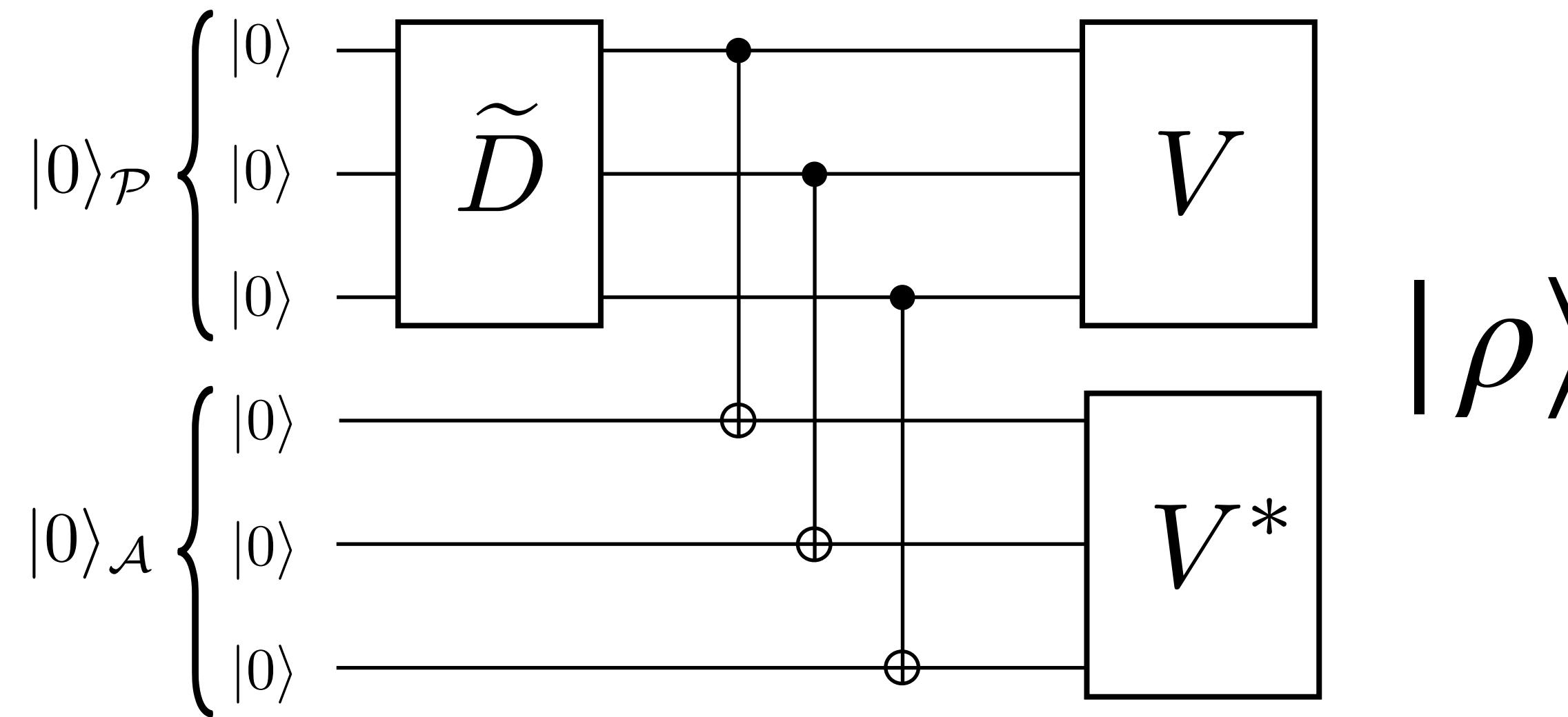


VARIATIONAL QUANTUM ALGORITHM FOR MARKOVIAN OPEN QUANTUM SYSTEMS



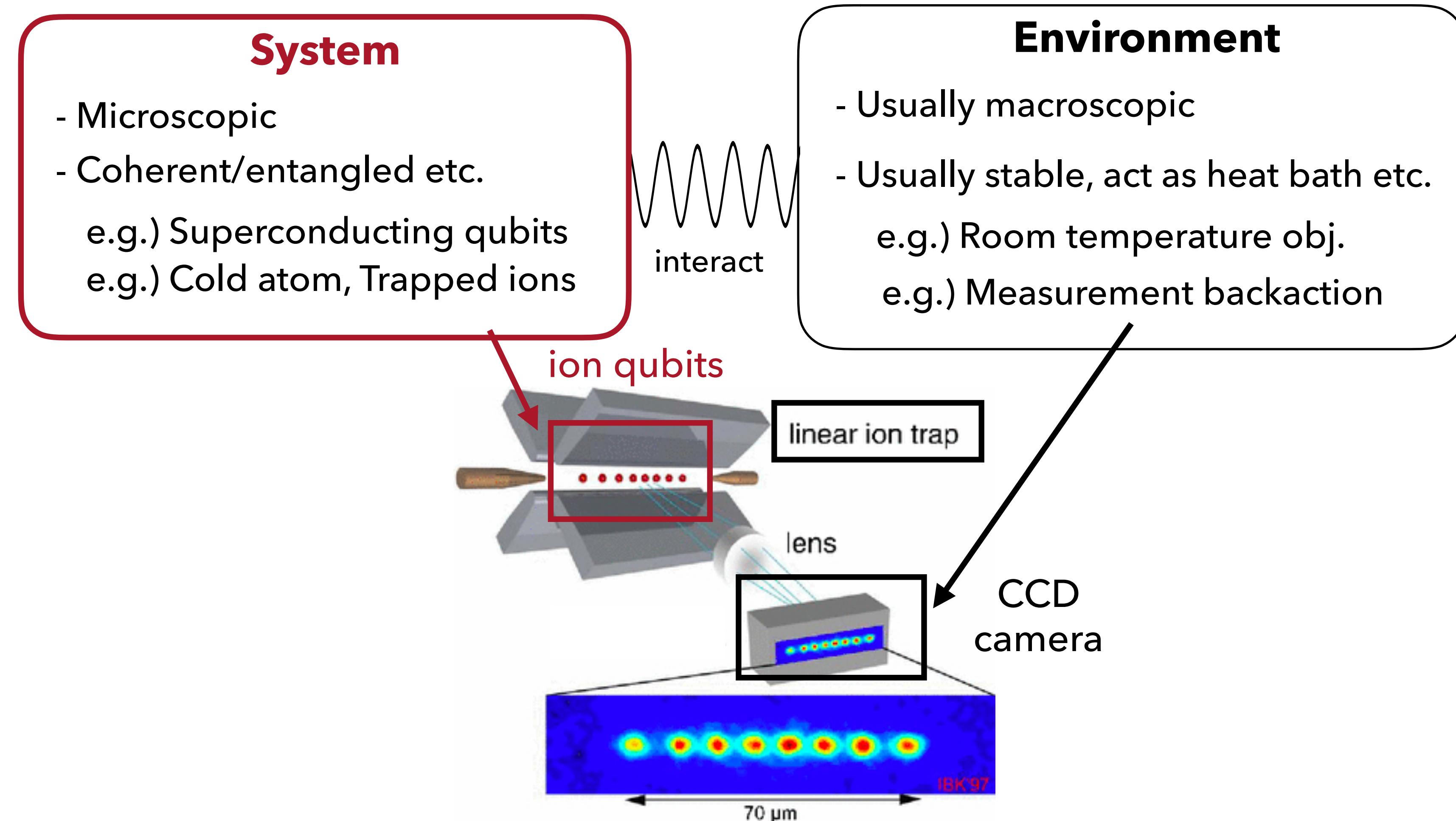
N. Yoshioka (U. Tokyo), Y. O. Nakagawa (QunaSys)



K. Mitarai, and K. Fujii (Osaka U.)

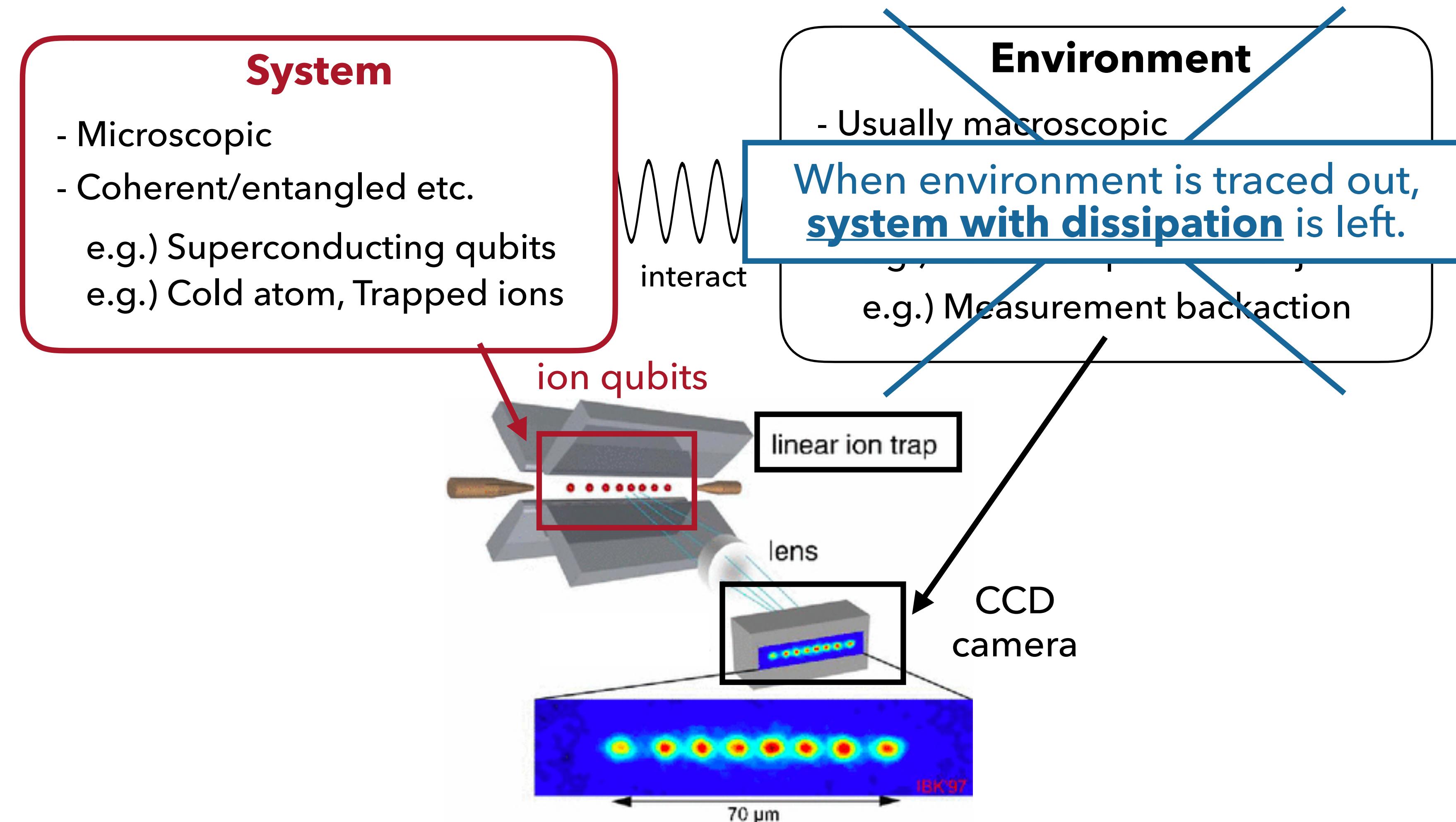
arXiv: 1908.09836

Introduction: Open quantum system



Taken from Garcia-Ripoll et al., J. Phys B ('05)

