.

## 2. Methods

Pyranometer data from the NSA are used to determine the cloud transmission  $T_c$ , which is defined here as the ratio of the measured downwelling short-wave flux (insolation) at the surface to the equivalent surface flux under cloud-free skies and background aerosol (with all other atmospheric conditions constant). Pyranometer measurements (Michalsky et al., 1999) are made by the Sky Radiometers on Stand for Downwelling Radiation (SKYRAD) instrument suite. The broad-band surface albedo ( $\alpha_{\text{sw}}$ ) measurements are derived by ratio of these measurements with the upwelling flux made by the Ground Radiometers on Stand for Upwelling Radiation (GNDRAD) suite. Data used are recorded at 20 s and averaged to 1 min. The data used in this study encompass the years 1998–2003. Both suites are standard ARM program measurements kept in continuous operation at the NSA. To compute the cloud-free downwelling surface flux,  $F_{\text{cf}}$ , we use a 179-band discrete-ordinates radiative transfer model (Stamnes et al., 1988) optimized for the Arctic atmosphere (Lubin and Vogelmann, 2007), which corrects for the annual cycle in Earth–sun distance and incorporates spectral surface albedo for various Arctic conditions (Grenfell and Perovich, 1984; Perovich et al., 2002).  $F_{\text{cf}}$  is estimated using the following combinations of spectral surface albedos that are determined based on the measured value of  $\alpha_{\text{sw}}$ : snow-covered sea ice for  $\alpha_{\text{sw}} > 0.91$ , a linearly weighted mixture of bare and snow-covered sea ice for  $0.65 \leq \alpha_{\text{sw}} \leq 0.91$ , a similar mixture of tundra and bare sea ice for  $0.14 \leq \alpha_{\text{sw}} < 0.65$  and tundra for  $\alpha_{\text{sw}} < 0.14$ . These approximations are appropriate for a coastal high latitude station, given the spatial variability in multiple photon reflection

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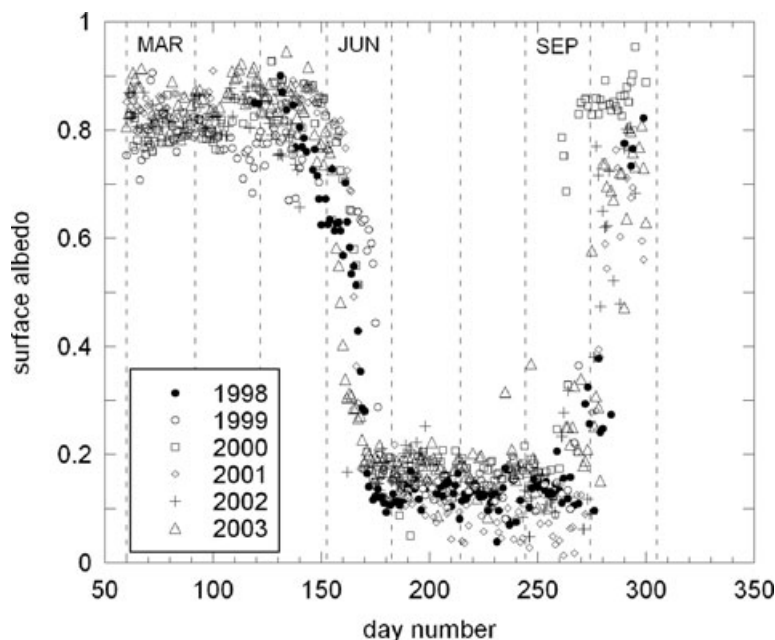


Fig. 1. Daily average broadband surface albedo measured at the ARM NSA site at Barrow, AK during 1998–2003. Vertical dashed lines identify the months.

effects for the full range of surface albedo (Podgorny and Lubin, 1998).

Before examining the differences between cloud transmissions under clean-air versus all other conditions, we need to account for any dependence on solar zenith angle  $\theta_0$ . Solar zenith angle differences as small as  $5^\circ$  can yield differences in  $T_c$  comparable to those we expect from the IDE<sub>1</sub>. Therefore, we examined the entire data set for solar zenith angle dependence, and derived a normalization  $T_c = T_{c0} + a(0.5 - \mu_0)$ , where  $T_{c0}$  is the uncorrected cloud transmission,  $\mu_0 = \cos\theta_0$ , and from a linear regression  $a = 0.514, 0.538, 0.473$  and  $0.620$  for March–May, June, July–August and September–October, respectively. With the  $T_c$  values thus normalized to  $\theta_0 = 60^\circ$ , we can search for a short-wave manifestation of the IDE<sub>1</sub> without worrying about possible artefacts from systematic differences in mean solar zenith angle.

Figure 1 shows the seasonal cycle in daily average  $\alpha_{sw}$ . The albedo is uniformly high during spring until late May, when it begins to decrease with the onset of the melt season. The timing of the melt season and  $\alpha_{sw}$  decrease varies considerably from one year to the next during June, as does the timing of the autumn freeze-up and  $\alpha_{sw}$  increase during September and October. Based on this seasonal cycle in  $\alpha_{sw}$ , we subdivide our analysis of IDE<sub>1</sub> into four time periods, (a) March–May, (b) June, (c) July and August and (d) September and October. During June, we only interpret the  $T_c$  determinations for which  $\alpha_{sw} > 0.4$ , when there are a sufficient number of data points.

