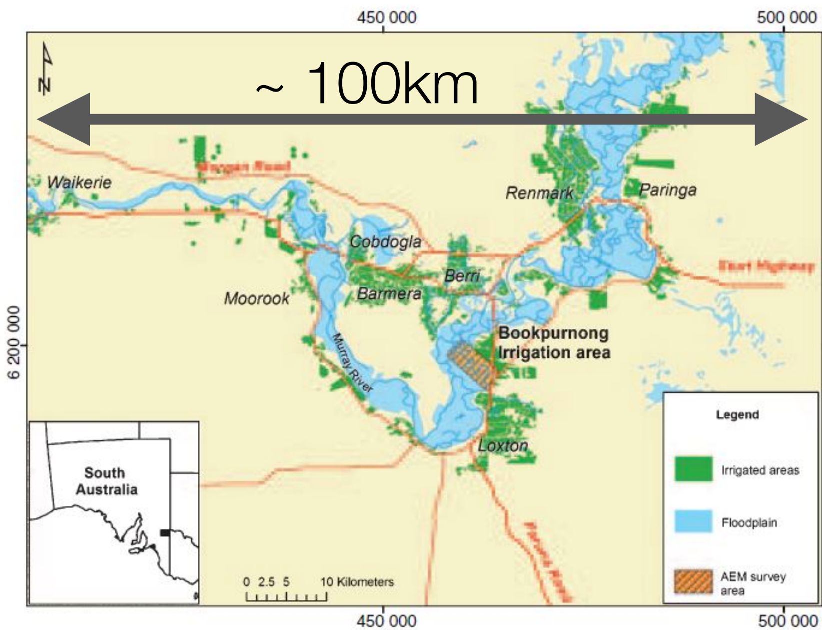
EM Fundamentals

DISC 2017
Electromagnetics
Fundamentals and Applications

Motivation: applications difficult for DC

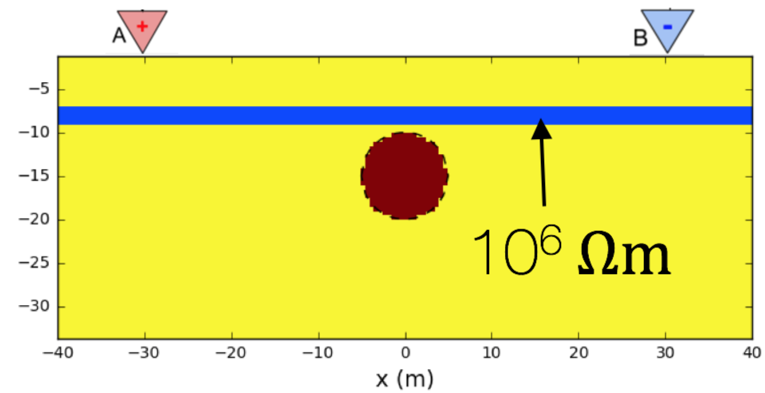
Large areas to be covered

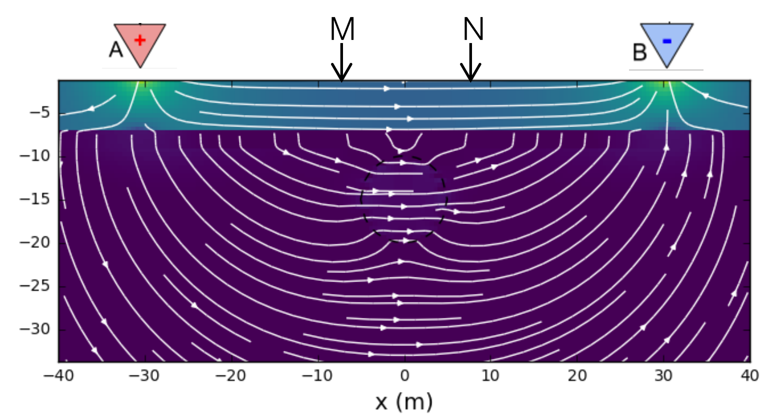


Rugged terrain

Hard to inject

Resistive layer “shields” target



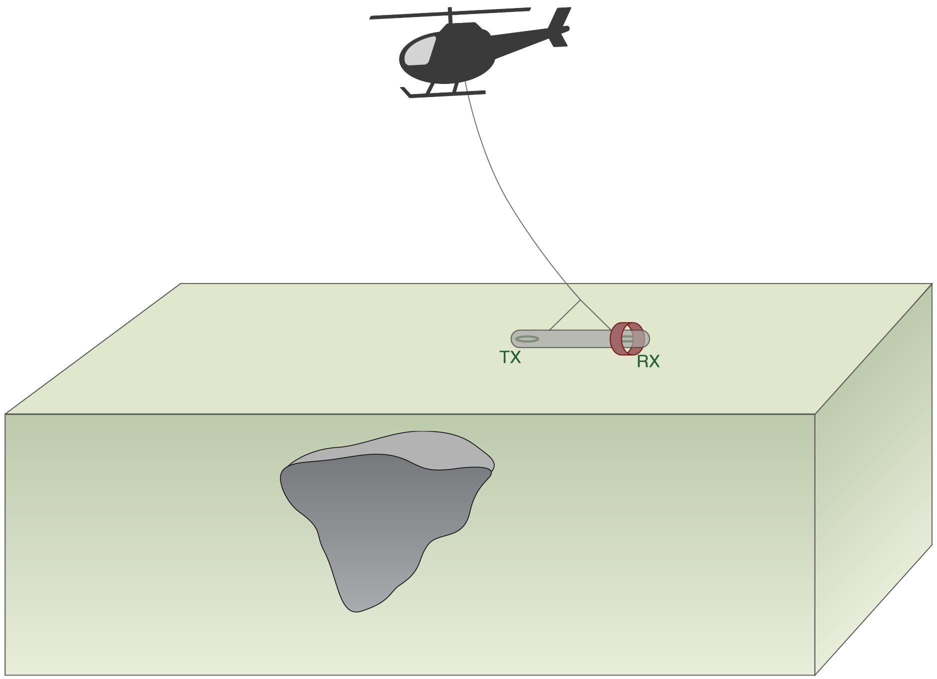


Outline

- Basic Survey
- Ampere's and Faraday's Laws (2-coil App)
- Circuit model for EM induction
- Frequency and time domain data
- Sphere in homogeneous earth
- Cyl code
- Energy losses in the ground