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Time evolution and eigenstates in isolated/open systems



System	Time evolution equation	Eigenstates	Remarks
Isolated	Pure $i\hbar \frac{d \psi\rangle}{dt} = \hat{H} \psi\rangle$ Hamiltonian	$\hat{H} \psi\rangle = E \psi\rangle$ <u>Ground/excited states</u>	<ul style="list-style-type: none"> - Application of NISQ to <ul style="list-style-type: none"> - Quantum chemistry - Finance - Numerous investigations done
Open	Mixed $\frac{d\hat{\rho}(t)}{dt} = \mathcal{L}\hat{\rho}(t)$ Liouvillian	$\mathcal{L}\hat{\rho} = 0$ <u>Steady state</u> Target today $\mathcal{L}\hat{\rho} = \lambda\hat{\rho}$ <u>Decaying modes</u>	<ul style="list-style-type: none"> - Important to understand <ul style="list-style-type: none"> - transports in quantum systems - non-equilibrium topo. phase - Few algorithms, No demonstration in NISQ

Quantum master equation assumes

1. time-homogeneous, Markovian process
2. CPTP (completely-Positive and Trace-Preserving)

Our proposal: dissipative-system VQE (dVQE) algorithm

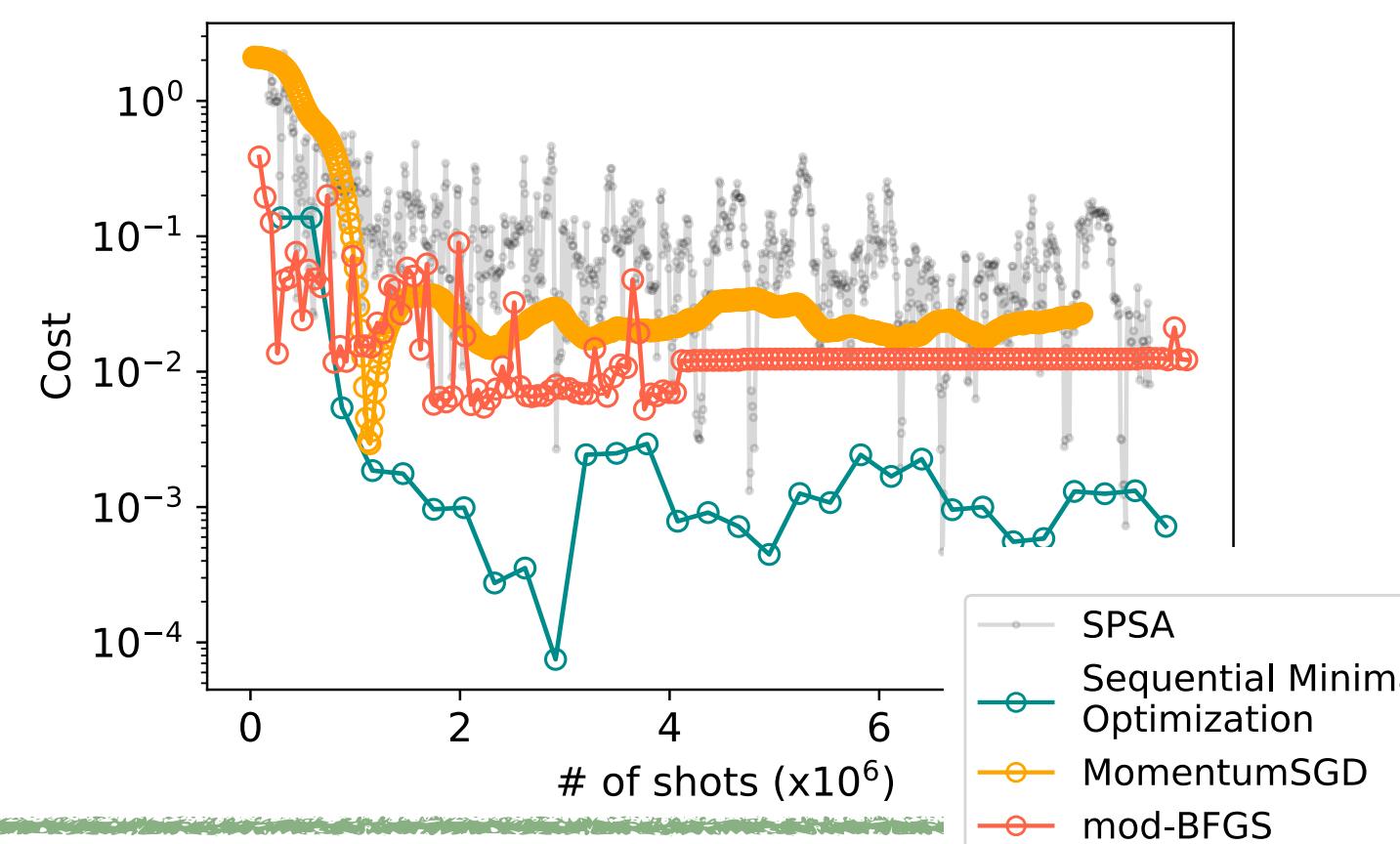
1. Encode mixed state in quantum circuit

$$|\rho_\theta\rangle = [V(\theta_v) \otimes V^*(\theta_v)] \\ \times \left(\prod_{n=1}^N \text{CNOT}_{n,n+N} \right) \tilde{D}(\theta_d) |0\rangle$$

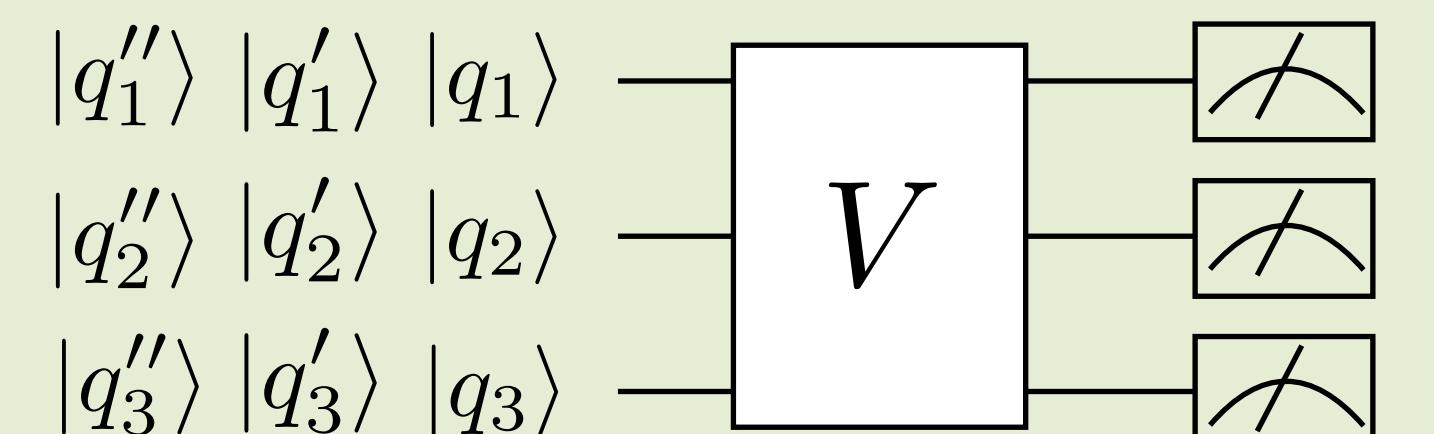
2. Define cost function

$$\arg \min_{\rho} \frac{\langle \rho | \hat{\mathcal{L}}^\dagger \hat{\mathcal{L}} | \rho \rangle}{\langle \rho | \rho \rangle}$$

Execute optimization



3. Measure observables



$$\rightarrow \langle A \rangle = \sum_{\mathbf{q}} \lambda_{\mathbf{q}} \langle \mathbf{q} | \hat{V}^\dagger \hat{A} \hat{V} | \mathbf{q} \rangle$$

1. Encoding mixed state into variational quantum circuit

Physical requirements for density matrices (general mixed state)

$$(I) \rho^\dagger = \rho$$

(Hermiticity)

$$(II) \langle \psi | \rho | \psi \rangle \geq 0, \forall |\psi\rangle \in \mathcal{H}$$

(Positivity)

$$(III) \text{Tr}[\rho] = 1$$

(Unit trace)

Encoded in ansatz

**Assured in actual measurement
(Step 3)**

Requirements (I),(II) in **matrix/vector** representation

Matrix rep.

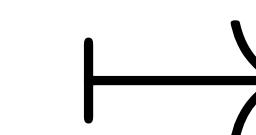
$$\rho = \sum_{ij} \rho_{ij} |i\rangle\langle j|$$

(I), (II) is equivalent to rewriting

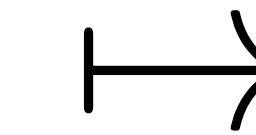
$$\rho = V D V^\dagger$$

$$D = \text{diag}(\{\lambda_q\})$$

where $\forall \lambda \geq 0$



Choi
isomorphism



Vector rep.

$$|\rho\rangle = \sum_{ij} \frac{\rho_{ij}}{C} |i\rangle_{\mathcal{P}} \otimes |j\rangle_{\mathcal{A}}$$

C: normalization
const.