The subset of Arctic clouds sampled herein was identified as being single-layered by the ARM Active Remotely-Sensed Cloud Locations algorithm (ARSCL; Clothiaux et al., 2000).

Figure 2 shows  $T_c$  as a function of ARSCL-measured cloud base altitude. As expected, higher clouds are generally optically thinner. However from March to July, clouds with bases in the lowest part of the boundary layer have on average a slightly smaller optical thickness (and larger  $T_c$ ) than those just above. The ARSCL data also reveal (not shown) that these lowest clouds have a tendency to be geometrically thinner than those just above, where 75% of them have a geometrical extent less than 500 m versus 65% for the clouds just above. Thus, with so many of these clouds contained within the boundary layer and also having relatively lower optical thickness, we can expect some degree of susceptibility to cloud nucleation (and concomitant changes in cloud radiative properties) from aerosol contained within the boundary layer. With the available data, only clouds having cloud base and cloud thickness both under 1000 m are used for statistical analysis of the IDE<sub>1</sub>. (Due to various gaps in the pyranometer data set and the location's cloud climatology, there are considerably fewer determinations of  $T_c$  for higher clouds than for those below 1000 m.) This restriction of our IDE<sub>1</sub> analysis to the lowest clouds favours liquid water during the Arctic spring and summer. Figure 3 shows the distribution of base height and geometrical thickness for the clouds considered in our IDE<sub>1</sub> analysis. Most of these clouds are likely to be within the boundary layer, even during summer.

Measurements of downwelling zenith middle-infrared radiance from the Atmospheric Emitted Radiance Interferometer (AERI) (Knuteson et al., 2004) are used to retrieve cloud effective droplet radius,  $r_e$  and liquid water path,  $LWP$ , using established radiative transfer algorithms (Guo et al., 2005; Lubin and

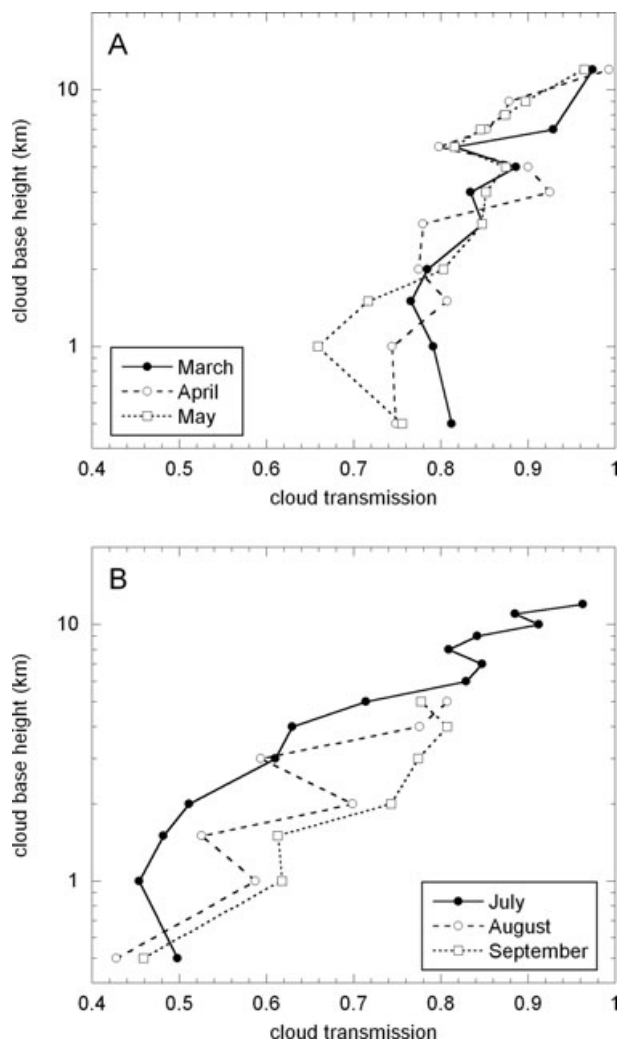


Fig. 2. Average cloud transmission  $T_c$  derived from the 1998–2003 NSA pyranometer data for single-layered clouds, as a function of cloud base altitude, during (A) spring and (B) summer and early autumn. Results for June are similar to those of May, and are omitted for clarity.

Vogelmann, 2006). AERI-retrieved  $LWP$  and  $r_e$  are considered spurious if  $LWP > 50 \text{ g m}^{-2}$ , due to lack of spectral variability in mid-IR cloud emission (Turner, 2005). For  $LWP > 50$ ,  $r_e$  is not retrieved and  $LWP$  is specified from the microwave radiometer (MWR) data (Westwater et al., 2001). MWR data are, in turn, not reliable for smaller instantaneous  $LWP$  values due to large radiometric uncertainty (Westwater et al., 2001). To eliminate potential discontinuity at the merger of these two  $LWP$  retrievals, the MWR-retrieved  $LWP$  was regressed against the AERI-retrieved  $LWP$ . This revealed a bias and slope of  $12.841 \text{ g m}^{-2}$  and  $1.092$ , respectively, in the MWR retrievals as compared to the AERI retrievals, and the bias was subtracted from all MWR  $LWP$  values.

