

Non-linear Transmission Line (NTL) Study of Electromagnetic (EM) Modes in Capacitive Discharges

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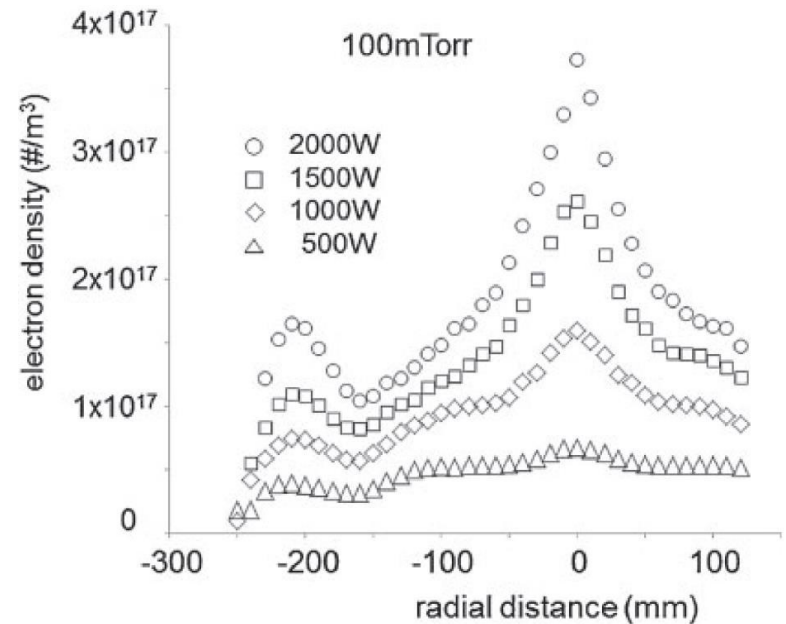
On-line LTP Seminar, 1/25/2022

**Ref: Jian-Kai Liu, E. Kawamura, M.A. Lieberman, Yong-Nian Wang,
Phys. Plasmas 29, 013058 (2022)**

Introduction

Larger area, higher frequency capacitively coupled plasma (CCP) reactors are being used in Si wafer and flat panel processing to **maximize production**

- When $\lambda \gg R$ is no longer satisfied, finite wavelength or **standing wave** effects can lead to center-high plasma non-uniformity, especially near a **spatial resonance**
- **Harmonics** excited in the plasma due to the **non-linear sheaths** also enhance non-uniformity due to their **shorter wavelengths**
- These **non-linearly excited harmonics** can be enhanced by the **series resonance**

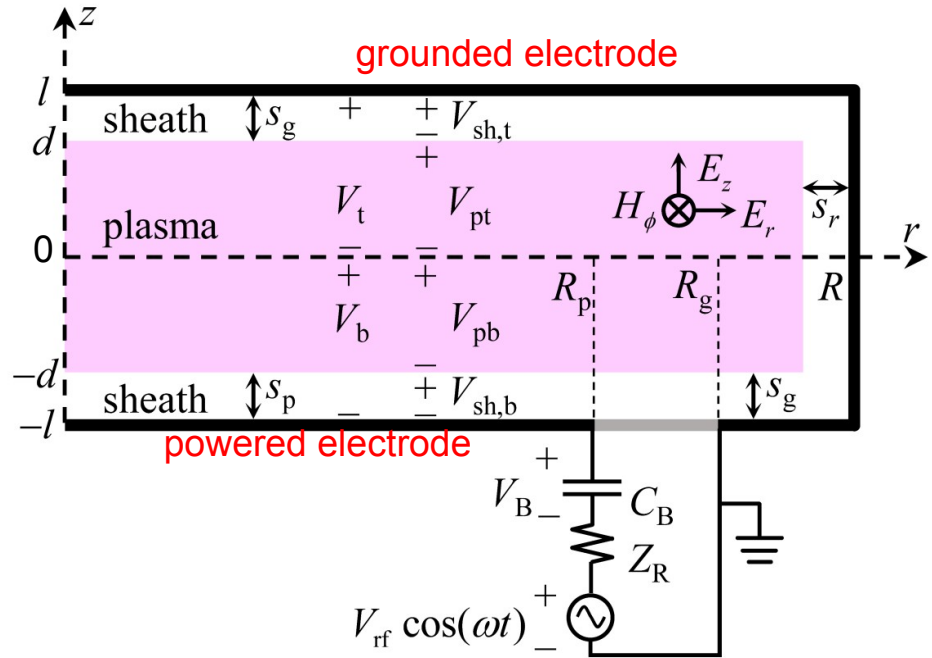


Mid-gap n_e profiles for a 100 mTorr Ar discharge with radius $R=26$ cm and gap $L=2.5$ cm, operated at $f=60$ MHz.