Basic Experiment

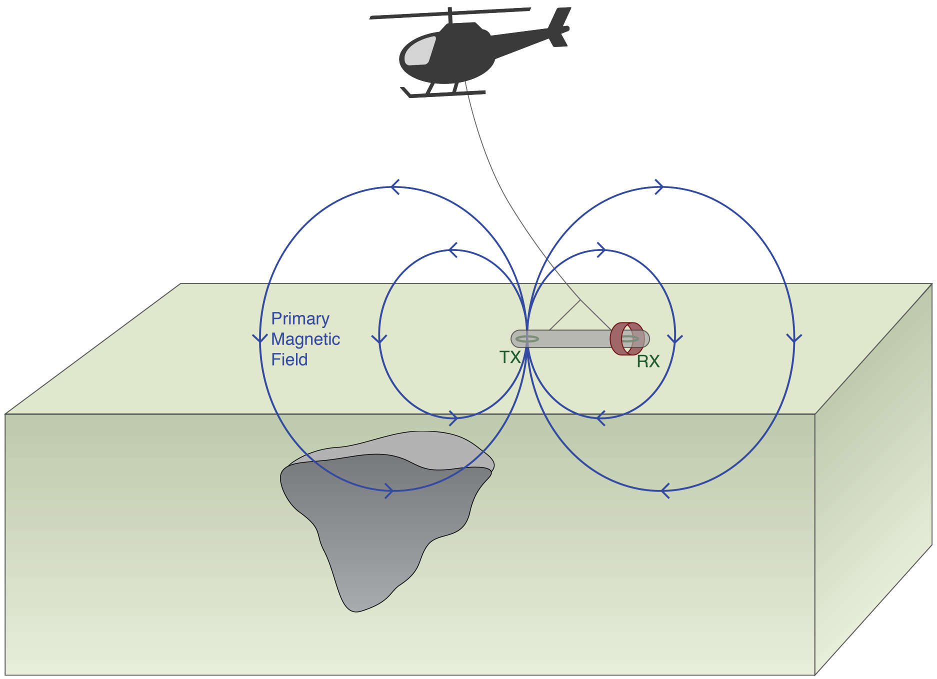
- **Setup:**
  - transmitter and receiver are in a towed bird



Basic Experiment

- **Setup:**
  - transmitter and receiver are in a towed bird

- **Primary:**
  - Transmitter produces a primary magnetic field



Basic Experiment

- **Setup:**
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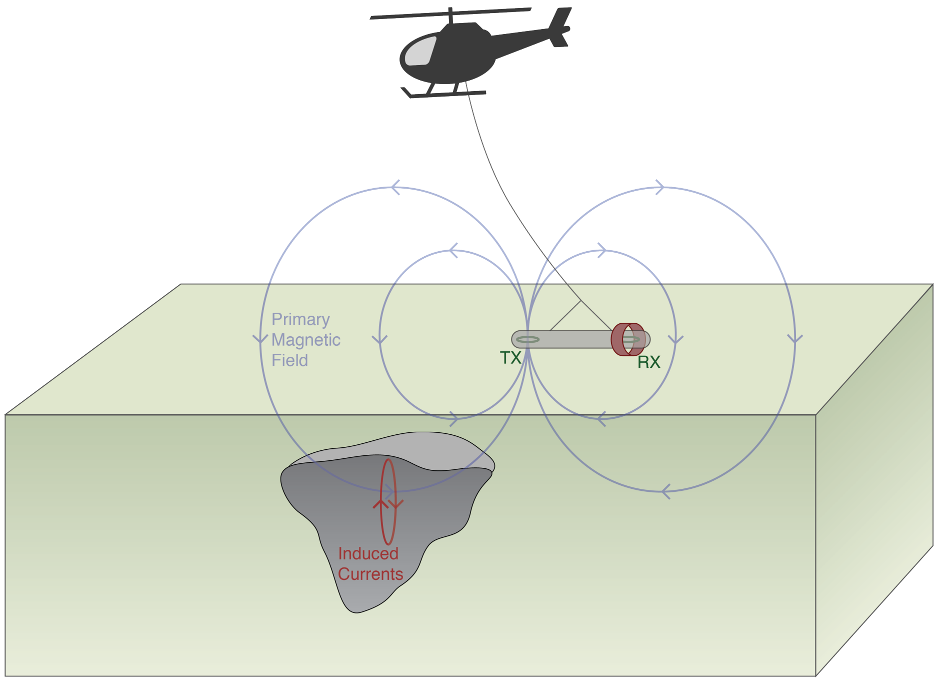
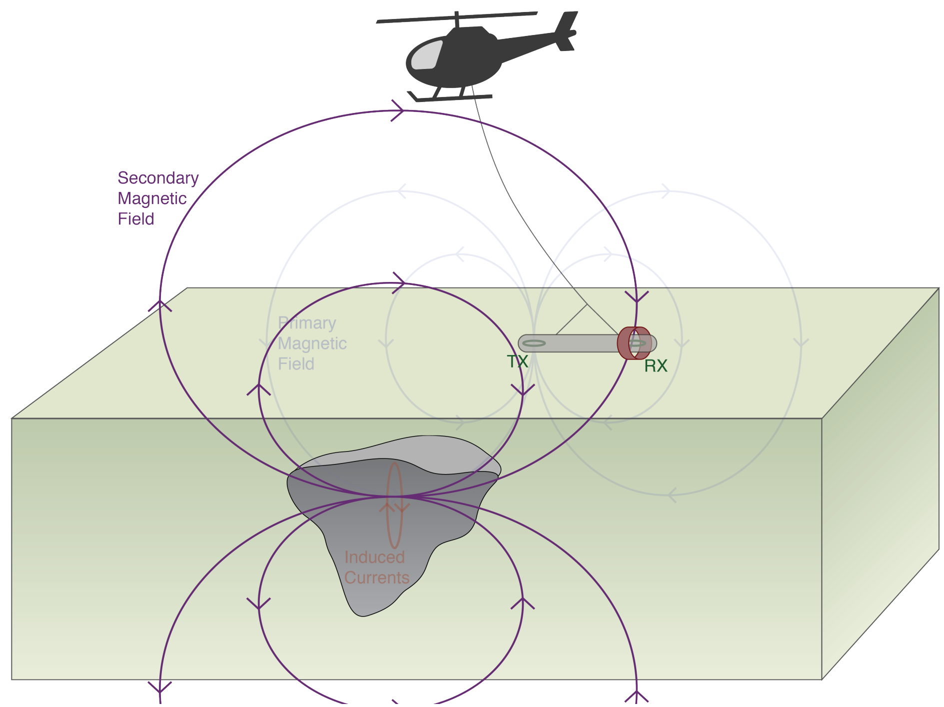
- **Primary:**
  - Transmitter produces a primary magnetic field
- **Induced Currents:**
  - Time varying magnetic fields generate electric fields everywhere and currents in conductors