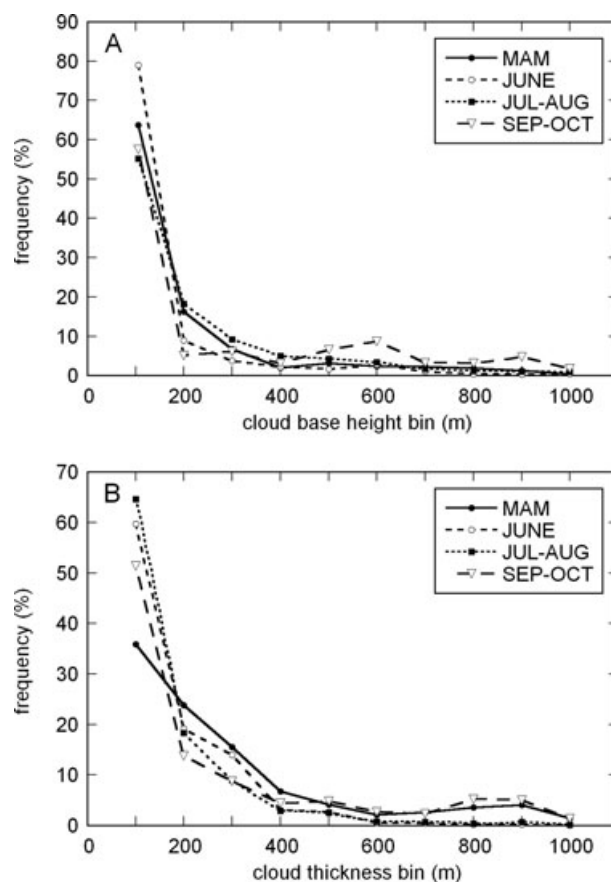


Fig. 3. Histograms of cloud base height (A) and cloud geometrical thickness (B), as determined from ARSCL data, for the cloud cases considered in this study.

### 3. Results

Several observational and theoretical studies have shown that anthropogenic aerosol in the Arctic troposphere, when entrained in clouds, decreases the effective radius  $r_e$  of cloud liquid water droplets (Morrison et al., 2005; Garrett and Zhao, 2006; Lubin and Vogelmann, 2006). The common occurrence of this  $\text{IDE}_1$  for optically thin, low-level clouds at Barrow, AK can be seen using AERI retrievals of  $r_e$  and  $LWP$ . We collocate these retrievals with hourly-averaged measurements of aerosol condensation nuclei (CN) concentrations at Barrow (Delene and Ogren, 2002). During spring, CN can often serve as a proxy for cloud condensation nuclei (CCN) concentrations (e.g. Garrett et al., 2004). During summer, we do not expect a straightforward relationship between CN and CCN, as new particle formation by nucleation can yield large CN concentrations (of order  $100\text{--}1000 \text{ cm}^{-3}$ ) that do not contribute to CCN, even though CCN concentrations of order  $100 \text{ cm}^{-3}$  can be observed sometimes during the Arctic summer (Garrett et al., 2004). In the absence of a clear relationship between CN and CCN during summer, we use the surface-measured CN concentrations as a simple discriminator

between pristine clean-air conditions (no measurable  $\text{IDE}_1$  expected) and all other conditions, which may or may not contain CCN (and therefore may or may not exhibit an  $\text{IDE}_1$ ). Here the natural background range is taken to be  $0\text{--}50\text{ cm}^{-3}$ . Although higher CN does not necessarily indicate high CCN during summer, we make the defensible assumption that, for low clouds, background CN concentrations implies background CCN concentrations.

Figure 4a shows that, in nearly all cases,  $r_e$  values decrease with increasing CN concentration above the natural background range of  $0\text{--}50\text{ cm}^{-3}$ . Because this result prevails when we hold  $LWP$  nearly fixed (open circles in 4a), it is consistent with the  $\text{IDE}_1$  (Twomey, 1977). The  $\text{IDE}_1$  appears in optically thin clouds during the Barrow spring (Fig. 4a) and summer months (Figs. 4b and 4c), but is not perceptible during autumn (Fig. 4d). The summertime result is surprising given that the aerosol size distribution in the high Arctic typically exhibits a seasonal shift from having an accumulation mode dominate the distributions in spring and autumn to having a nucleation/Aitken mode dominate the summertime distribution (e.g. Garrett et al., 2004; Heintzenberg et al., 2006; Engvall et al., 2008), and therefore the relationship between surface-measured CN and CCN should be much less clear, as discussed above. Nevertheless, our result of Figs. 4b and c for the summer months appears when using two completely independent measurements—a surface particle counter for CN and spectroscopic remote sensing retrieval for  $r_e$ —and remains to be explained.