States are mapped as

$$|\rho\rangle = [V \otimes V^*] |D\rangle$$

$$|D\rangle = \tilde{D}|0\rangle = \sum_q \frac{\lambda_q}{C} |q\rangle_{\mathcal{P}} \otimes |q\rangle_{\mathcal{A}}$$

1. Encoding mixed state into variational quantum circuit

Physical requirements for density matrices (general mixed state)

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Requirements (I),(II) in matrix/vector representation

Vector rep.

$$|\rho\rangle = \sum_{ij} \frac{\rho_{ij}}{C} |i\rangle_{\mathcal{P}} \otimes |j\rangle_{\mathcal{A}}$$

Choi isomorphism maps states to

$$|\rho\rangle = [V \otimes V^*] |D\rangle$$

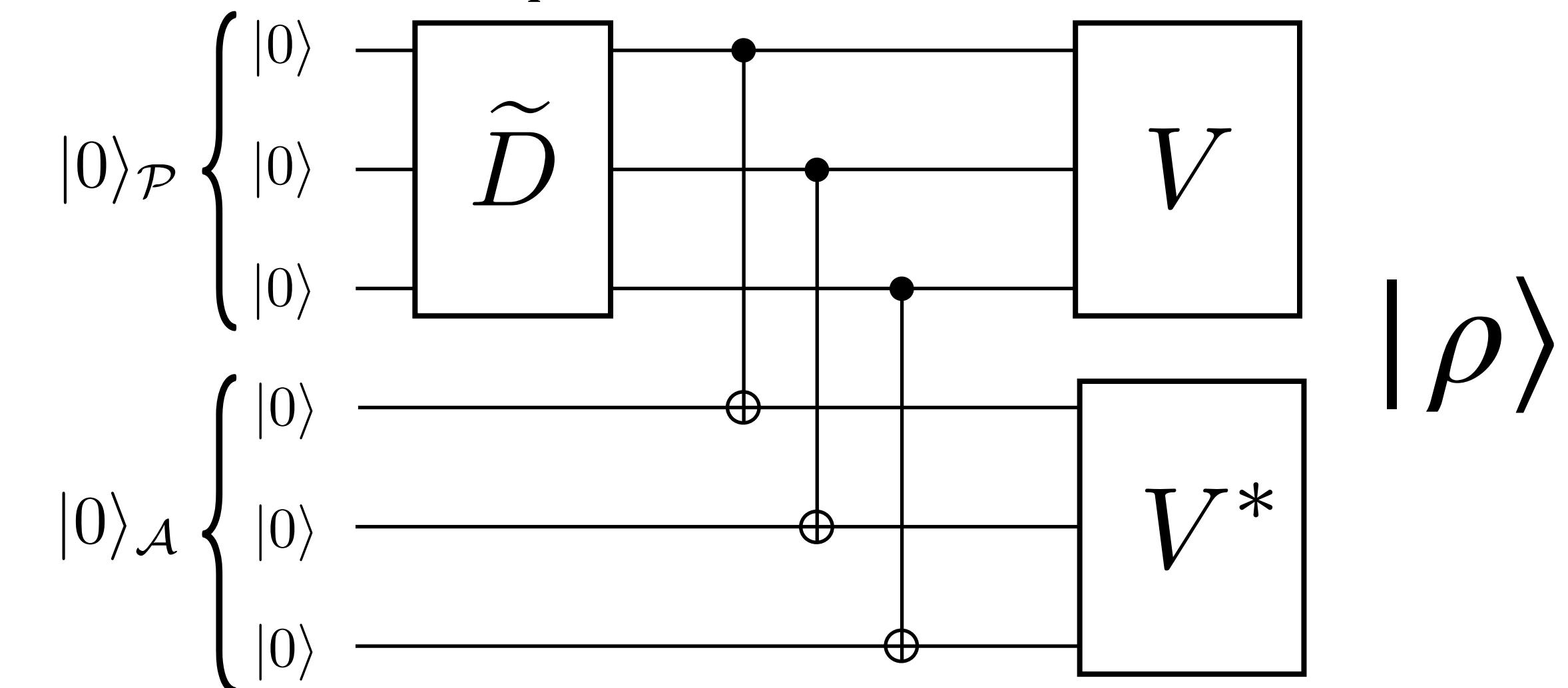
$$|D\rangle = \tilde{D}|0\rangle = \sum_{\mathbf{q}} \frac{\lambda_{\mathbf{q}}}{C} |\mathbf{q}\rangle_{\mathcal{P}} \otimes |\mathbf{q}\rangle_{\mathcal{A}}$$

Our ansatz

Generate $\{\lambda_{\mathbf{q}}\}$

Entangle as $|\mathbf{q}\rangle \otimes |\mathbf{q}\rangle$

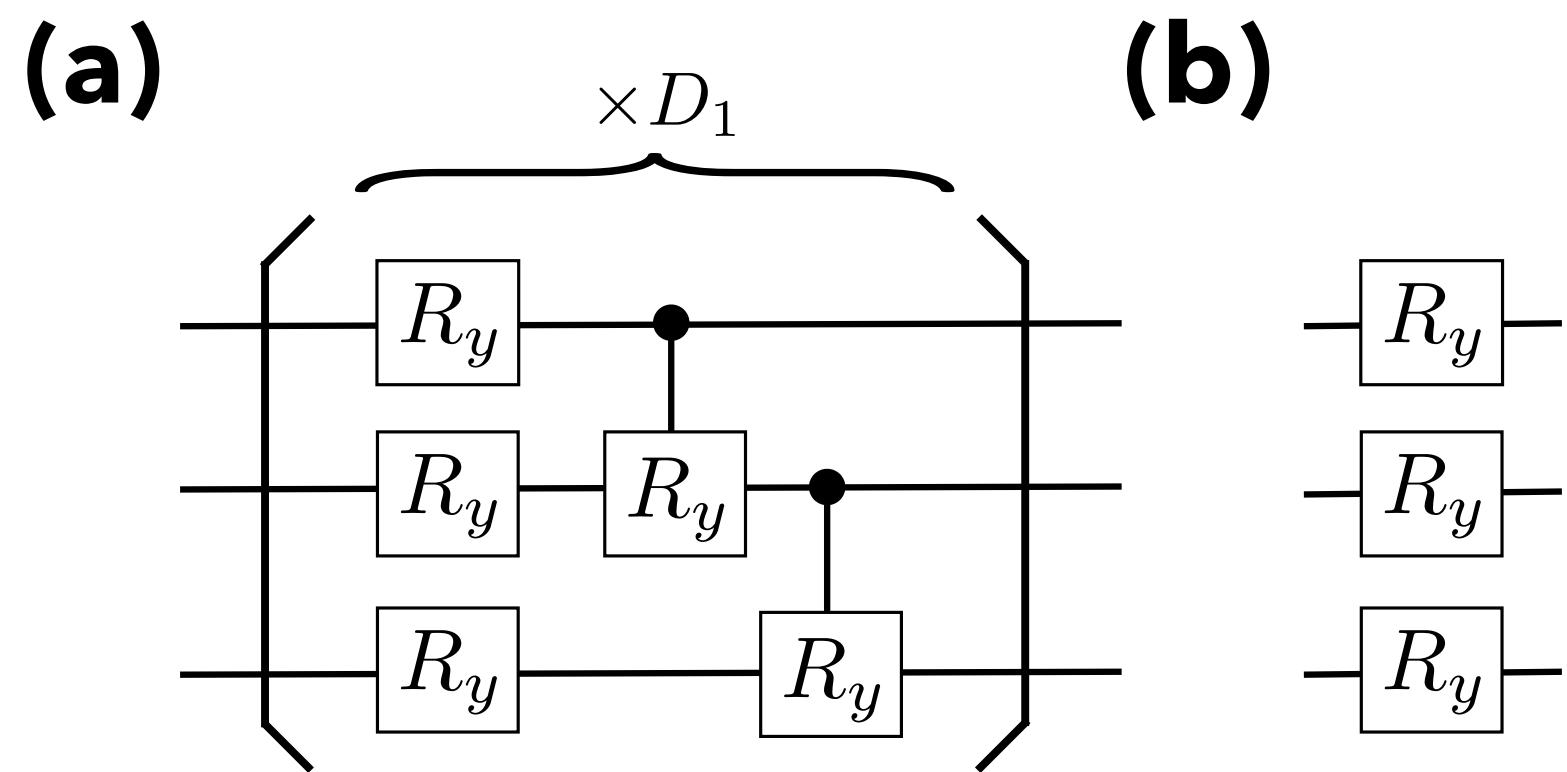
Basis rot.



\tilde{D} is chosen to produce positive wavefunction (see next page)