Sawada et al, Japan J. App. Phys. 53, 03DB01 (2014)

Asymmetric Cylindrical Reactor Geometry

- Capacitive EM fields are in the transverse magnetic (TM) mode: (E_r, E_z, H_ϕ)
- At **low pressure** p_g , relative dielectric constant of plasma: $\kappa_p = 1 - \omega_p^2 / \omega^2 \approx -\omega_p^2 / \omega^2$, where $\omega_p = [ne^2 / (m\epsilon_0)]^{1/2}$ and $\omega = 2\pi f$ are the plasma and applied radian frequencies
- Matching plasma inductance L_p with sheaths' capacitance C_s gives the series resonance: $\omega_{\text{res}} = 1 / (L_p C_s)^{1/2} = \omega_p [(s_g + s_p) / 2d]^{1/2}$



- $R = 25 \text{ cm}$ and $2l = 5 \text{ cm}$
- $R_p = 15 \text{ cm}$ and $R_g = 16 \text{ cm}$
- RF voltage source $V_{\text{rf}} \cos(\omega t)$
- Source resistance $Z_R = 0.5 \Omega$
- Blocking capacitor C_B

Symmetric and Antisymmetric Modes

- Two radially propagating surface wave modes in the plasma:

Symmetric: $E_{zs}(z) = E_{zs}(-z)$

Antisymmetric: $E_{za}(z) = -E_{za}(-z)$

- At low pressure, the transverse wavenumber $k_t \approx j \omega_p / c \equiv j \alpha \Rightarrow$
 $E_{zs}(z) \sim \cos(k_t z) \approx \cosh(\alpha z)$ and
 $E_{za}(z) \sim \sin(k_t z) \approx \sinh(\alpha z)$

- Wavelengths of the modes:

$$\lambda_s \approx \lambda_{\text{vac}} [(s_p + s_g)/2]^{1/2}, \lambda_{\text{vac}} = c/f$$

$$\lambda_a \approx \lambda_{\text{vac}} (\omega_p/c) [2d s_g s_p / (s_p + s_g)]^{1/2}$$

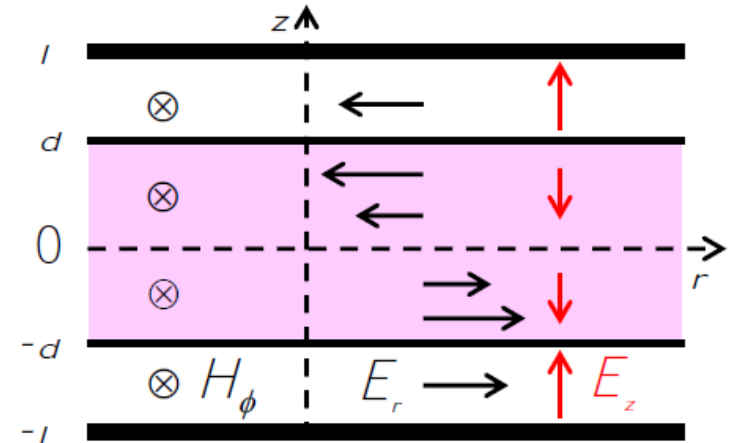
- First spatial resonances occur at:

$$k_s R = (2\pi/\lambda_s) R = \chi_{01}, \quad k_a R = (2\pi/\lambda_a) R = \chi_{11}$$

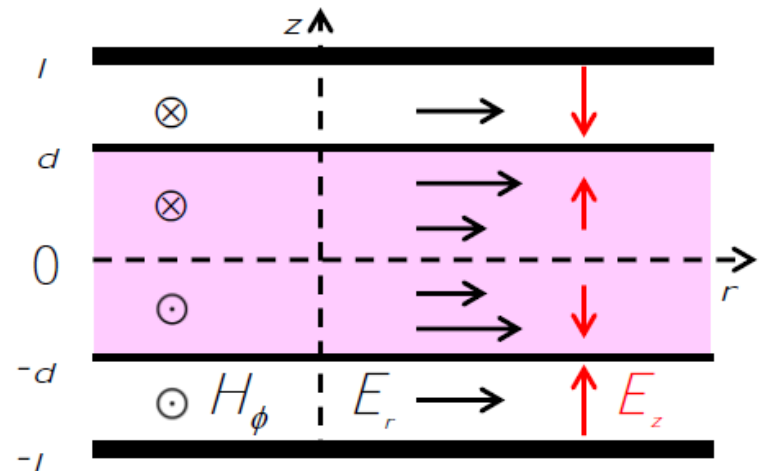
with $\chi_{01} \approx 2.405$, $\chi_{11} \approx 3.832$ the 1st

zeroes of $J_0(k_s r)$ and $J_1(k_a r)$

Symmetric Mode:



Antisymmetric Mode:

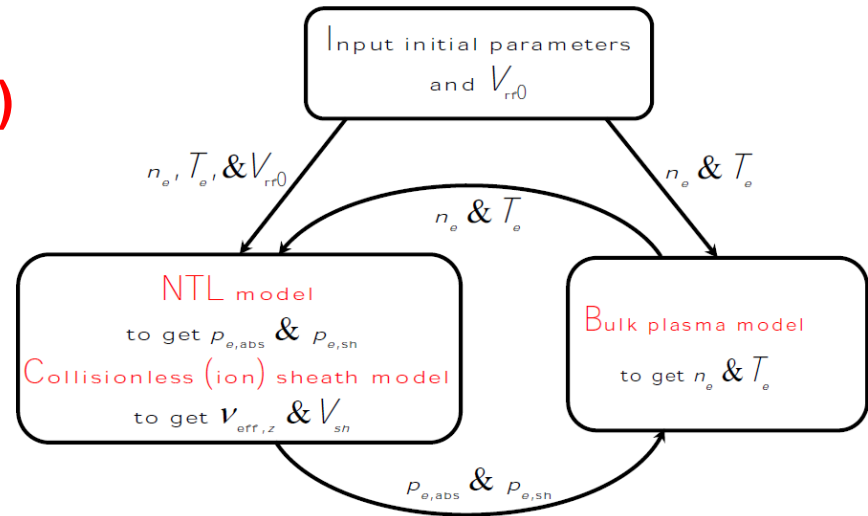


Low pressure Capacitive Discharge Model

- Analytical axial field solve:

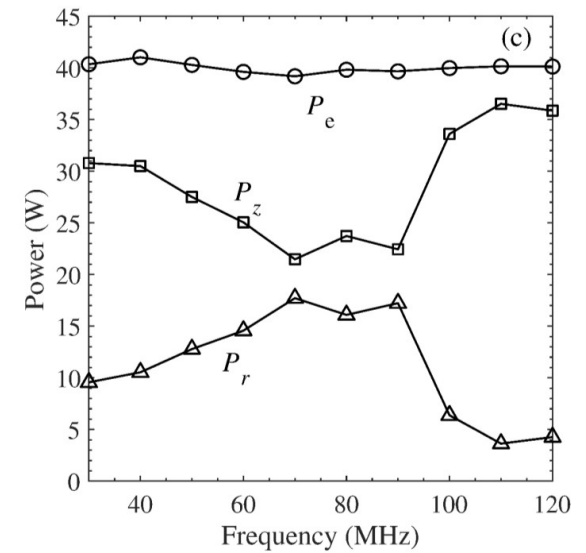
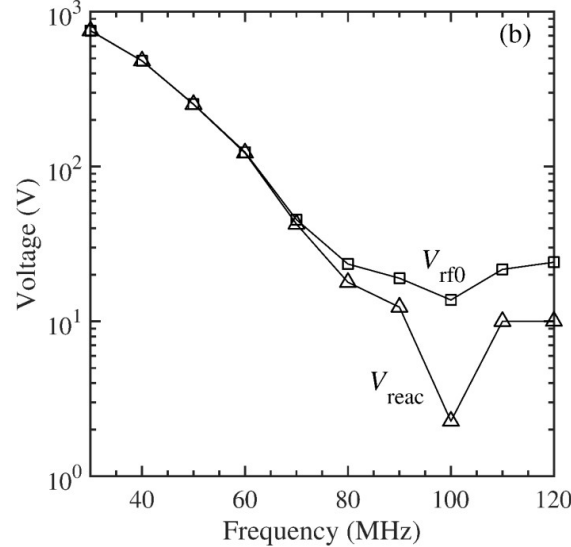
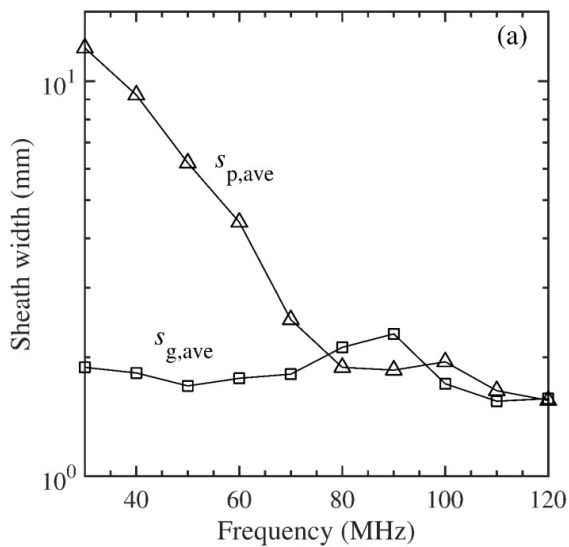
$$E_z(z,r,t) = (\cosh(\alpha z)/\cosh(\alpha d))E_{zs}(r,t) + (\sinh(\alpha z)/\sinh(\alpha d))E_{za}(r,t)$$

- 1D radial NTL model: Maxwell's equations and electron momentum balance equation => set of non-linear **PDE's in (r,t)** to solve numerically



- Analytical collisionless (ion) Child Law sheath model: Finds sheath voltages, sheath widths and an anisotropic effective collision frequency: $v_{eff,z} = v_m + v_{stoc} + v_{ohm,sh}$ ($v_{eff,r} = v_m$)
- 2D quasi-neutral, ambipolar bulk fluid model obtains $\langle n_e \rangle, \langle T_e \rangle$