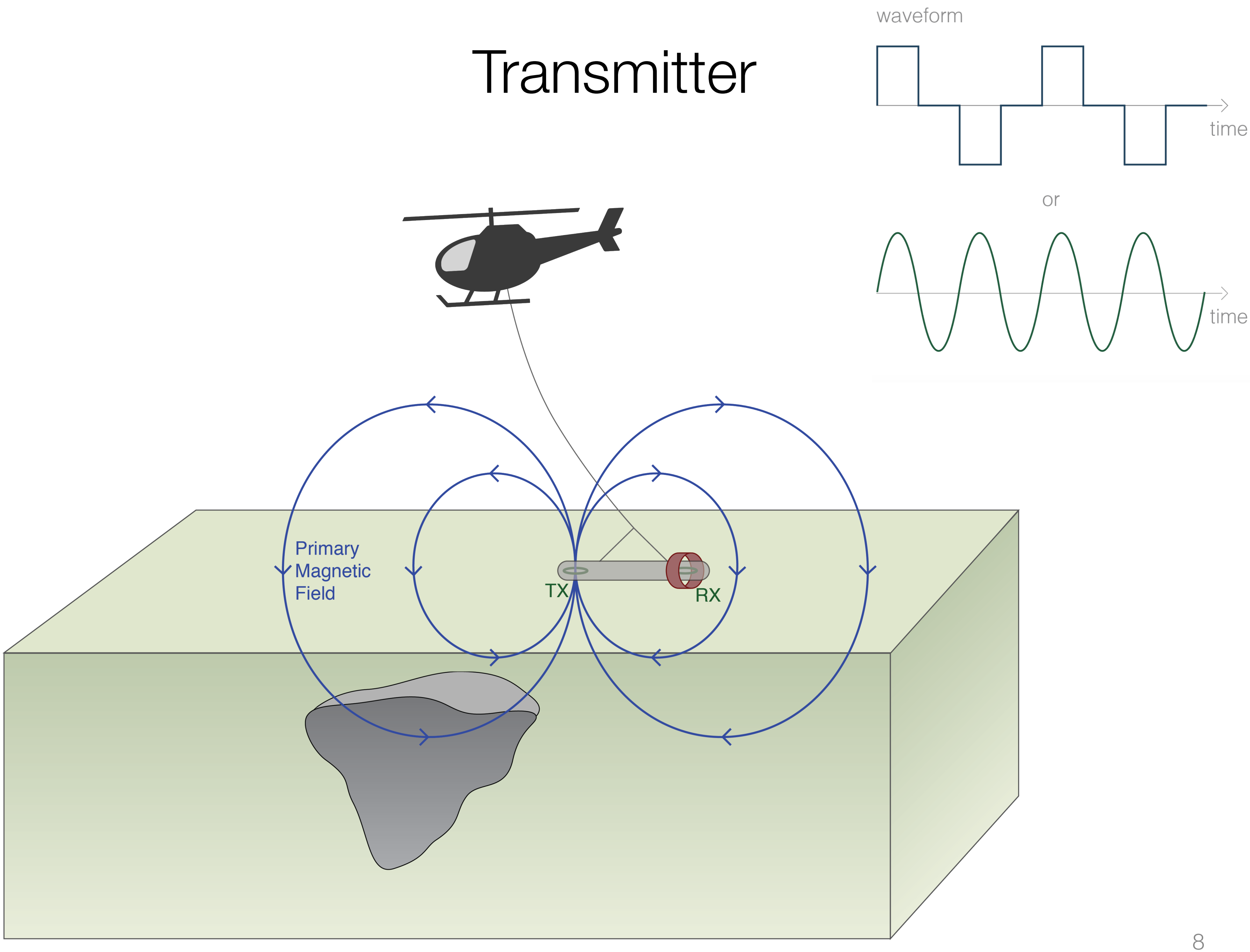


Basic Experiment

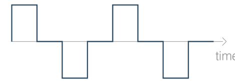
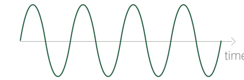
- **Setup:**
  - transmitter and receiver are in a towed bird
- **Primary:**
  - Transmitter produces a primary magnetic field
- **Induced Currents:**
  - Time varying magnetic fields generate electric fields everywhere and currents in conductors
- **Secondary Fields:**
  - The induced currents produce a secondary magnetic field.