Examples of variational quantum circuits

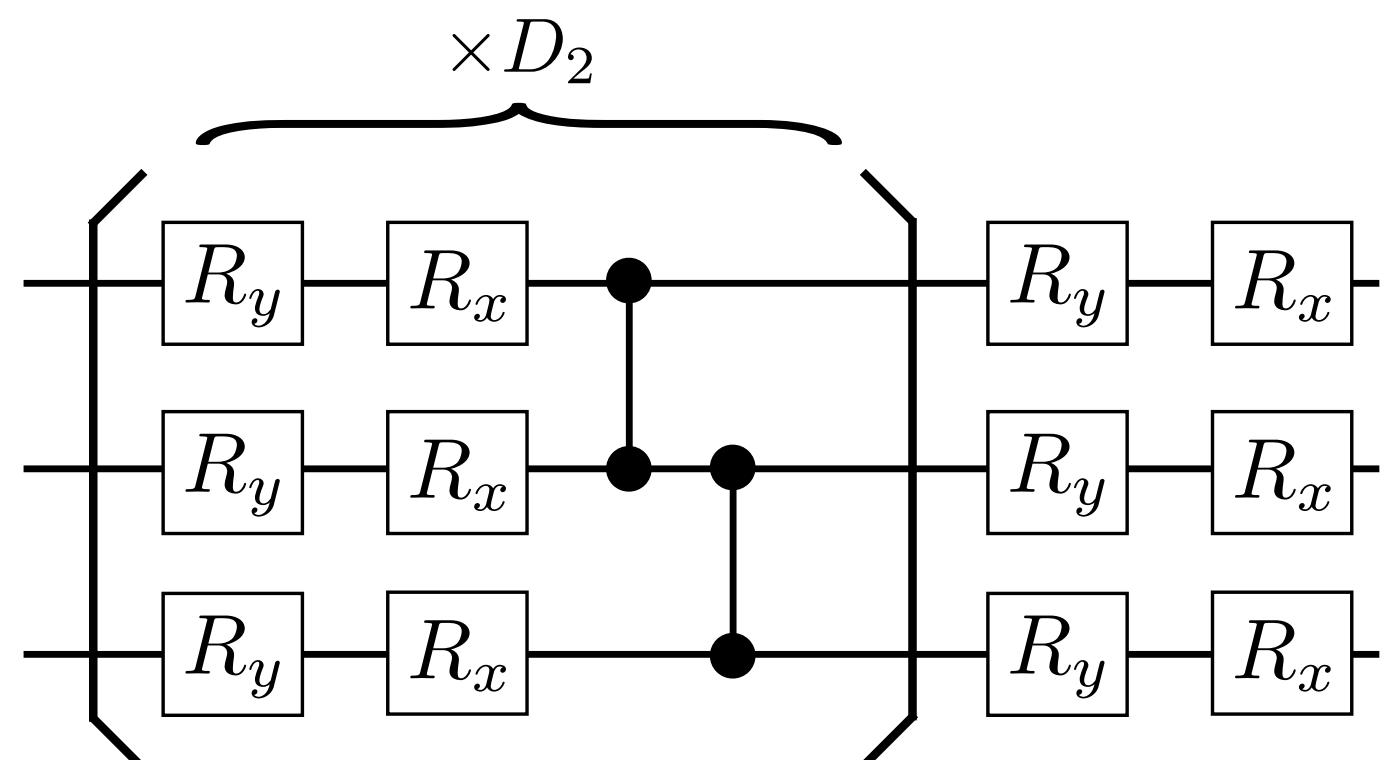
Circuit for eigenvalues $\{\lambda_q\}$: \tilde{D}



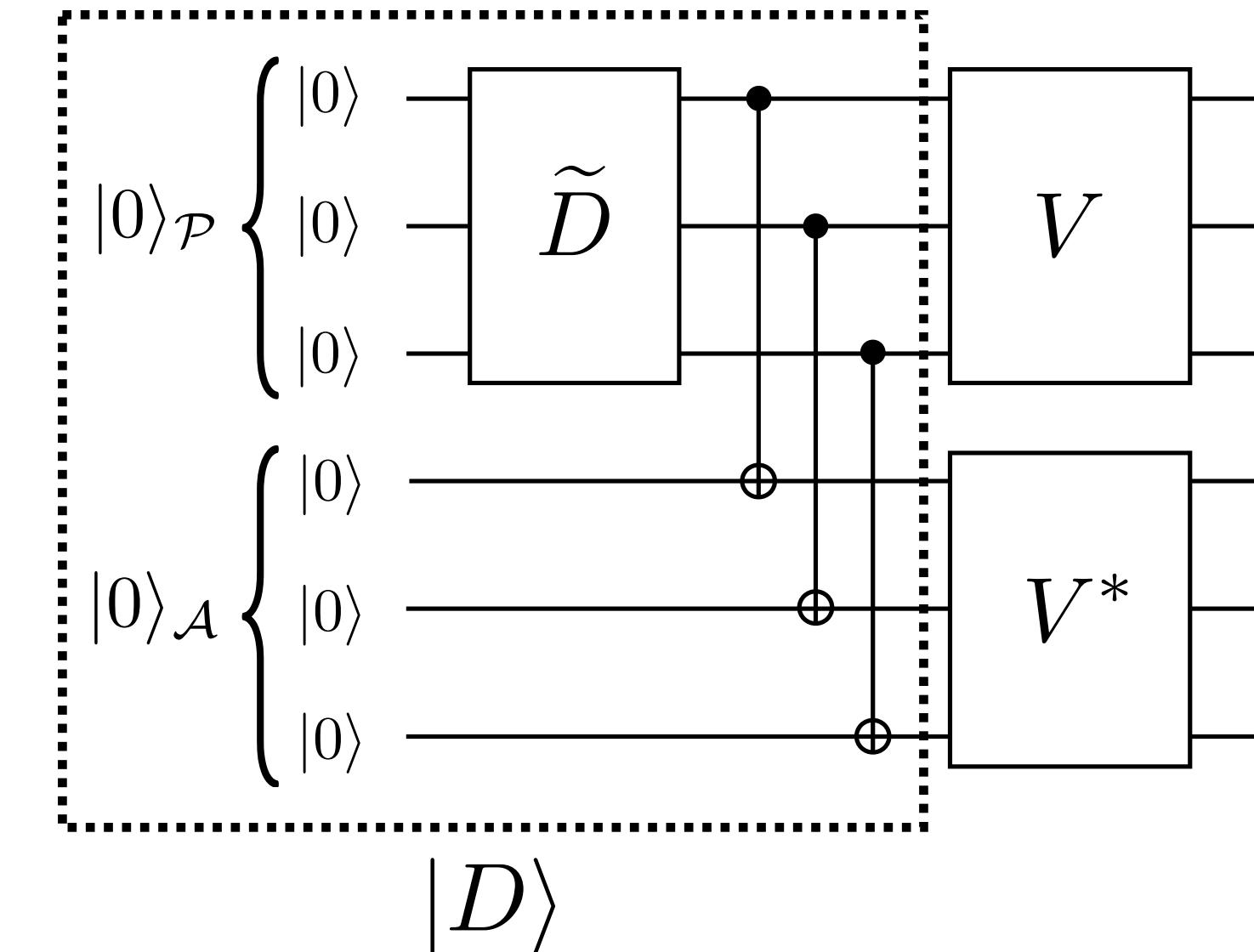
where angle θ of each $Ry = \exp(-i\theta Y/2)$ satisfies $0 \leq \theta \leq \pi$
so as to $\lambda_q \geq 0$

Circuit for basis transformation: V

e.g. hardware-efficient ansatz



Our ansatz



$$|\rho\rangle = [V \otimes V^*] |D\rangle$$

$$|D\rangle = \tilde{D}|0\rangle = \sum_{\mathbf{q}} \frac{\lambda_{\mathbf{q}}}{C} |\mathbf{q}\rangle_{\mathcal{P}} \otimes |\mathbf{q}\rangle_{\mathcal{A}}$$

$$\lambda_q \geq 0$$

2. Define cost function – variational search of steady states

Markovian master equation in Lindblad form

$$\hat{\mathcal{L}}|\rho(t)\rangle = \left(-i \left(\hat{H} \otimes \hat{1} - \hat{1} \otimes \hat{H}^T \right) + \sum_i \gamma_i \hat{\mathcal{D}}[\hat{\Gamma}_i] \right) |\rho(t)\rangle$$

Unitary evol. Non-unitary evol.

where

$$\hat{\mathcal{D}}[\hat{\Gamma}_i] = \hat{\Gamma}_i \otimes \hat{\Gamma}_i^* - \frac{1}{2} \hat{\Gamma}_i^\dagger \hat{\Gamma}_i \otimes \hat{1} - \hat{1} \otimes \frac{1}{2} \hat{\Gamma}_i^T \hat{\Gamma}_i^*$$

(e.g. $\hat{\Gamma}_i = \sigma_i^-$ for spontaneous emission, T₁ effect)

Stationary state as kernel

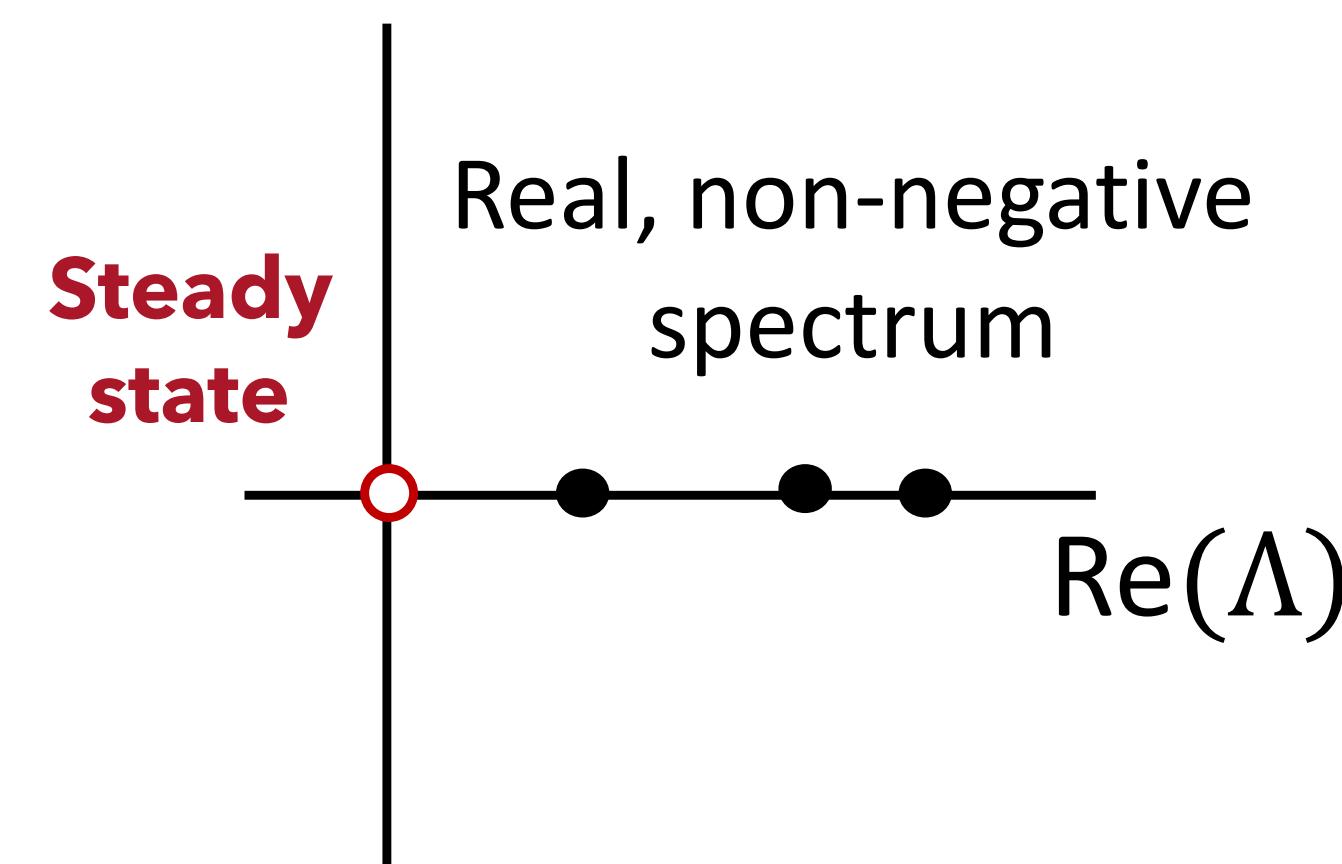
Cui et al. PRL ('15)

$$\hat{\mathcal{L}}|\rho\rangle = 0 \rightarrow \hat{\mathcal{L}}^\dagger \hat{\mathcal{L}}|\rho\rangle = 0$$