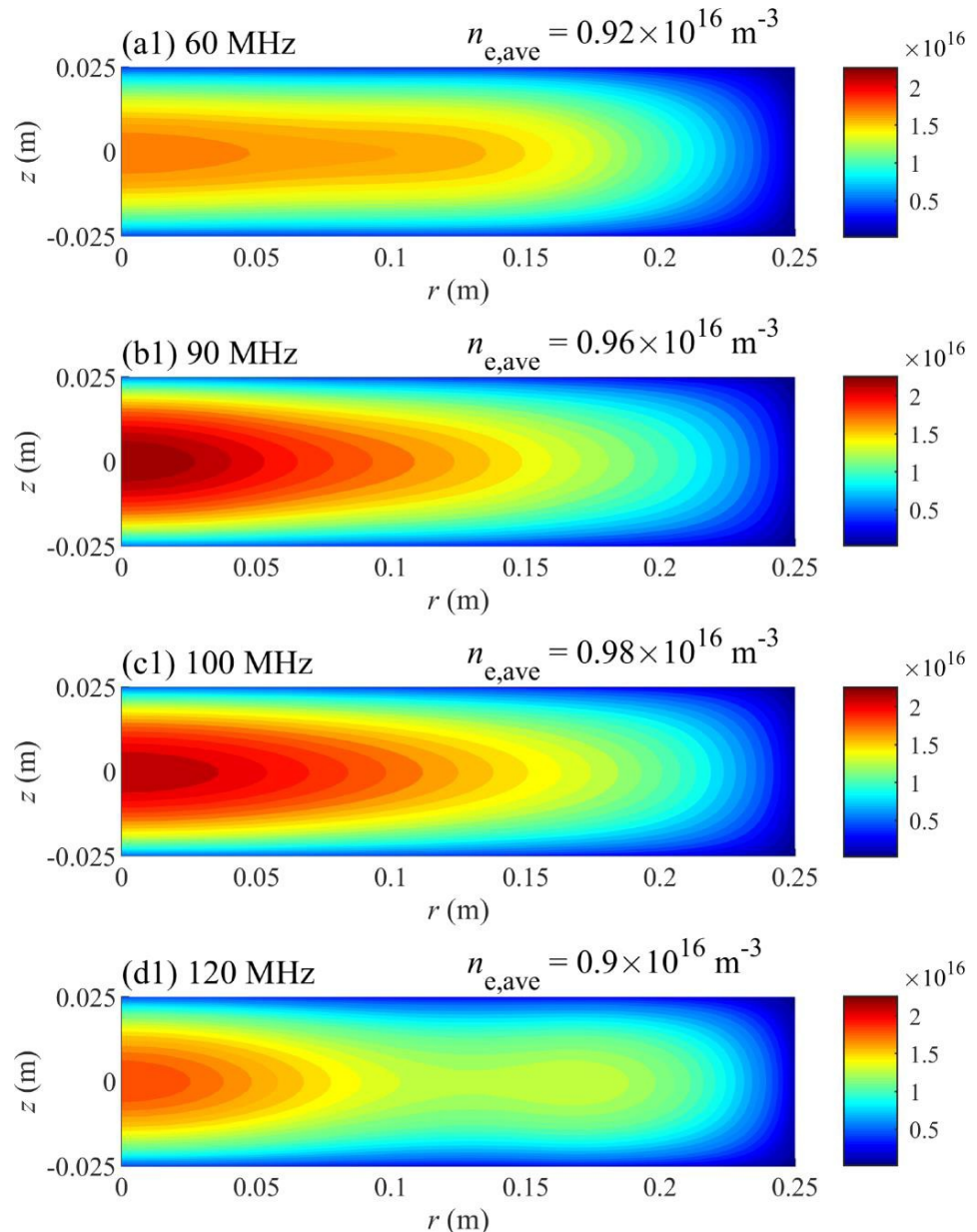
20 mTorr Ar Results for $P_e=40$ W at $f=30$ -120 MHz