One possibility is that this result is merely an artefact of air mass correlation, which has been shown to confound remote sensing-based identification of the  $\text{IDE}_1$  (e.g. Mauger and Norris, 2007). If differences in CCN are mainly due to differences in prevailing wind direction and aerosol origin, and retrievals obtained from clouds in a moist air mass with high CCN (one wind direction) are arbitrarily juxtaposed with retrievals obtained from clouds in a less moist air mass with lower CCN (another wind direction), then an artificially strong  $\text{IDE}_1$  will appear. We attempted to identify possible air mass correlation using the available NSA data. From surface meteorological observations in July and August, we identified two prevailing wind regimes, a westerly regime with wind directions  $240^\circ\text{--}320^\circ$  and an easterly regime with wind directions  $20^\circ\text{--}100^\circ$ . The former is accompanied by higher surface temperatures and lower surface pressures than the latter. Together, these regimes comprise 83% of the data containing the low clouds considered in this study. Figure 5 shows the observed inverse relationship between retrieved  $r_e$  and CN, for July and August, appearing for both wind regimes. The easterly regime overlaps with the preferred ‘clean air’ direction for air sampling at Barrow by NOAA (Polissar et al., 2001). Measurements made during westerlies might conceivably involve local pollution artefacts from Barrow, Alaska although, as a precaution, data points with wind direction within  $\pm 10^\circ$  of the direction of Barrow (2 km distant) have been omitted from analysis. There appear to be systematic differences in

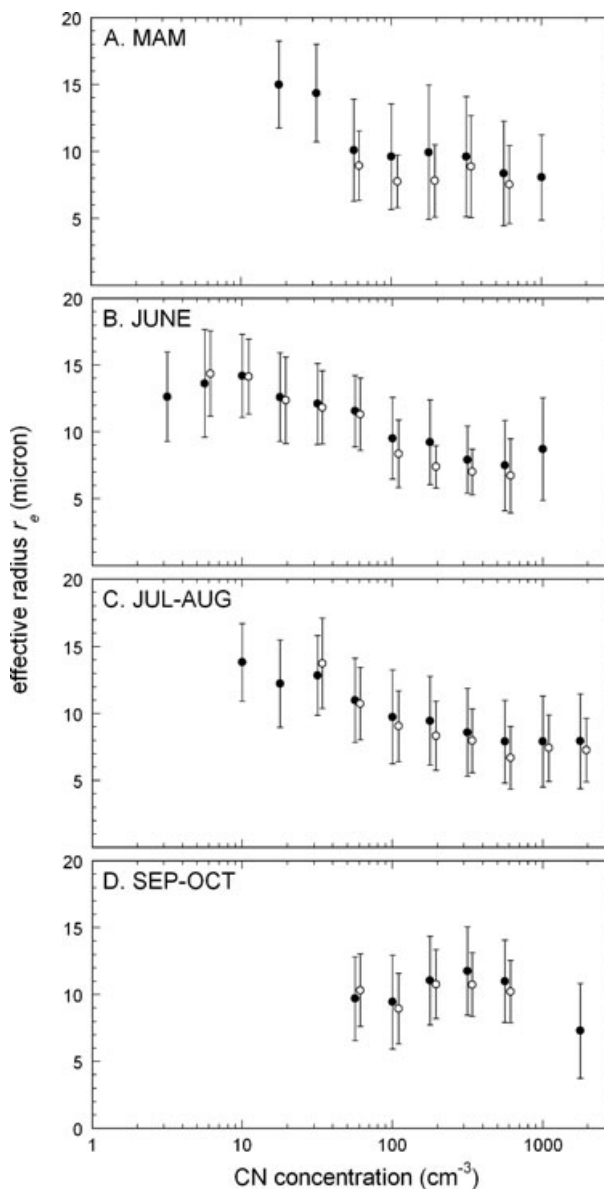
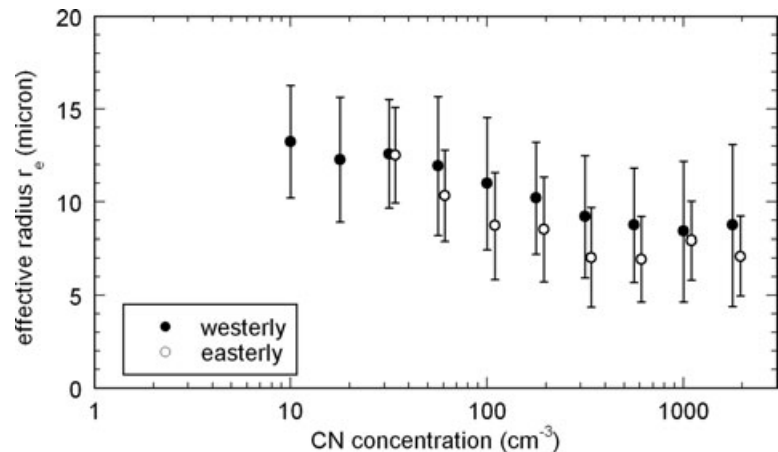


Fig. 4. Instantaneous AERI retrievals of  $r_e$  in low-level, optically thin ( $LWP < 50\text{ g m}^{-2}$ ) stratiform cloud over Barrow, AK, from years 1998–2003 and for the months (A) March, April and May (MAM), (B) June, (C), July and August and (D) September and October, as a function of the corresponding CN hourly averaged concentration. The  $r_e$  retrievals are averaged over CN concentration bins whose outer bounds are specified in the plot. Vertical error bars are  $\pm 1\sigma$ . Solid points are the averages over all applicable data; open circles are the averages over the  $LWP$  range  $5\text{--}15\text{ g m}^{-2}$ .