This allows us to obtain steady state via

$$\arg \min_{\rho} \frac{\langle \rho | \hat{\mathcal{L}}^\dagger \hat{\mathcal{L}} | \rho \rangle}{\langle \rho | \rho \rangle}$$

Spectrum of $\hat{\mathcal{L}}^\dagger \hat{\mathcal{L}}$



3. Measure Observables – efficient way

Naive way:
Vector Representation using $2N$ qubit

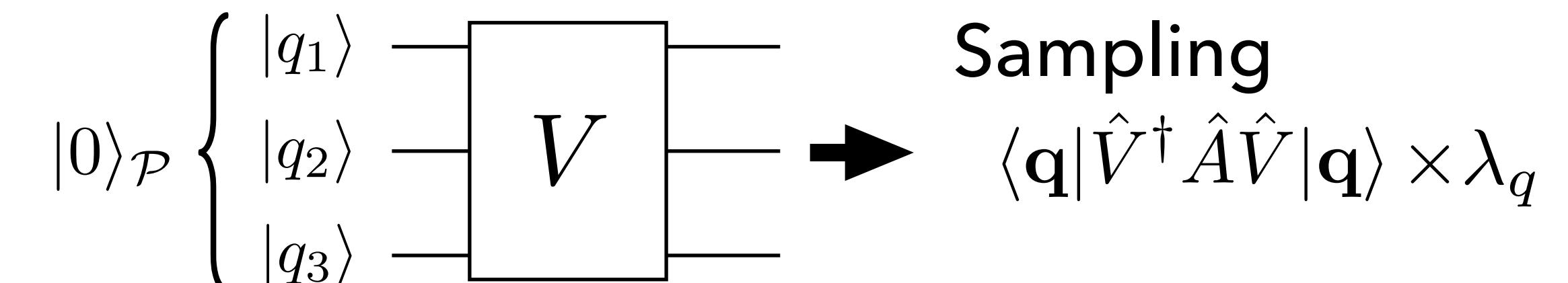
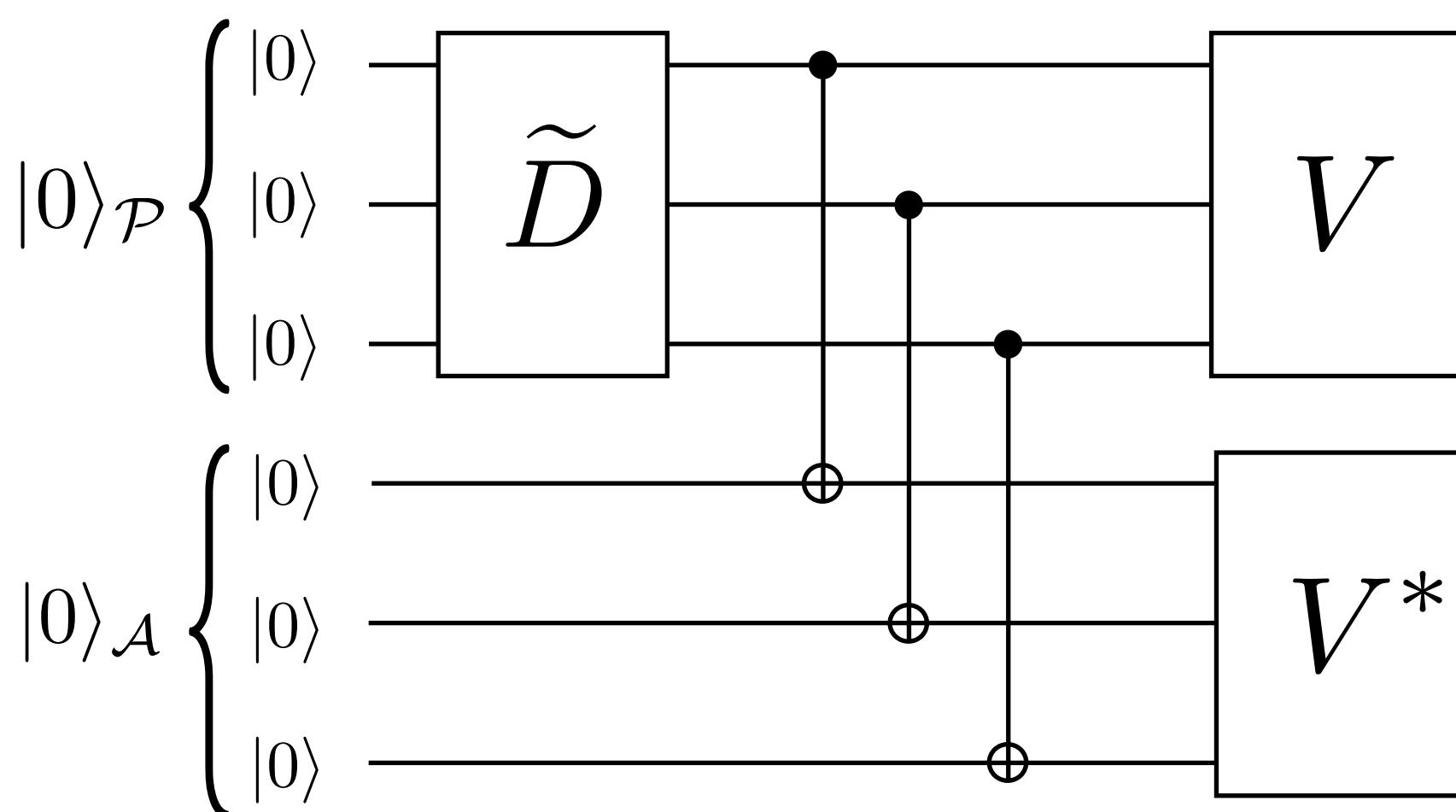
$$\begin{aligned}\langle \hat{A} \rangle &= \text{Tr}[\hat{\rho} \hat{A}] \\ &= \left(\sum_i \langle i |_{\mathcal{S}} \otimes \langle i |_{\mathcal{A}} \right) \hat{A} \otimes \hat{1} |\rho\rangle,\end{aligned}$$

- Post-select “diagonal terms”
- Exponentially inefficient

Our proposal:
Matrix Representation using N qubit

$$\begin{aligned}\langle A \rangle &= \text{Tr}[\hat{\rho} \hat{A}] = \sum_{\mathbf{q}} \lambda_{\mathbf{q}} \langle \psi_{\mathbf{q}} | \hat{A} | \psi_{\mathbf{q}} \rangle \\ &= \sum_{\mathbf{q}} \lambda_{\mathbf{q}} \langle \mathbf{q} | \hat{V}^\dagger \hat{A} \hat{V} | \mathbf{q} \rangle\end{aligned}$$

- Determine $\{\lambda_{\mathbf{q}}\}$ from sampling
- Measure with initial bit \mathbf{q} with weight $\lambda_{\mathbf{q}}$



Demonstration of our proposal in quantum Ising model

Model

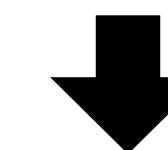
Hamiltonian: ZZ coupling and traverse-field

$$H = \frac{1}{2} \sum_i \sigma_i^z \sigma_{i+1}^z + g \sum_i \sigma_i^x$$

Dissipation: Local damping & dephasing

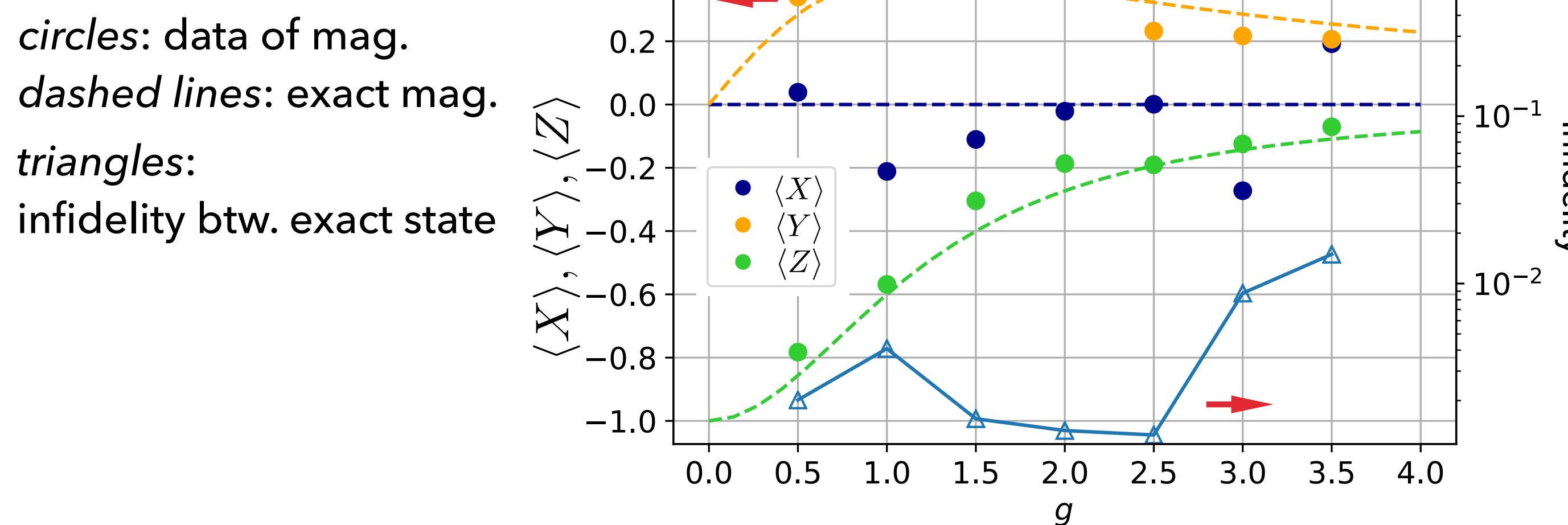
$$c_i^{(1)} = \sigma_i^- \quad c_i^{(2)} = \sigma_i^z \quad \text{with amplitudes } \gamma_1 = 1, \gamma_2 = 0.5$$

$$\hat{\mathcal{L}} |\rho\rangle = \left(-i(H \otimes \mathbb{1} - \mathbb{1} \otimes H^T) + \sum_a \sum_{i=0}^{N-1} \gamma_a \hat{\mathcal{D}}[c_i^{(a)}] \right) |\rho\rangle \quad \hat{\mathcal{D}}[\hat{\Gamma}_i] = \hat{\Gamma}_i \otimes \hat{\Gamma}_i^* - \frac{1}{2} \hat{\Gamma}_i^\dagger \hat{\Gamma}_i \otimes \hat{1} - \hat{1} \otimes \frac{1}{2} \hat{\Gamma}_i^T \hat{\Gamma}_i^*$$