Transmitter



Basic Equations: Quasi-static

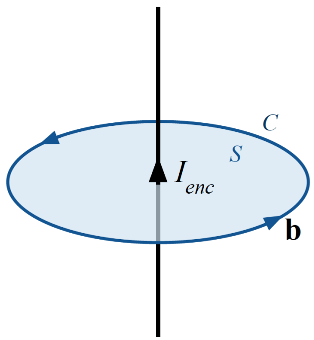
	Time 	Frequency 
Faraday's Law	$\nabla \times \mathbf{e} = -\dfrac{\partial \mathbf{b}}{\partial t}$	$\nabla \times \mathbf{E} = -i\omega \mathbf{B}$
Ampere's Law	$\nabla \times \mathbf{h} = \mathbf{j} + \dfrac{\partial \mathbf{d}}{\partial t}$	$\nabla \times \mathbf{H} = \mathbf{J} + i\omega \mathbf{D}$
No Magnetic Monopoles	$\nabla \cdot \mathbf{b} = 0$	$\nabla \cdot \mathbf{B} = 0$
Constitutive Relationships (non-dispersive)	$\mathbf{j} = \sigma \mathbf{e}$ $\mathbf{b} = \mu \mathbf{h}$ $\mathbf{d} = \varepsilon \mathbf{e}$	$\mathbf{J} = \sigma \mathbf{E}$ $\mathbf{B} = \mu \mathbf{H}$ $\mathbf{D} = \varepsilon \mathbf{E}$

* Solve with sources and boundary conditions

Ampere's Law

$$\nabla \times \mathbf{H} = \mathbf{J}$$

Wire



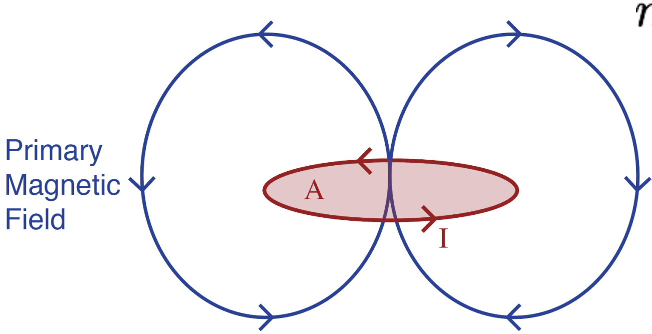
$$\mathbf{B} = \frac{\mu_0 I_{enc}}{2\pi r}\hat{\phi}$$

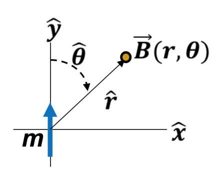
Right hand rule

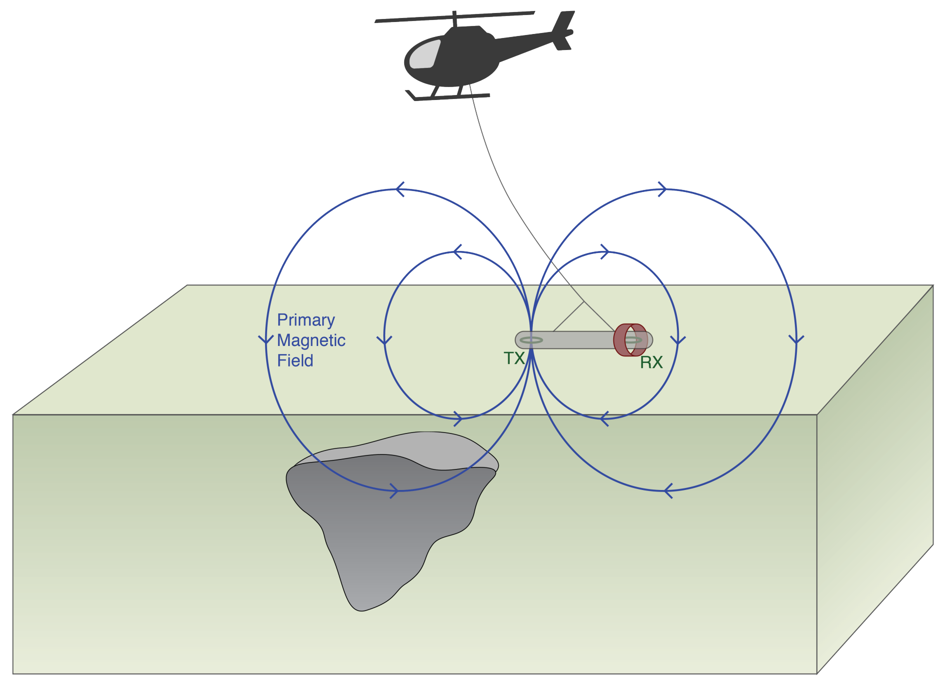
Current loop

$$\mathbf{B} = \frac{\mu_0 m}{4\pi r^3}(2\cos\theta\hat{\mathbf{r}} + \sin\theta\hat{\boldsymbol{\theta}})$$