$r_e$  between the two wind regimes, but a relationship between  $r_e$  and CN consistent with an  $\text{IDE}_1$  appears with similar consistency for either regime. Thus, with the readily available NSA data, we find no obvious air mass correlation altering the result of Figs 4b and c.

Fig. 5. Instantaneous AERI retrievals of  $r_e$  in low-level, optically thin ( $LWP < 50 \text{ g m}^{-2}$ ) stratiform cloud cover over Barrow, Alaska, from years 1998–2003 and for the months of July and August, for two prevailing wind regimes, westerly ( $240\text{--}320^\circ$ ) and easterly ( $20\text{--}100^\circ$ ), as a function of corresponding hourly-averaged CN concentration, as in Fig. 4.



Research on the  $IDE_1$  at lower latitudes, where the energy budget is dominated by short-wave radiation, has emphasized the surface cooling by enhanced cloud reflection (Lohmann and Feichter, 2005; Penner et al., 2006) as opposed to the long-wave warming manifestation identified in the Arctic (Garrett and Zhao, 2006; Lubin and Vogelmann, 2006). To determine the total impact of the  $IDE_1$  on the Arctic surface energy budget and climate, we must examine its impact on both short-wave and long-wave radiation. The long-wave surface-warming component of the  $IDE_1$  for thin clouds has been estimated previously at  $+3.4 \text{ W m}^{-2}$  during spring (Lubin and Vogelmann, 2006). Here we observationally determine the short-wave component by using pyranometer measurements of downwelling short-wave surface irradiance simultaneously with the  $LWPs$  from AERI and MWR data.

Figure 6a shows the difference in  $T_c$  between clean-air and all other conditions in four bins encompassing the full range of observed cloud  $LWP$  in our data from March–May. A decrease in  $T_c$  signifies a reduction in short-wave flux at the surface, or a cooling effect. Decreased  $T_c$  in the presence of enhanced aerosols appears throughout the data set, and is more statistically significant at higher  $LWP$ . In contrast, the long-wave manifestation of the springtime  $IDE_1$  occurs only for clouds with  $LWP < \sim 50 \text{ g m}^{-2}$ , since optically thicker clouds radiate essentially as blackbodies (Turner, 2005; Garrett and Zhao, 2006; Lubin and Vogelmann, 2006). June (Fig. 6b) also shows an  $IDE_1$  manifestation, with differences between clean-air and all other conditions appearing as more statistically significant for smaller  $LWP$ . During July–August (Fig. 6c), statistically significant differences appear for all  $LWP$ . During September–October (Fig. 6d), there is no evidence of a short-wave manifestation of the  $IDE_1$ .

Figure 6 is highly suggestive of a short-wave manifestation of the  $IDE_1$ , since the  $t$ -tests for the various  $LWP$  bins indicate consistent statistical significance in the differences between the clean-air and the other cases across nearly the full range of  $LWP$ . However, Fig. 6 is not entirely conclusive because the  $LWP$  bins are quite wide. In Fig. 7, we perform a more stringent test

for persistence of the  $IDE_1$  at constant  $LWP$  by calculating the average  $T_c$ , in both clean-air and all other conditions, in averages of 30 data points about each  $LWP$  value shown. The result shown in Fig. 7 is conceptually consistent with the  $IDE_1$  if cloud  $T_c$  differences are found between clean-air and all other conditions for fixed  $LWP$ . In both spring and summer, the  $T_c$  difference shows a local minimum around  $LWP = 30 \text{ g m}^{-2}$ , which is consistent with radiative transfer simulations that indicate the  $IDE_1$  should be harder to detect in surface Arctic data in this moderate  $LWP$  range (Lubin and Vogelmann, 2007). For all other  $LWP$  during spring, the  $T_c$  difference between clouds under clean-air versus all other conditions is positive and above the instrument detection uncertainty, which signifies a reduction in surface insolation under cloud when  $CN > 50 \text{ cm}^{-3}$ . A similar result is found for optically thinner clouds during summer. This result is less pronounced in June, where  $T_c$  differences hover just above or about the detection level for  $30 < LWP < 80 \text{ g m}^{-2}$  and then drop below the detection uncertainty for larger  $LWP$ . However, during July and August, pronounced  $T_c$  differences are found for  $LWP < 100 \text{ g m}^{-2}$ , while for larger  $LWP$  values the differences are consistently less and, in many cases, drop below the detection uncertainty. During autumn the  $T_c$  differences are positive but almost entirely below the detection uncertainty, as may be expected from Figs. 4d and 6d.