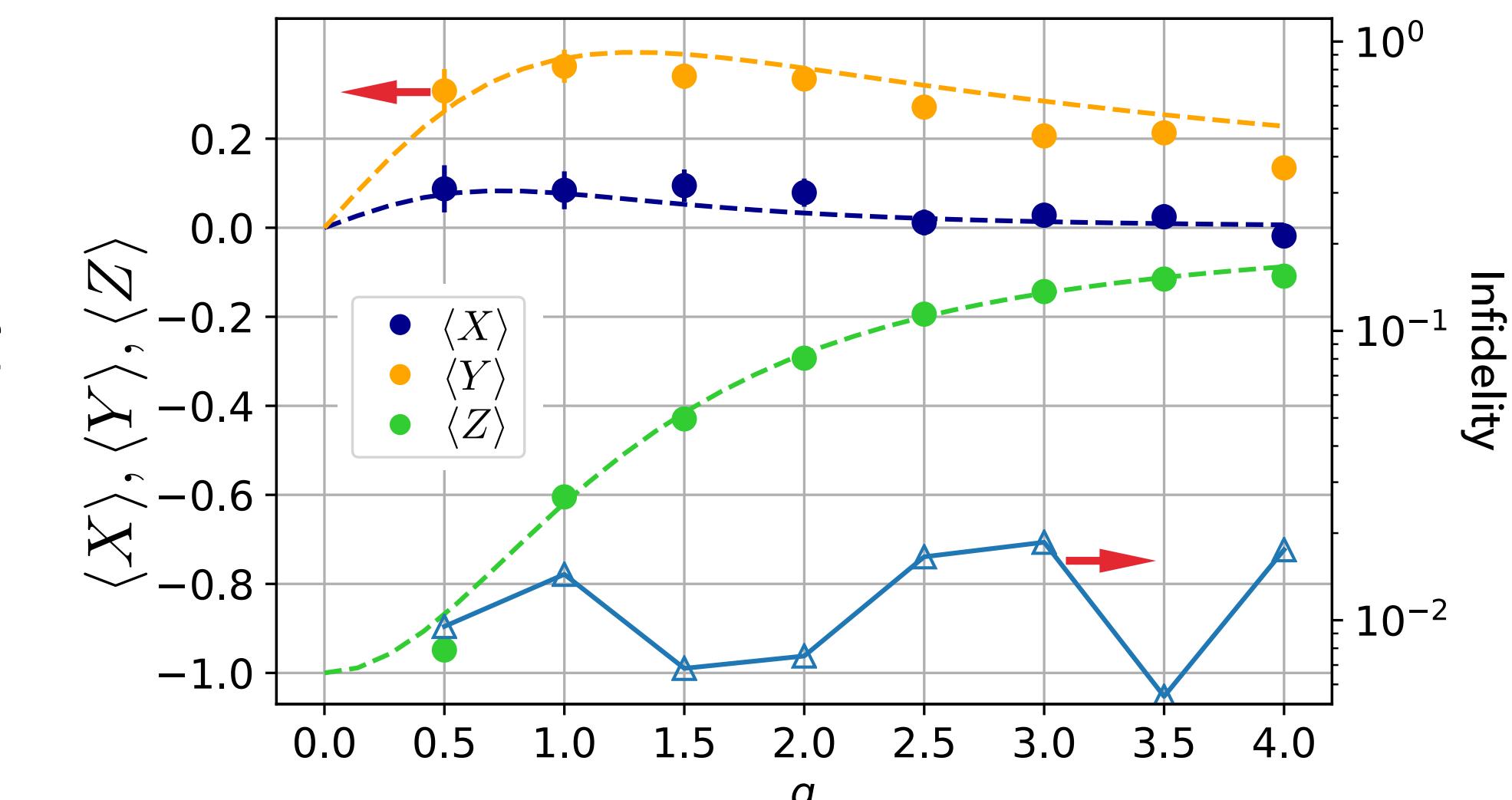


Calculating magnetization curve and density matrix fidelity by dVQE

Quantum simulation on real NISQ device
(Aspen-4, Rigetti, 2N=2 qubits)

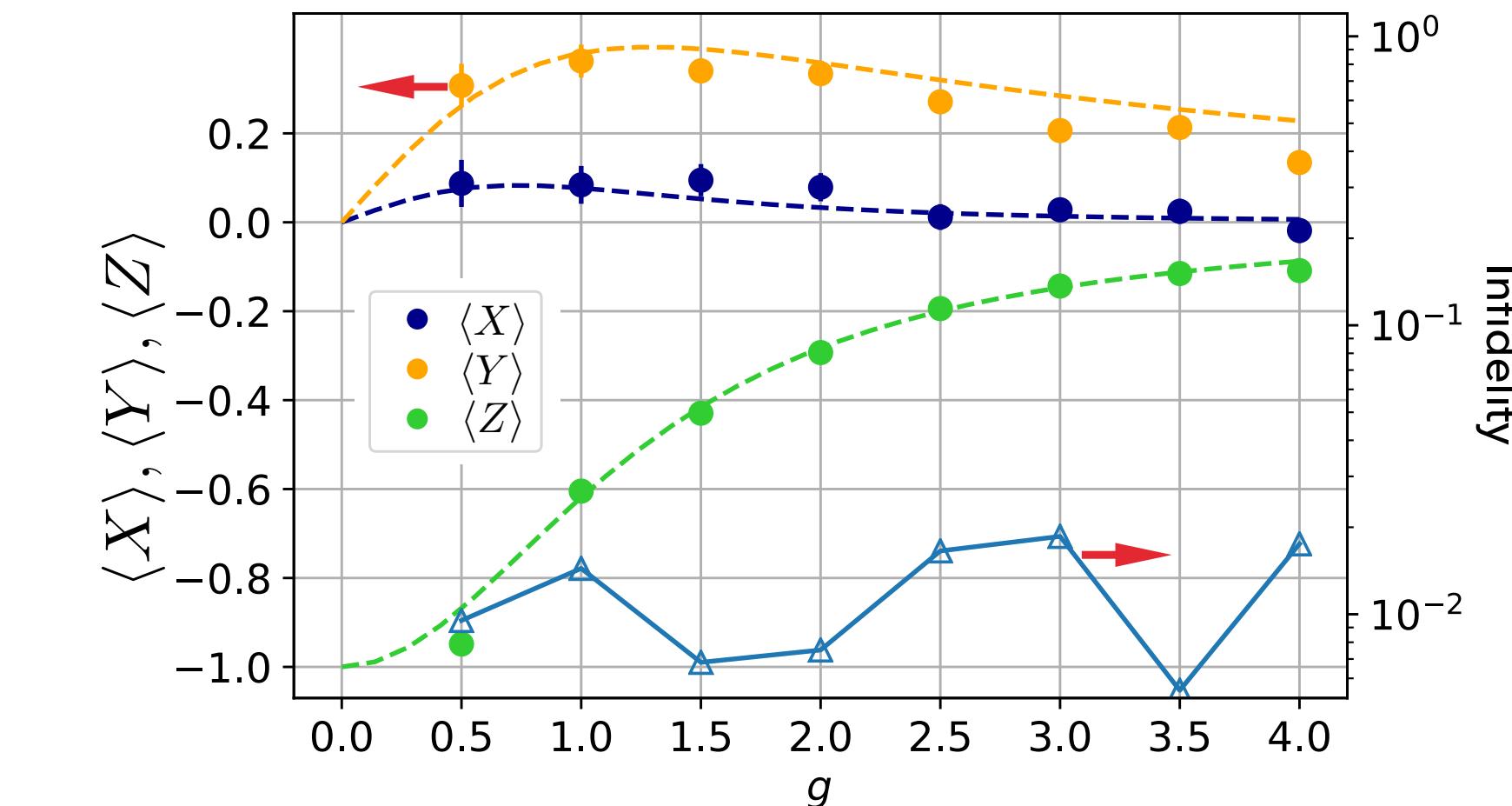
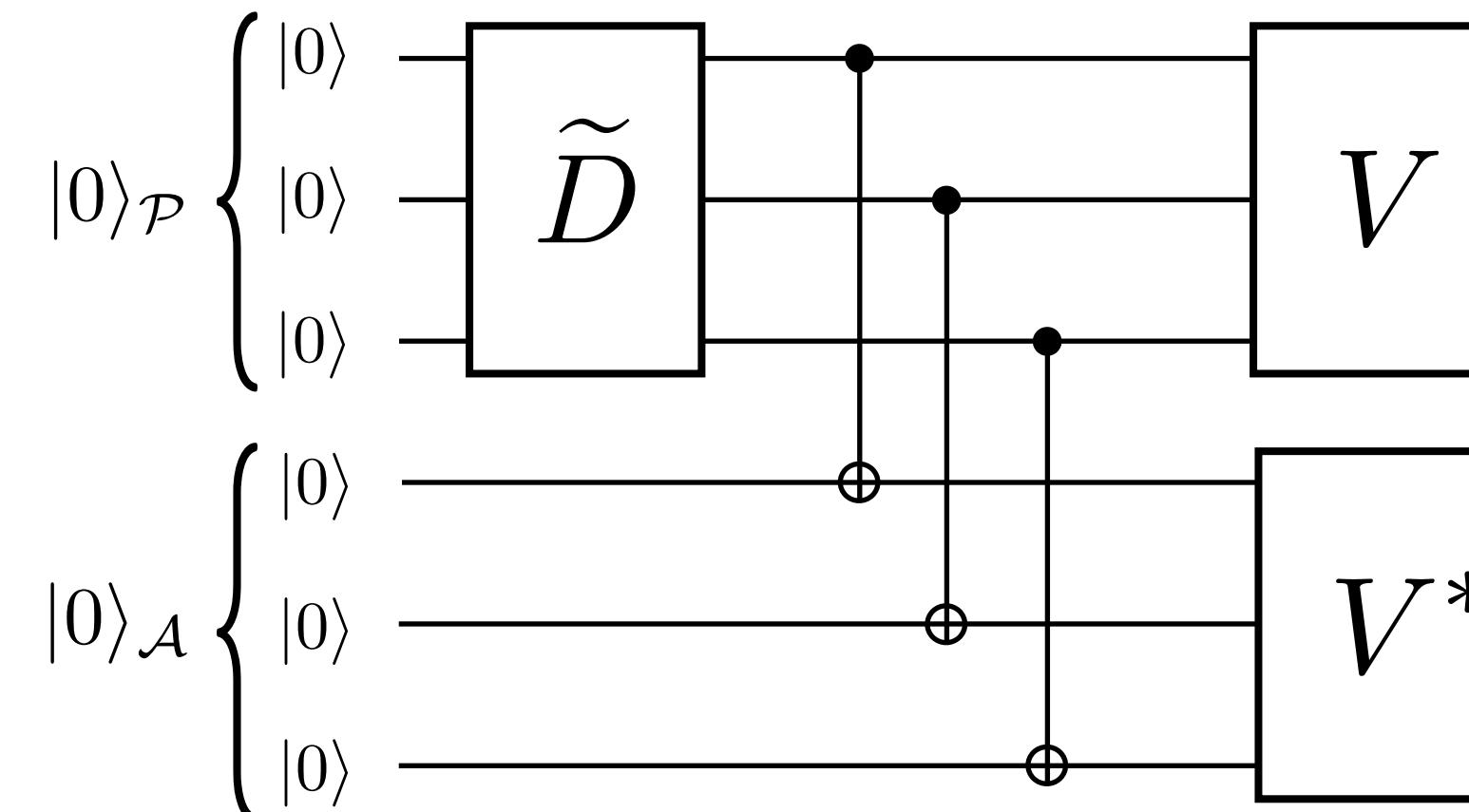