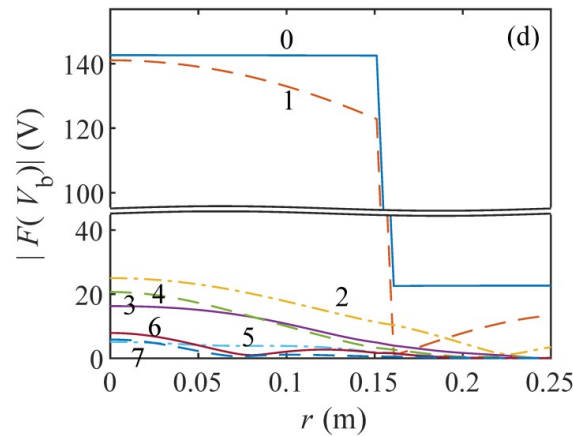
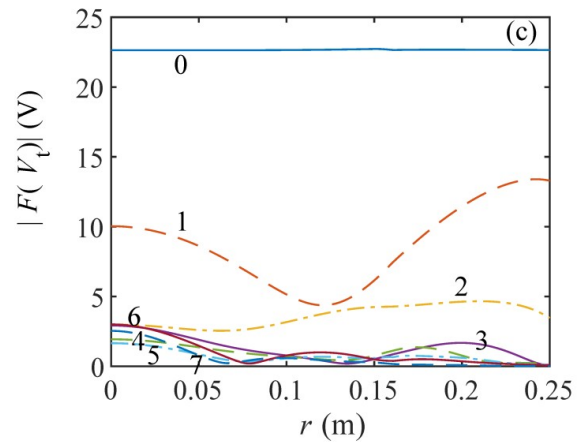
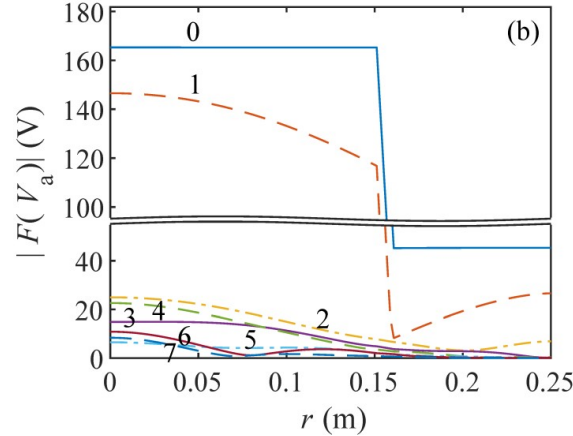
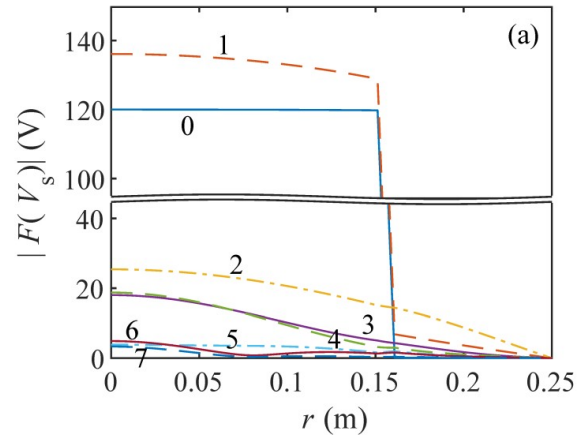
- P_r dominated by **anti-sym** mode; P_z dominated by **sym** mode
- Source voltage and s_p tends to decline with increasing f at fixed P_e
- At **lowest f** , **sym mode dominates** so that $s_p > s_g$ and $P_r \ll P_z$
- As $f \rightarrow f_{a1} \sim 85$ MHz, $s_g > s_p$ and $P_r \rightarrow P_z$, and as $f \rightarrow f_{s1} \sim 100$ MHz the source voltage reaches a minimum and $P_z \rightarrow P_e$ and $P_r \rightarrow 0$

Plasma density at 60, 90, 100 and 120 MHz

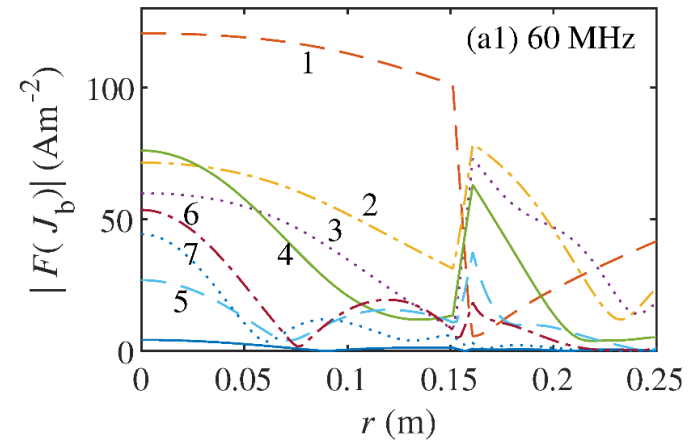
- 60 MHz, $f < f_{a1}$, f_{s1} : center peak due to finite λ effect and non-linear harmonics
- 90 MHz $\sim f_{a1}$, 100 MHz $\sim f_{s1}$: sharper peaks due to shorter λ at higher f
- 120 MHz $> f_{a1}$, f_{s1} : second peak since shorter λ allow secondary maxima of surface waves to appear within reactor