#### 4. Discussion

These results show observational evidence of a short-wave  $IDE_1$  in the Arctic for all observed cloud  $LWP$  during spring, and during late summer for  $LWP < 100 \text{ g m}^{-2}$ . The summertime result is surprising, given the conventional wisdom that the  $IDE_1$  is an anthropogenic springtime phenomena in the Arctic (Garrett and Zhao, 2006; Lubin and Vogelmann, 2006) and given that aerosol mass concentrations during the Barrow summer are typically —four to five times smaller than in spring (Quinn et al., 2002). However, this summertime result may be valid based the following two considerations. First, Dong and Mace (2003)

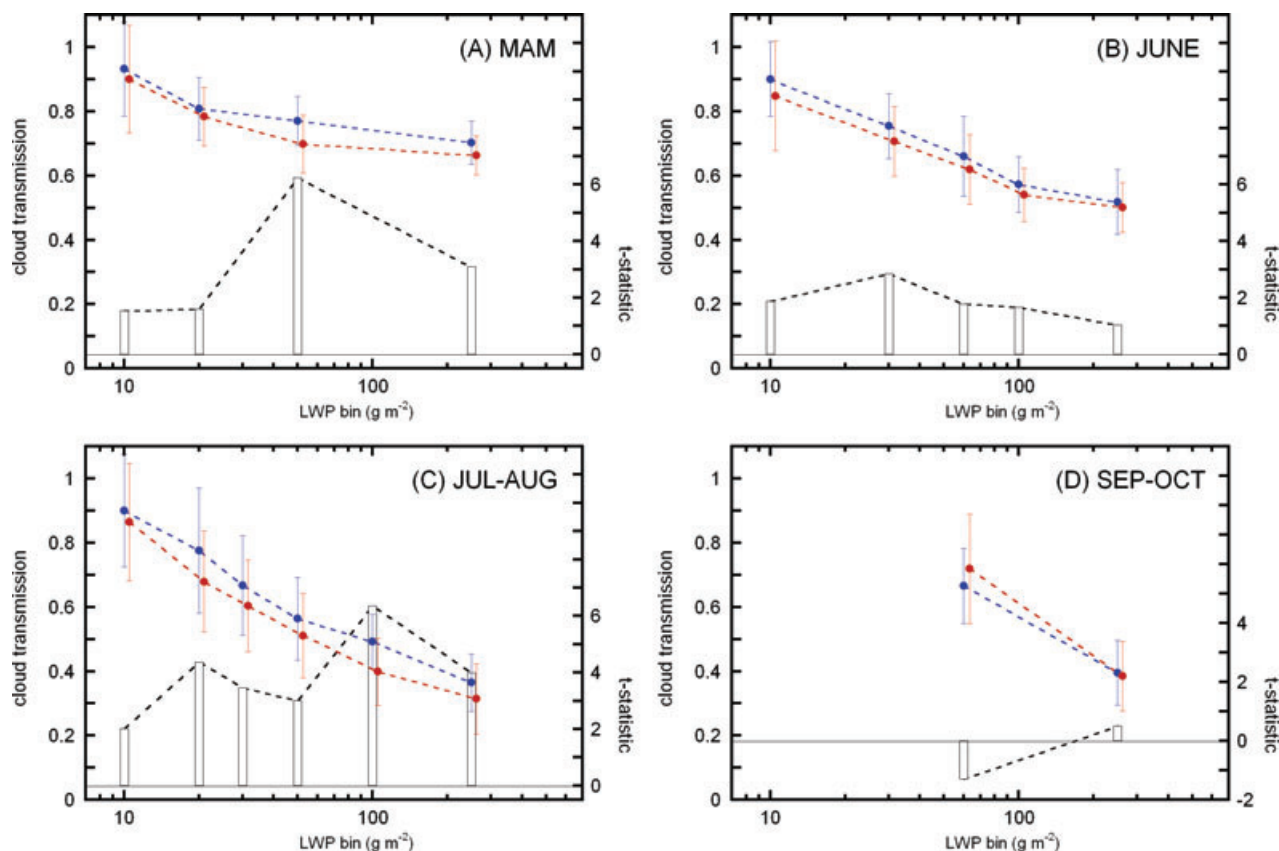


Fig. 6. Cloud transmissions  $T_c$  under clean-air (blue;  $CN < 50 \text{ cm}^{-3}$ ) and all other conditions (red;  $CN > 50 \text{ cm}^{-3}$ ), averaged over all data from (A) March–May (1998–2003), (B) June, (C) July–August and (D) September–October. LWP bin ranges are determined by the availability of enough data points ( $n \geq 30$ ) for meaningful statistics: (A) 0–10, 10–20, 20–50 and 50–250  $\text{g m}^{-2}$ ; (B) 0–10, 10–30, 30–60, 60–100  $\text{g m}^{-2}$  and 100–250  $\text{g m}^{-2}$ ; (C) 0–10, 10–20, 20–30, 30–50, 50–100, 100–250  $\text{g m}^{-2}$  and (D) 0–60, 60–250  $\text{g m}^{-2}$ . Vertical error bars are  $\pm 1\sigma$ , and the statistical significance of the average  $T_c$  differences between clean-air and all other conditions in each bin are indicated by Student's  $t$ -test. All  $T_c$  values used in this averaging have been normalized to a solar zenith angle of  $60^\circ$  to remove a small but non-negligible solar zenith angle dependence, as described in the text.