$$m = IA$$







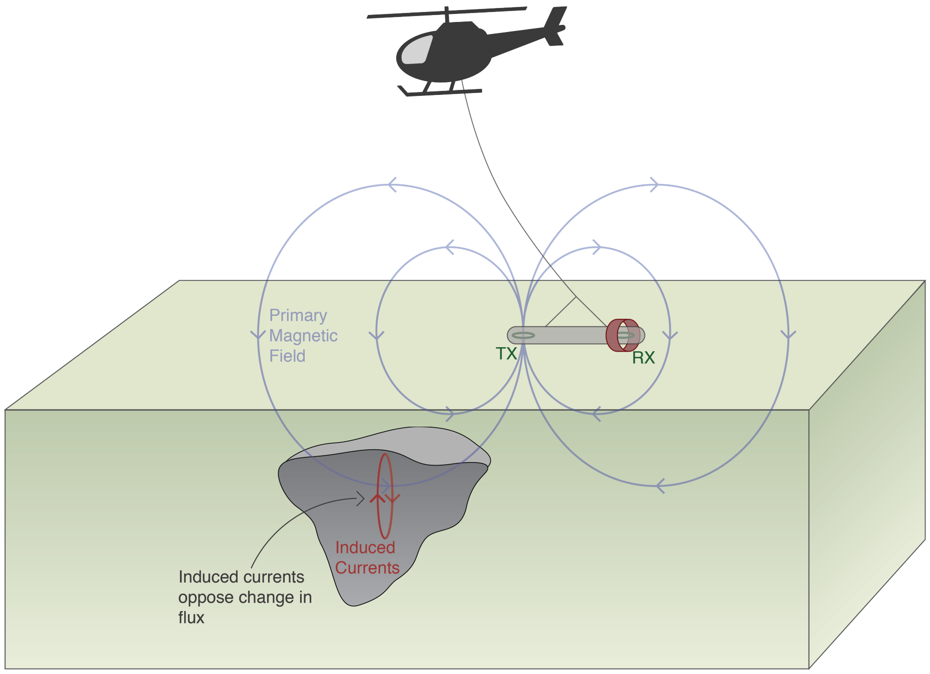
Faraday's Law

Lenz'
Law

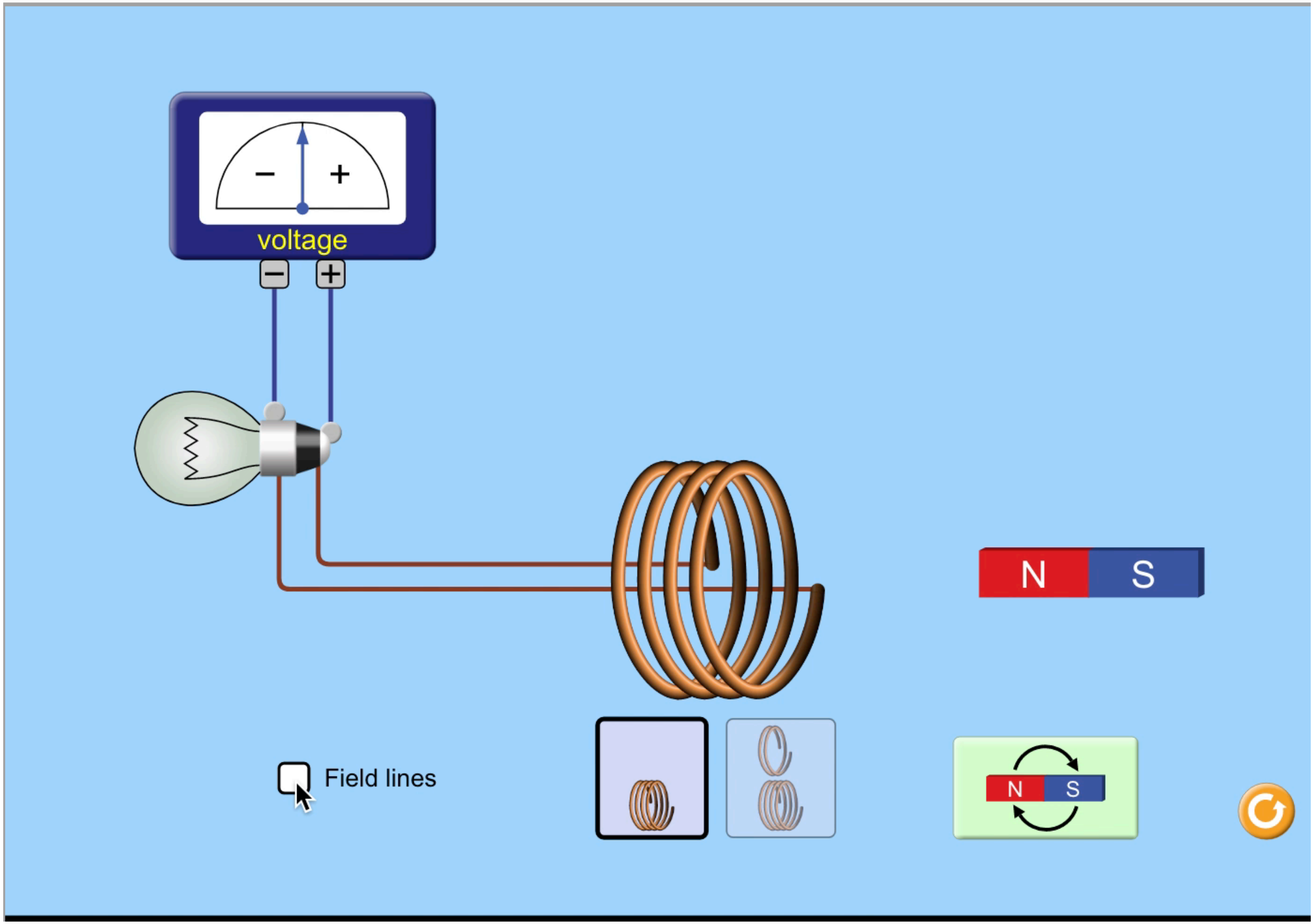
$$\nabla \times \mathbf{e} = -\frac{\partial \mathbf{b}}{\partial t}$$