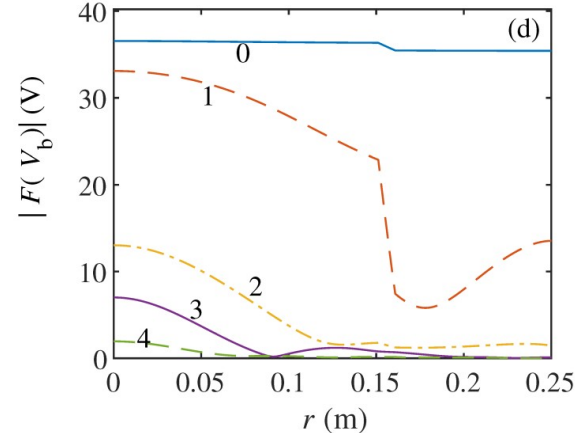
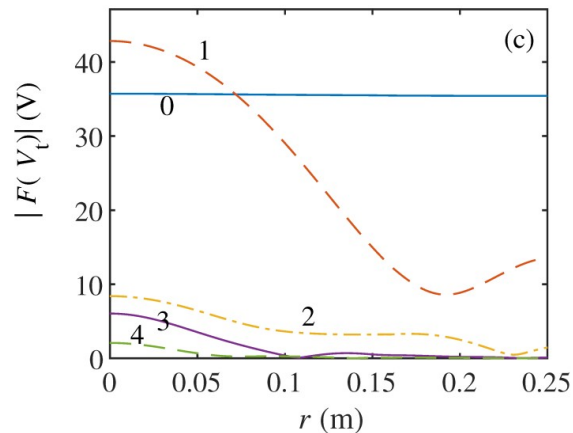
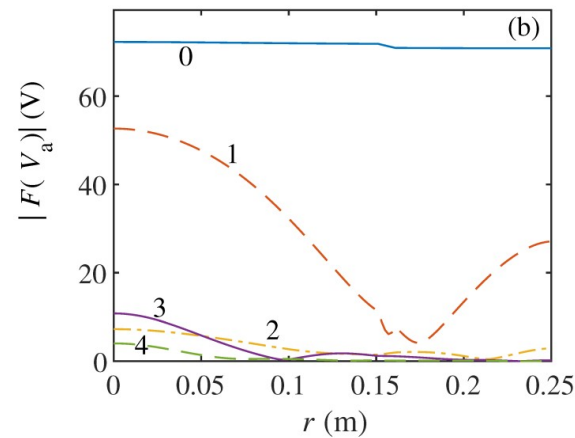
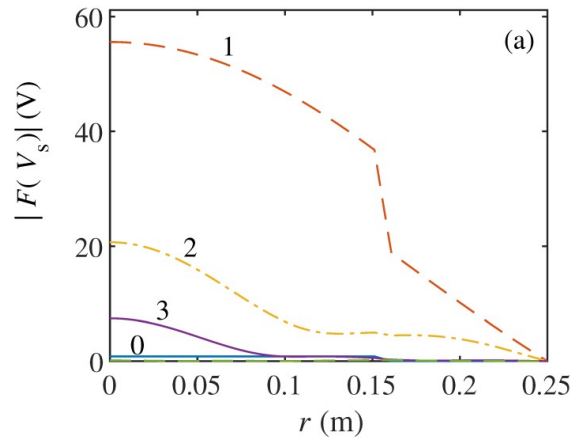
Fourier Transform of Voltages at 60 MHz



- 60 MHz, $f < f_{a1}, f_{s1}$:
center-high non-uniformity is enhanced by non-linear harmonics
- $f_{\text{res}} \sim 400$ MHz is between 6th and 7th harmonic

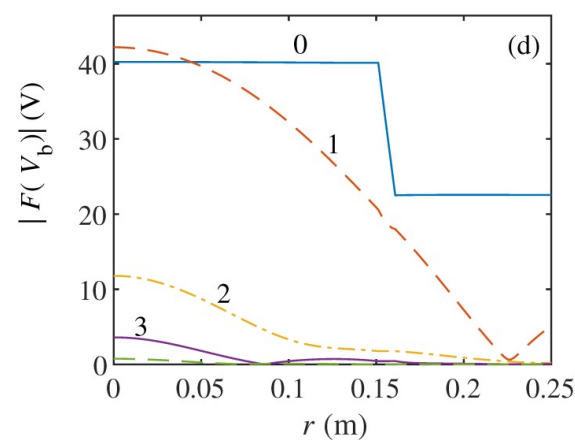
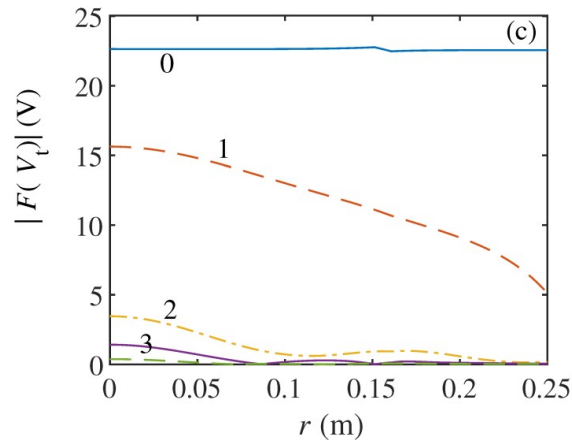
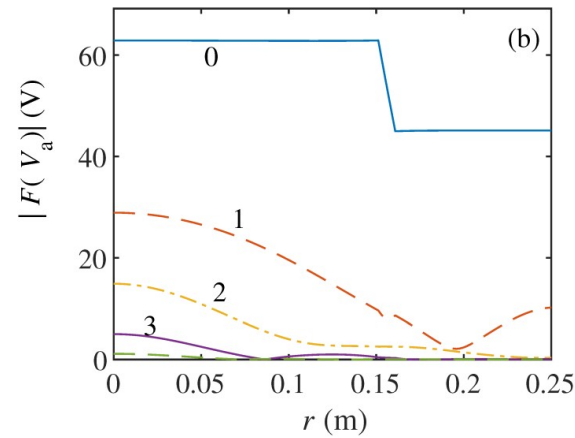
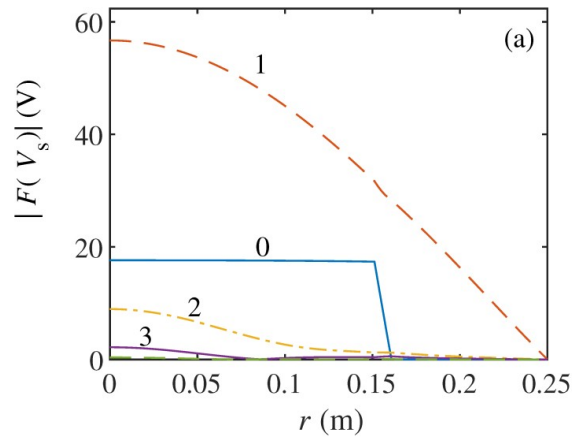