used NSA cloud radar, pyranometer and microwave radiometer data to retrieve  $r_e$  for NSA clouds, where LWP is determined by the microwave radiometer. Their retrievals show  $r_e$  in the range  $7\text{--}10 \mu\text{m}$  during spring and  $10\text{--}15 \mu\text{m}$  during summer. The latter are consistent with our summertime values for clean air in Fig. 4 ( $CN < 50 \text{ cm}^{-3}$ ), but are larger than our retrieved  $r_e$  under  $CN > 50 \text{ cm}^{-3}$ . We infer from this that summertime aerosol do not affect optically thicker clouds sufficiently to change their climatologically averaged  $r_e$  and, hence, the surface radiation balance under them. It remains to be explained why optically thicker clouds appear to be unaffected. Second, a recent study of spring and summertime Arctic aerosol, using well-tested global aerosol and chemical transport models (Korhonen et al., 2008), shows that updraft velocities frequently observed in summertime Arctic boundary layer clouds, of order  $15 \text{ cm s}^{-1}$  (Lawson et al., 2001), can induce cloud nucleation with the locally produced sulfuric acid particles in the Aitken mode. Arctic Ocean observations of Aitken-mode particles and CCN support this hypothesis (Lohmann and Leck, 2005). Thus, there a plausible

mechanism for the summertime IDE<sub>1</sub> that we observe primarily in optically thinner clouds. Further, this effect may also represent a negative feedback mechanism (net cooling effect) as the Arctic climate warms, since Gabric et al. (2005) have shown the potential for significant increases in dimethylsulfide (DMS) flux via retreating sea ice cover and warming sea surface temperatures in the Arctic. We note that this summertime effect is at slight variance from the usual definition of IDE<sub>1</sub> given by Twomey (1977), who originally described it with reference to pollution aerosol. Here, if the above-described mechanism for the summer IDE<sub>1</sub> applies, it manifests from an additional flux of natural particles in a warming climate. Finally, we note that our sample cloud distribution is heavily weighted toward geometrically thin clouds very close to the surface (Fig. 3), and our results concerning to the IDE<sub>1</sub> in all seasons pertain only to these clouds and not to cloud cover at higher altitudes. However, Fig. 3, based on 6 yr of ARM data, indicates that thin low level clouds are very frequent in the Arctic in a climatological sense.

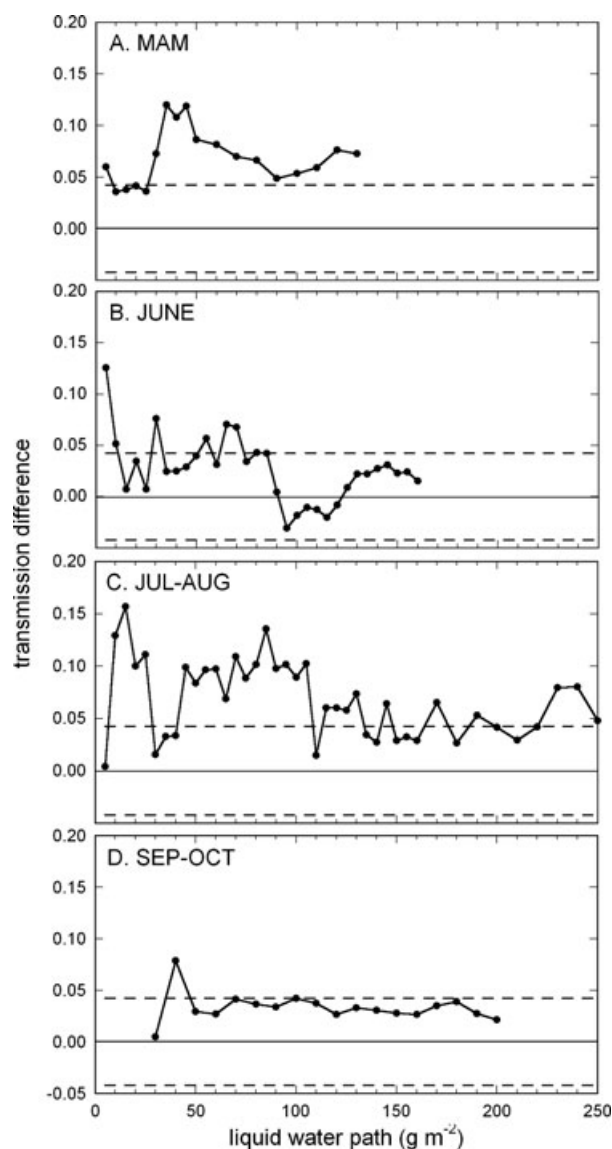


Fig. 7. Cloud transmission differences (clean-air minus all other conditions) evaluated by a moving average constant LWP window for the months (A) March, April and May (MAM), (B) June, (C), July and August and (D) September and October. Each point depicts the difference in mean  $T_c$  for the nearest 30 data points having the LWP average shown. Dashed lines depict the uncertainty in the  $T_c$  difference, which are computed based on the insolation levels at the particular times of year and the  $\pm 5 \text{ W m}^{-2}$  uncertainty in pyranometer-measured flux that is estimated for the ARM program's data collection and reduction methods (Michalsky et al., 1999). All  $T_c$  values used in this averaging have been normalized to a solar zenith angle of  $60^\circ$  to remove a small but non-negligible solar zenith angle dependence, as described in the text.