Ohm's Law

$$\mathbf{j} = \sigma \mathbf{e}$$



Faraday's Law



Faraday's Law

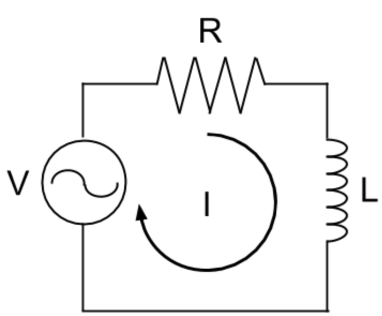
$$\nabla \times \mathbf{e} = -\frac{\partial \mathbf{b}}{\partial t}$$

Magnetic Flux

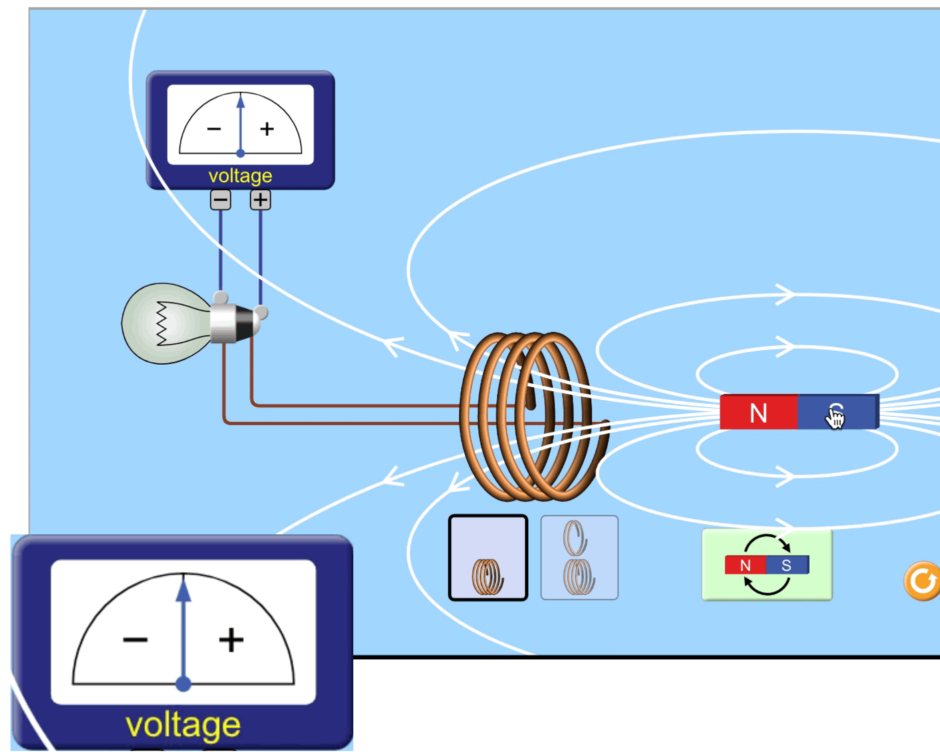
$$\phi_{\mathbf{b}} = \int_A \mathbf{b} \cdot \hat{\mathbf{n}} \; da$$

Induced EMF

$$V = EMF = -\frac{d\phi_{\mathbf{b}}}{dt} \mathbf{= 0}$$



$$\phi_b : \text{constant}$$



Faraday's Law

$$\nabla \times \mathbf{e} = -\frac{\partial \mathbf{b}}{\partial t}$$

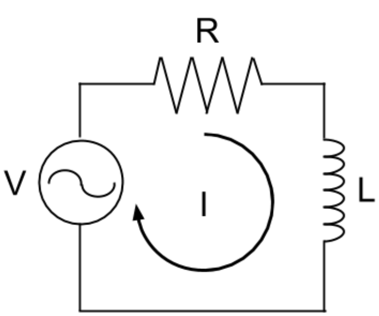
Magnetic Flux

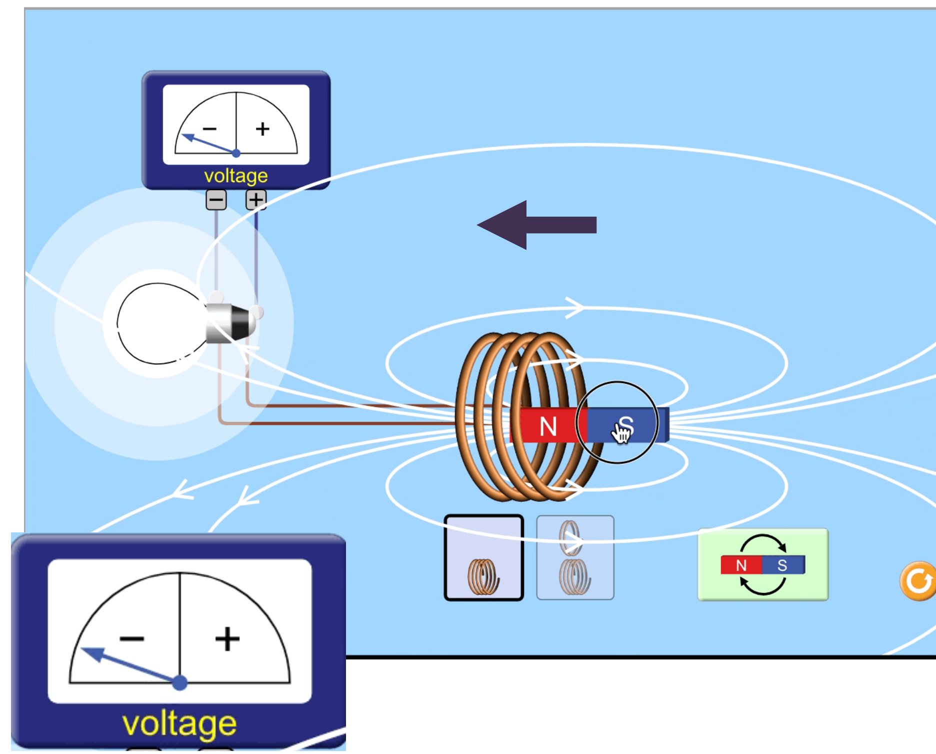
$$\phi_{\mathbf{b}} = \int_A \mathbf{b} \cdot \hat{\mathbf{n}} \; da$$