Fourier Transform of Voltages at 90 MHz



- 90 MHz, $f \sim f_{a1}$: Due to shorter wavelength at higher f , sharper center-high non-uniformity
- Fewer harmonics excited due to decreasing V_{rf} with f , and more power going to lower harmonics near their spatial resonances

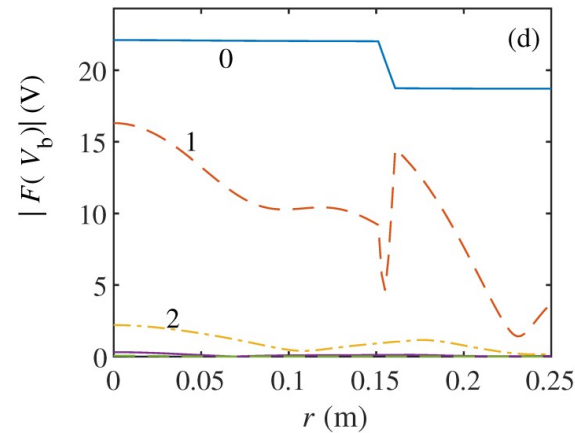
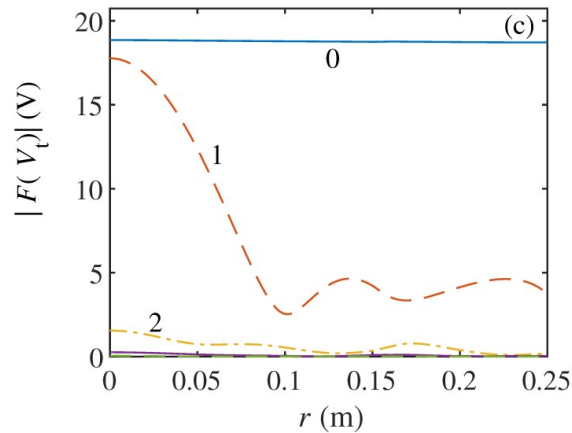
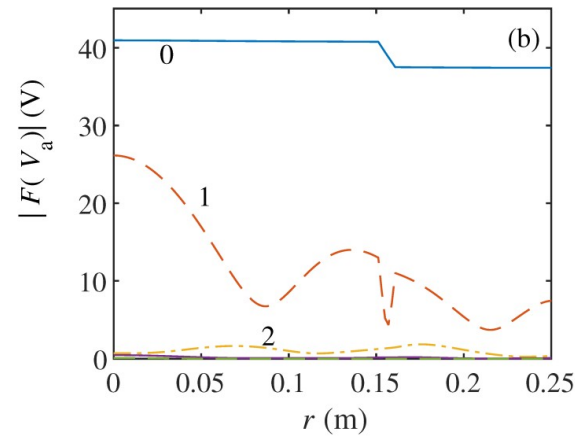
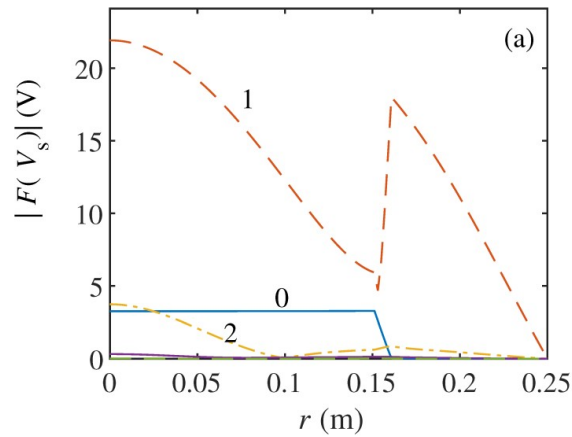
Fourier Transform of Voltages at 100 MHz