Although explanations for our summertime evidence of an IDE<sub>1</sub> can be proposed, we caution that this result remains counterintuitive based on our understanding of the seasonal variation in Arctic aerosol microphysics and chemistry. In demonstrating

this result with the available ground-based data, we cannot use CN to determine a threshold for detecting CCN during summer. Instead, we can only use CN to isolate the clean-air cases from all other cases, which may or may not involve CCN. We do not find an obvious airmass correlation artefact in the summertime result when we examine the available surface meteorological data (Fig. 5); however a more extensive backtrajectory analysis of this multiyear NSA data set could be beneficial. At this point, our results suggest that theoretical and field investigations should examine potential manifestations of an IDE<sub>1</sub> during the Arctic summer. However, we do not suggest that the magnitude of the summertime IDE<sub>1</sub> that appears in this study should be taken as a benchmark observation until the result is more fully explained and/or all conceivable artefacts are conclusively eliminated.

In summary, our observations show a short-wave IDE<sub>1</sub> during spring, evidence of a short-wave IDE<sub>1</sub> in cloud with smaller LWP during summer, and a barely perceptible IDE<sub>1</sub> during autumn. The autumn result is not surprising given that: (1) during refreeze of the sea ice the above-described DMS particle flux would decrease; (2) the overall aerosol burden is substantially reduced as compared with spring, after several months of removal by precipitation and influence of air masses from lower latitudes and (3) mixed-phase clouds are highly prevalent during autumn in the high Arctic (e.g. Verlinde et al., 2007), which reduces the sample sizes available for the liquid water cloud analysis described herein. Spring holds the greatest interest, due to the enormous seasonal increase in insolation that modulates the surface energy budget just prior to the melt season that often begins in late May (Perovich et al., 2002; also see Fig. 1).

We now estimate the total (short-wave plus long-wave) IDE<sub>1</sub> impact on the surface radiation budget by combining our short-wave result in Fig. 6a with our previous estimate of the long-wave warming from the IDE<sub>1</sub> (Lubin and Vogelmann, 2006) using our radiative transfer model and the observed distribution of springtime Arctic cloud LWP at Barrow (Lubin and Vogelmann, 2007). With these data, we calculate the time evolution of the total IDE<sub>1</sub> impact on the downwelling surface radiation budget at Barrow, Alaska and at latitudes farther north (Fig. 8), noting that shipboard transects of cloud observations as far as the pole (Lubin and Simpson, 1997) suggest that stratiform cloud opacities observed from Barrow's latitude are representative of more northerly latitudes as well. At most locations during March, the IDE<sub>1</sub> yields a surface warming. During April, the IDE<sub>1</sub> transitions to a slight cooling, whose magnitude varies with latitude, as the additional insolation attenuated by clouds with smaller  $r_e$  begins to exceed the enhanced cloud emission caused by the smaller  $r_e$ . Throughout May, surface cooling by the IDE<sub>1</sub> becomes more spatially uniform and its absolute magnitude increases from  $-7$  to  $-11 \text{ W m}^{-2}$ . Because both low-level cloud cover and anthropogenic aerosol are pervasive phenomena throughout the Arctic winter and spring, climate model simulations should account for this aerosol-cloud interaction which yields a climatologically significant surface warming



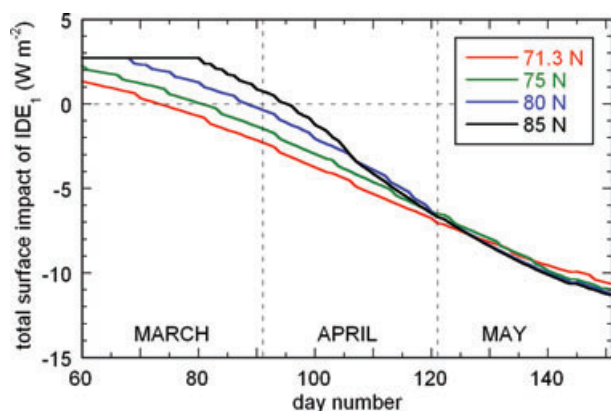


Fig. 8. Time evolution of the springtime total IDE<sub>1</sub> impact on the downwelling surface energy budget, based on the results of Fig. 2, and the climatological LWP distributions for low-level Arctic clouds (Lubin and Vogelmann, 2007) and a climatological cloud amount of 80% (Lubin and Vogelmann, 2006). Results are derived using the radiative transfer model for Barrow's latitude (71.3°N) and three higher latitudes. The IDE<sub>1</sub> impact is expressed as the change in diurnally averaged total downwelling surface flux (short-wave plus long-wave) due to an aerosol-induced reduction in  $r_e$ . The vertical dashed lines identify April 1 and May 1. Positive (negative) values indicate total surface warming (cooling).

during winter and early spring, followed by surface cooling just prior to the onset of sea ice melt.

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