Numerical simulation, 2N=4 qubits



Summary

- VQE-based algorithm for steady states in Markovian open quantum system (**dVQE**)
- Numerical/Experimental demonstration in dissipative quantum Ising model
- Applications to transport properties in nano devices/molecular systems
dissipative phase exploration in condensed matter systems are expected



Yoshioka, Nakagawa, Mitarai, and Fujii, [arXiv: 1908.09836](https://arxiv.org/abs/1908.09836)

SUPPLEMENTARY MATERIALS

Encoding mixed states into pure states

Choi isomorphism

$$\rho = \sum_{ij} \rho_{ij} |i\rangle\langle j| \quad \mapsto |\rho\rangle = \sum_{ij} \frac{\rho_{ij}}{C} |i\rangle_{\mathcal{P}} \otimes |j\rangle_{\mathcal{A}}$$
$$\hat{\rho} = \left(\begin{array}{c|c|c|c} \text{---} & \text{---} & \text{---} & \text{---} \\ \text{---} & \text{---} & \text{---} & \text{---} \\ \text{---} & \text{---} & \text{---} & \text{---} \\ \text{---} & \text{---} & \text{---} & \text{---} \end{array} \right) \mapsto |\rho\rangle\rangle = \left(\begin{array}{cccc} \text{---}^T & \text{---}^T & \text{---}^T & \text{---}^T \\ \text{---} & \text{---} & \text{---} & \text{---} \end{array} \right)$$

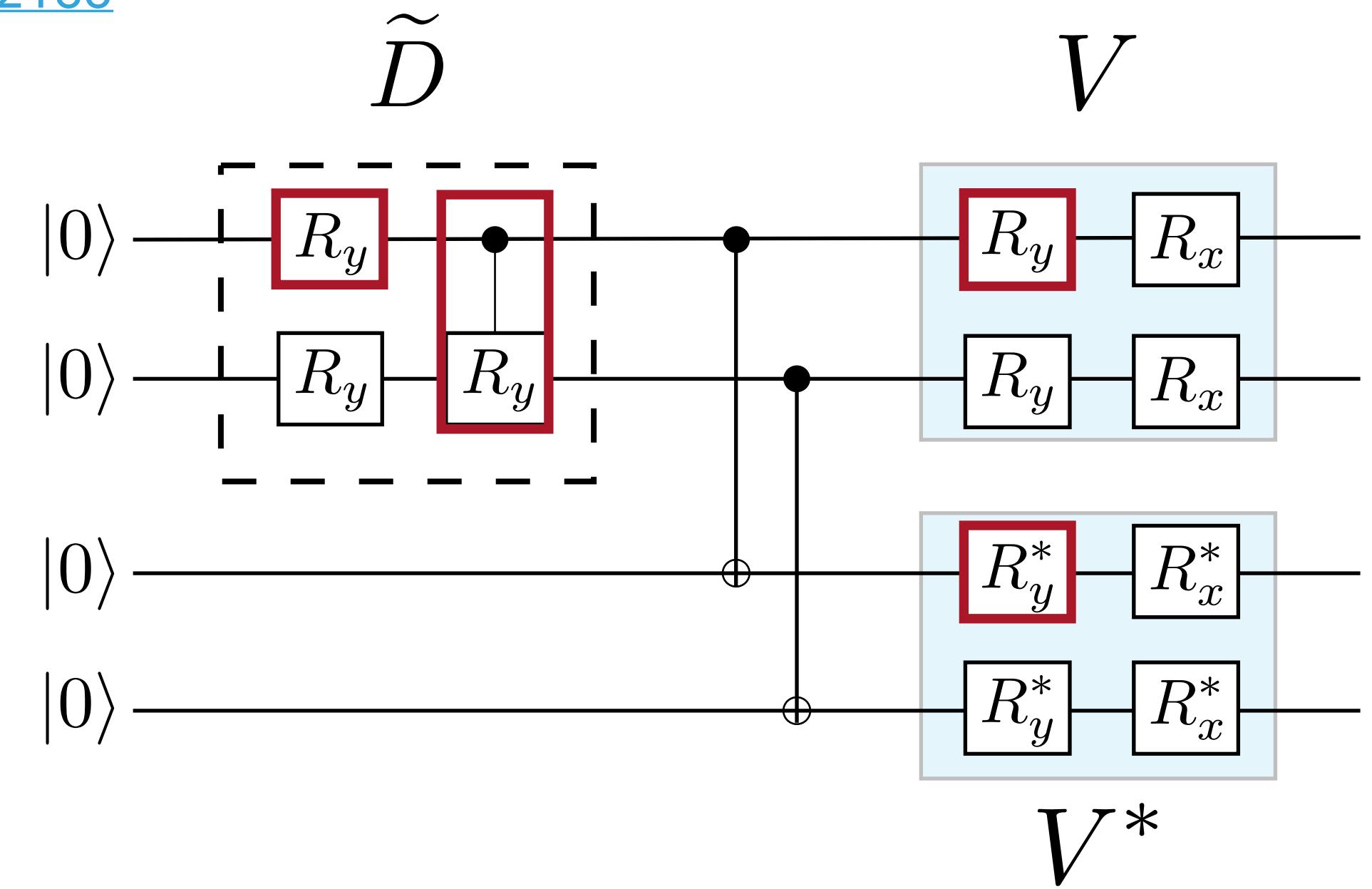
Map N qubit mixed state → 2N qubit pure state

