

## Appendix S1. Derivation of Eq. 5

We consider the case in which  $f_{min} = 0$  and  $f^{max} \rightarrow \infty$ . Substituting Eqs. 1 and 4 into Eq. 5,

$$\lambda_i(\alpha) = \int_0^\infty df \exp\left(-\frac{1}{2}\left(\frac{\ln(f/\phi_i)}{s}\right)^2\right) \frac{c}{z(\alpha)} f^{-\alpha}.$$

Denoting  $\ln(f/\phi_i)/s$  as  $y$ , since  $f = \phi_i e^{sy}$  and  $df = s\phi_i e^{sy} dy$ ,

$$\begin{aligned} \lambda_i(\alpha) &= cs \frac{\phi_i^{1-\alpha}}{z(\alpha)} \int_{-\infty}^{\infty} dy \exp\left(-\frac{y^2}{2} + (1-\alpha)sy\right) \\ &= cs \frac{\phi_i^{1-\alpha}}{z(\alpha)} \int_{-\infty}^{\infty} dy \exp\left(-\frac{1}{2}(y - (1-\alpha)s)^2 + \frac{(1-\alpha)^2 s^2}{2}\right). \end{aligned}$$

Finally, we have

$$\lambda_i(\alpha) = cs\sqrt{2\pi} \frac{\phi_i^{1-\alpha}}{z(\alpha)} \exp\left(\frac{(1-\alpha)^2 s^2}{2}\right).$$

Note that, in reality, for  $f^{max} \rightarrow \infty$  the normalization constant  $z(\alpha)$  takes a near-zero or an extremely large value depending on the value of  $\alpha$ ; in this paper, we consider the above calculation as an approximation of the situation in which the image resolution (i.e., the maximum spatial frequency) is sufficiently large.