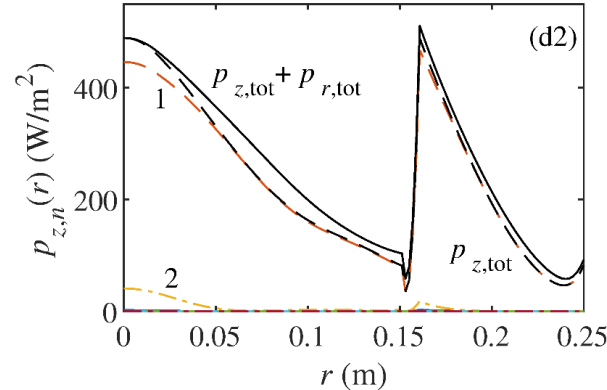
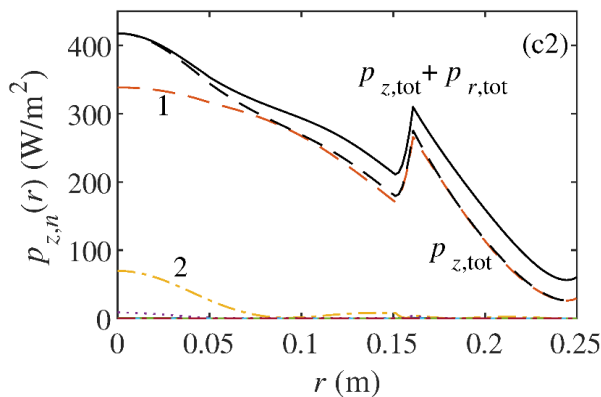
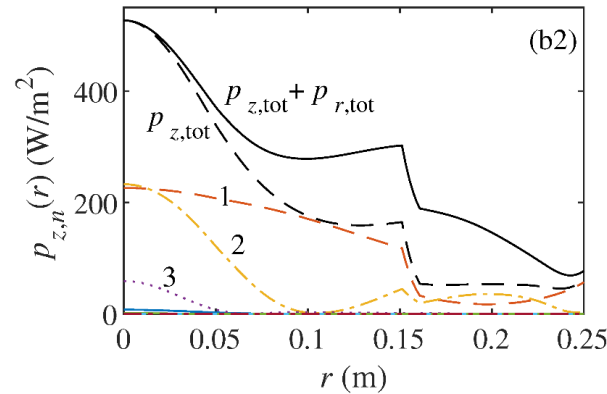
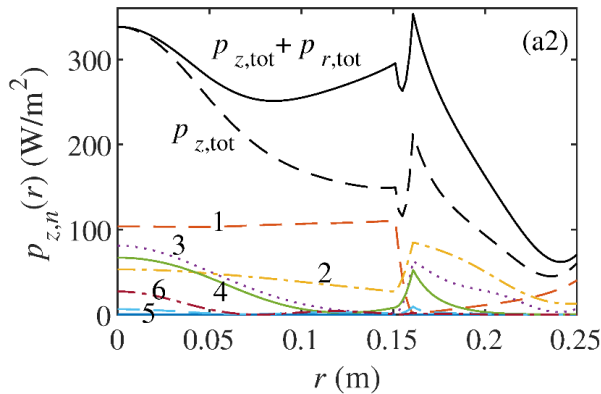
- 100 MHz, $f \sim f_{s1}$, but still near f_{a1}
- Source voltage has minimum at 100 MHz, resulting in few excited harmonics

Fourier Transform of Voltages at 120 MHz

- 120 MHz, $f > f_{s1}, f_{a1}$
- Observe secondary maxima in 1st harmonic due to shorter wavelengths at higher f



Electron Power Density at 60, 90, 100 and 120 MHz



- At 60 MHz significant center-high peak due to numerous higher harmonics
- As f increases, less higher harmonics are excited as V_{rf} decreases with f , and more P_e goes to lower harmonics near their spatial resonances
- At 120 MHz, see a significant secondary peak due to shorter wavelength

Conclusions

- The 1st antisymmetric and symmetric spatial resonances occur at $f_{a1} \sim 85$ MHz and $f_{s1} \sim 100$ MHz for a 20 mTorr Ar discharge with $2l=5$ cm, $R=25$ cm and fixed $P_e=40$ W
- Due to the different field symmetries, P_r is dominated by the **antisymmetric** mode while P_z is dominated by the **symmetric** mode
- At **lower f** , the **symmetric mode dominates** so that the sheath widths and voltages at the powered electrode are larger than at the grounded electrode, and $P_r \ll P_z$
- As $f \rightarrow f_{a1}$, $s_g > s_p$ and $P_r \rightarrow P_z$, and as $f \rightarrow f_{s1}$, the source voltage reaches a minimum and $P_z \rightarrow P_e$ and $P_r \rightarrow 0$
- Density shows a **sharp center peak** due to the standing wave effect, **nonlinearly excited harmonics** and spatial or series **resonances**
- For $f > f_{s1}$, the shortened wavelengths allow **secondary maxima of the modes** to occur within the reactor, causing **secondary n_e peaks**