Details of optimization – sequential minimal optimization

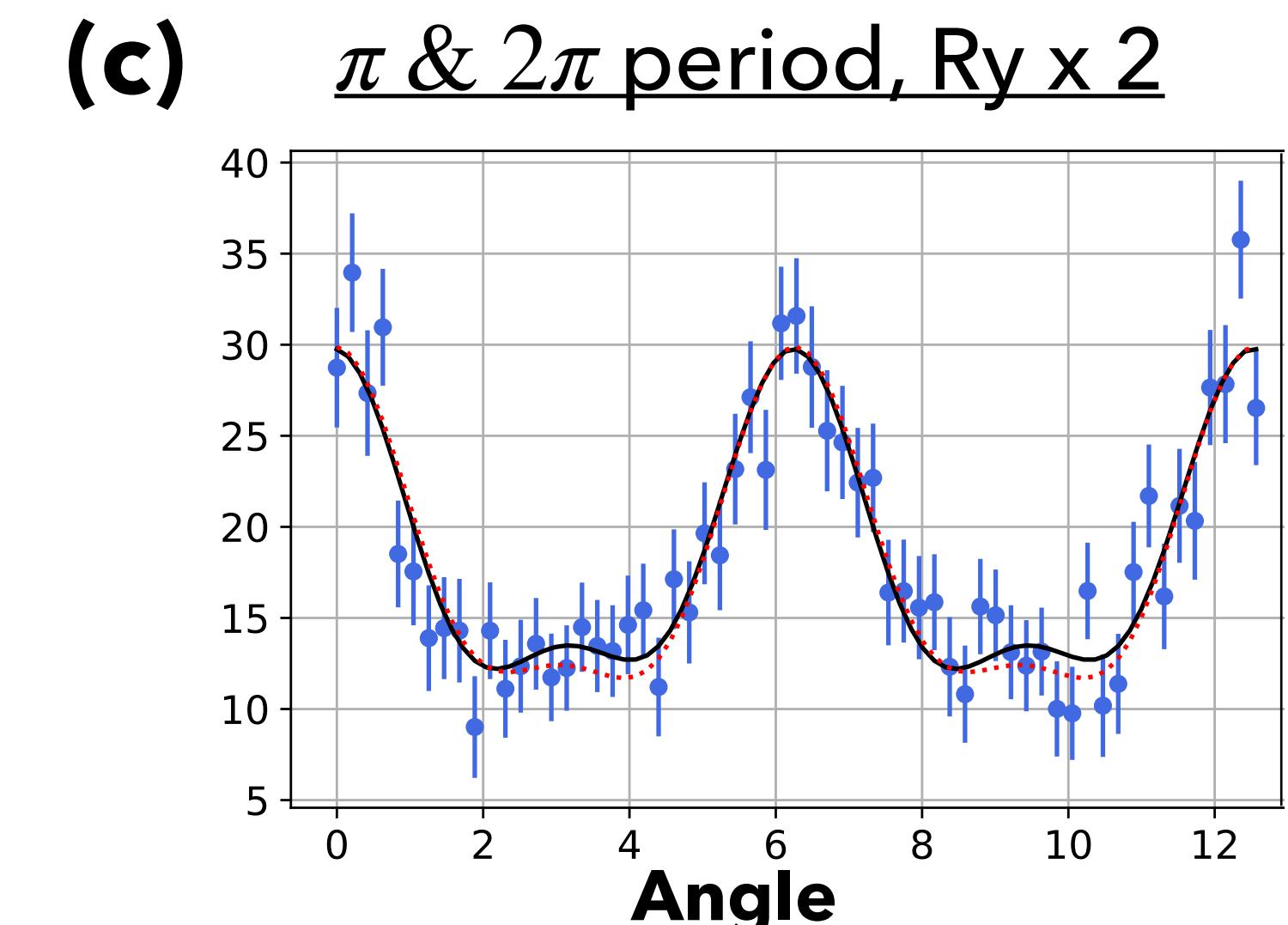
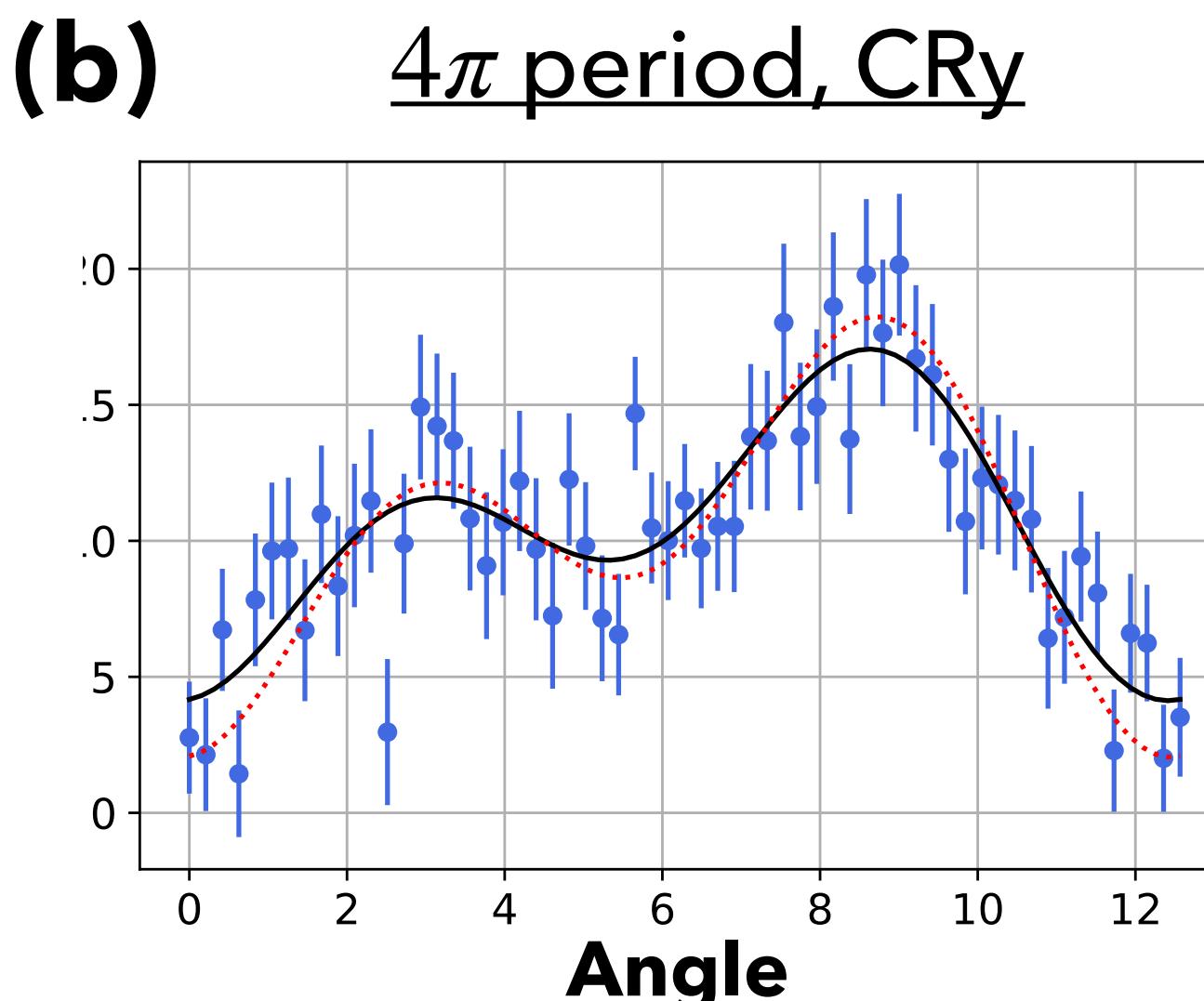
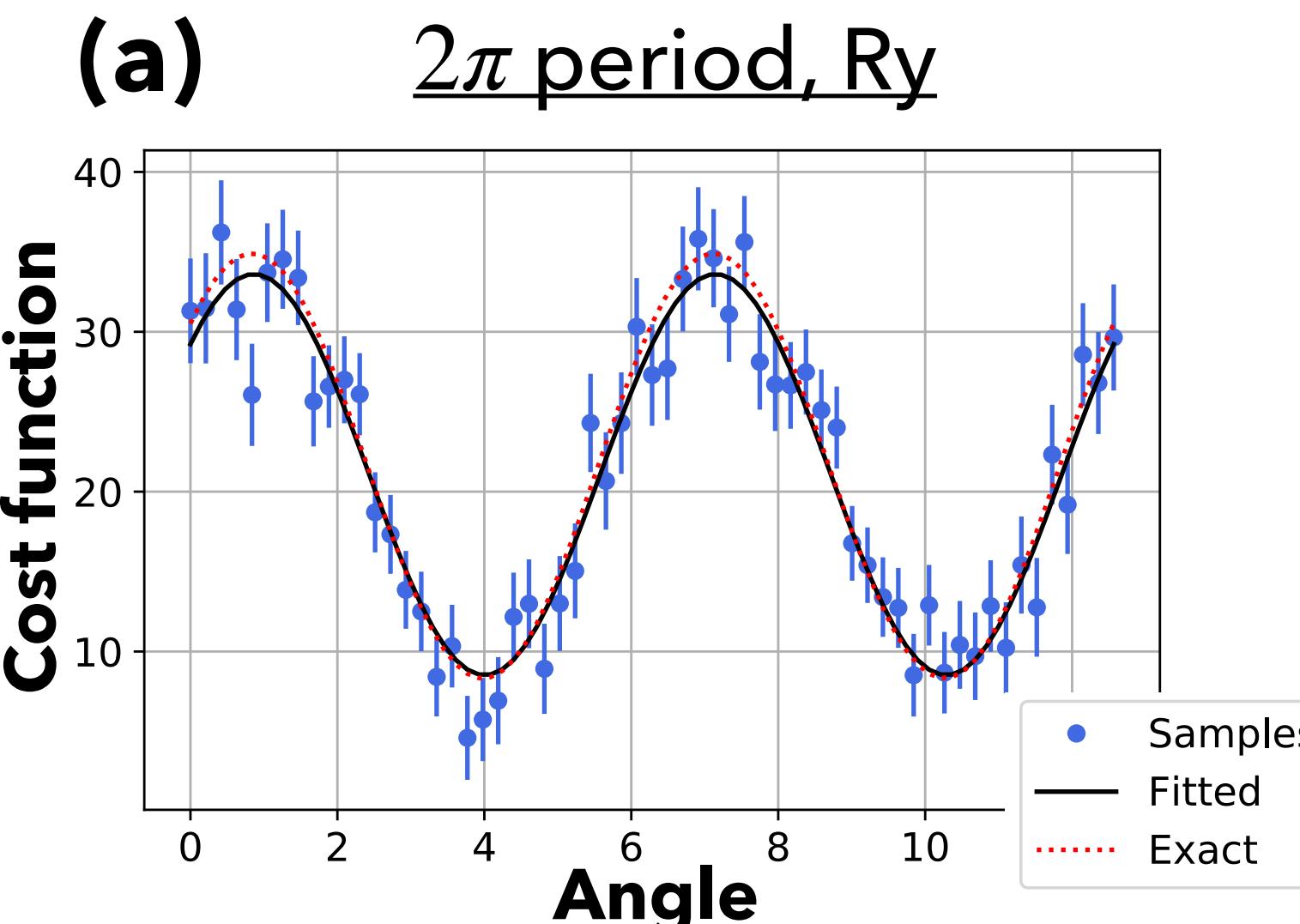
Nakanishi-Todo-Fujii method

Nakanishi, Fuji,&Todo, [arXiv:1903.12166](https://arxiv.org/abs/1903.12166)

- Use periodicity of cost function w.r.t. angles θ
- Sine-curve-like structure with periods:
 - 2π for ordinary rotation gates
 - $2^{M+1}\pi$ for $C^{(M)}$ -Rotation
 - $2\pi/M$ if θ appears M times



Cost function landscape in $2N=4$ qubits



Gate error mitigation

Quantum simulation with redundancy

Heya et al. , [arXiv:1904.08566](https://arxiv.org/abs/1904.08566)

1. Encode redundancy in the circuit using equation

$$R_x(\pm\pi/2) \equiv R_z(\pi) (R_x(\pm\pi/2))^{\mathcal{E}} R_z(\pi)$$

$$CZ \equiv CZ^{\mathcal{E}}$$

for odd $\mathcal{E} \in \{1, 3, 5, \dots\}$

2. Sample cost function
3. Mitigate by linear extrapolation w.r.t ϵ

Numerical simulation: depolarizing

1. Applying depolarizing channel for each k-qubit gates as

$$\rho \mapsto \rho' = (1 - p_k) \rho + \frac{4^k - 1}{4^k} p_k \mathbb{I}$$

rate

2. Sample cost function
3. Mitigate by linear extrapolation w.r.t. p_k