$$\phi_b :$$

Induced EMF

$$V = EMF = -\frac{d\phi_{\mathbf{b}}}{dt} < 0$$





Faraday's Law

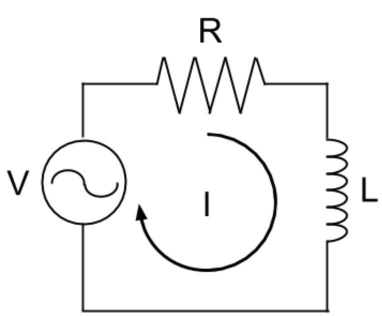
$$\nabla \times \mathbf{e} = -\frac{\partial \mathbf{b}}{\partial t}$$

Magnetic Flux

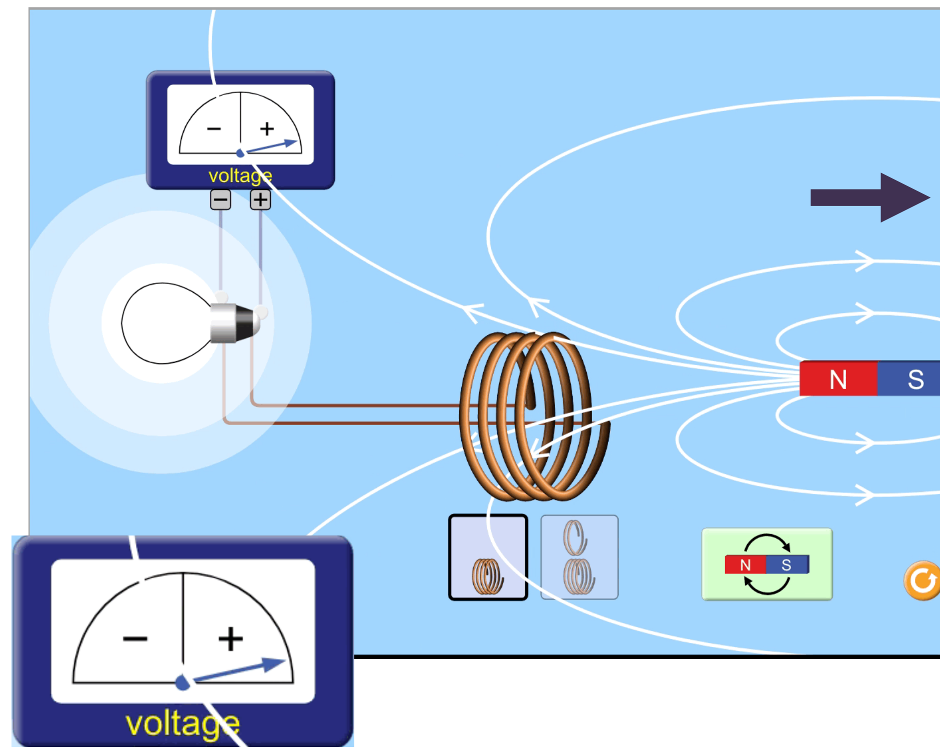
$$\phi_{\mathbf{b}} = \int_A \mathbf{b} \cdot \hat{\mathbf{n}} \; da$$

Induced EMF

$$V = EMF = -\frac{d\phi_{\mathbf{b}}}{dt} > 0$$



ϕ_b :



App for Faraday's Law

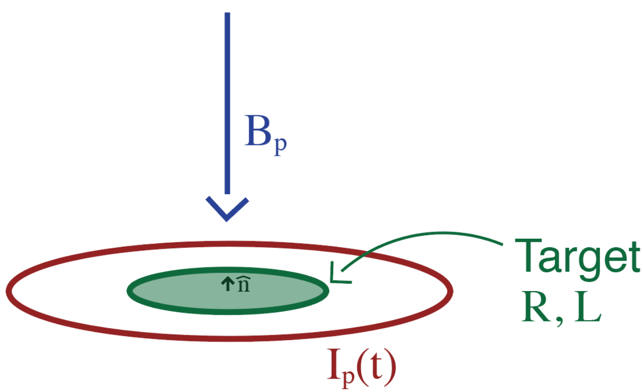
2 Apps:

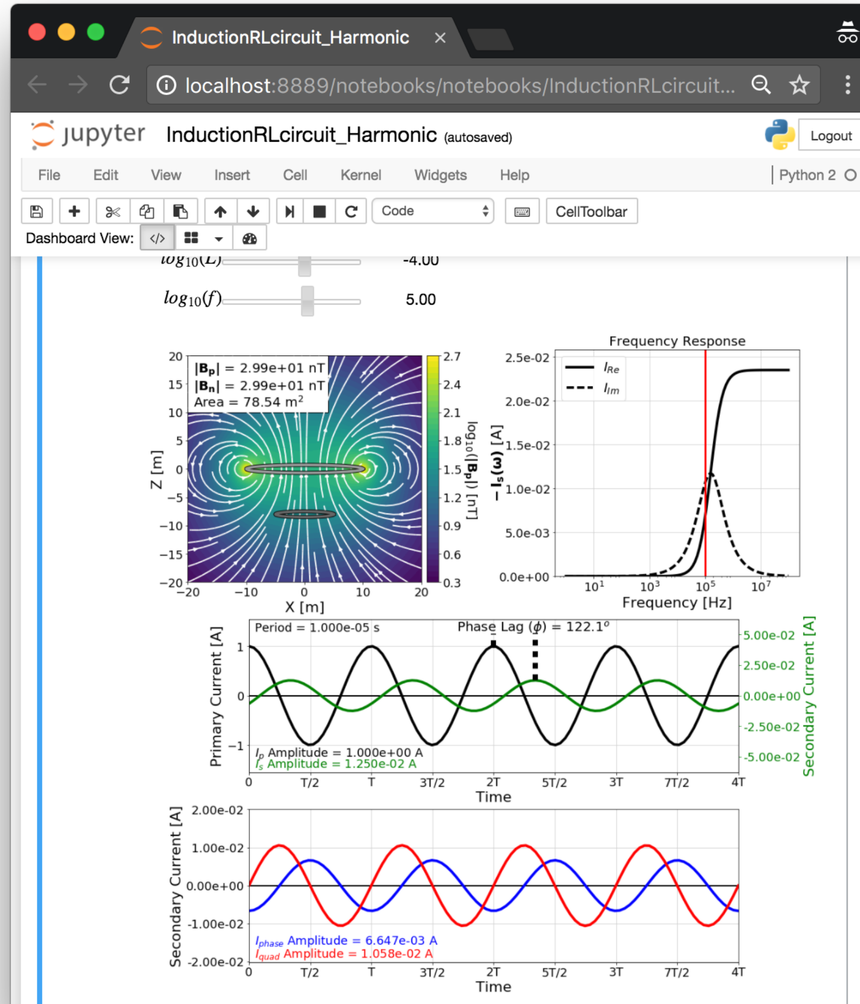
- Harmonic
- Transient

Lenz' Law

$$\nabla \times \mathbf{e} = -\frac{\partial \mathbf{b}}{\partial t}$$

$$\mathbf{j} = \sigma \mathbf{e}$$





http://em.geosci.xyz/apps.html

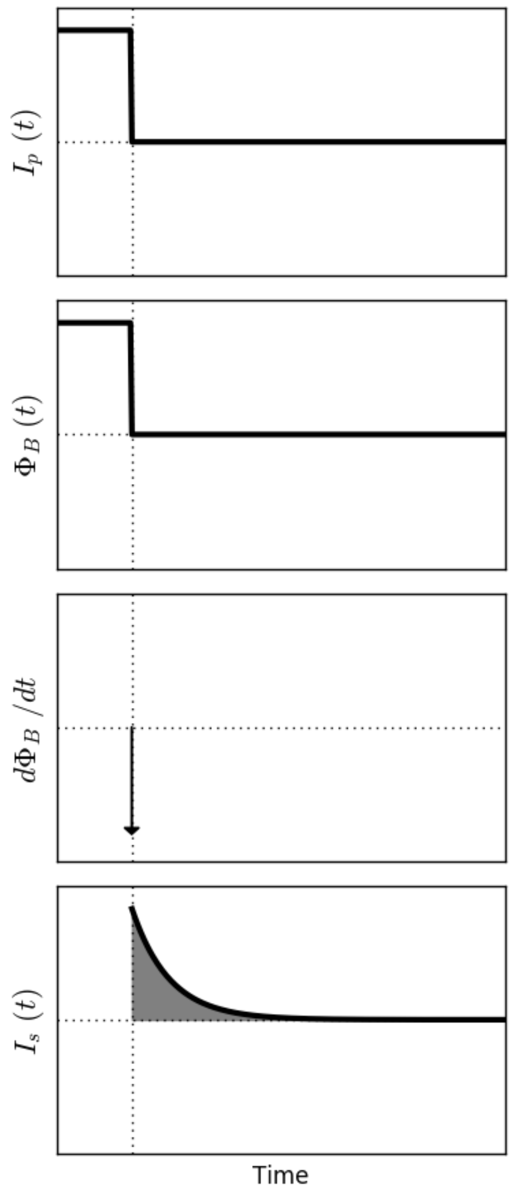
Two Coil Example: Transient

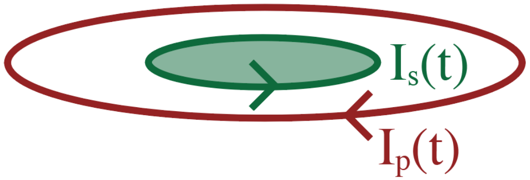
Primary currents

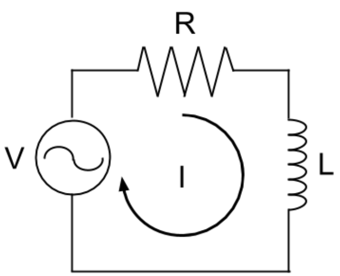
Magnetic flux

Time-variation of
magnetic flux

Secondary currents





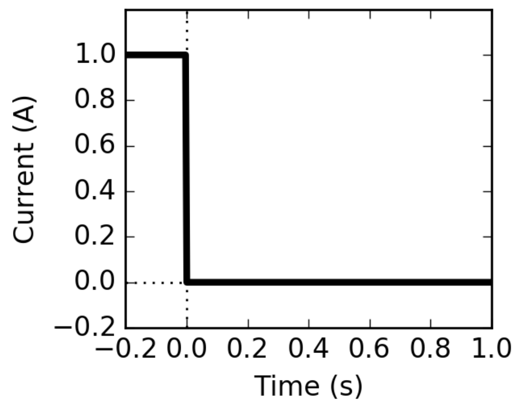


$$I_s(t) = I_s e^{-t/\tau}$$

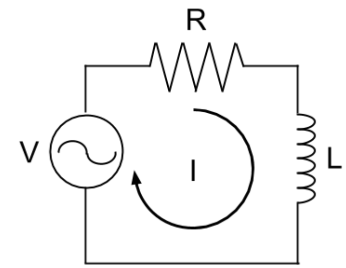
$$\tau = L/R$$

Response Function: Transient

Step-off current in Tx

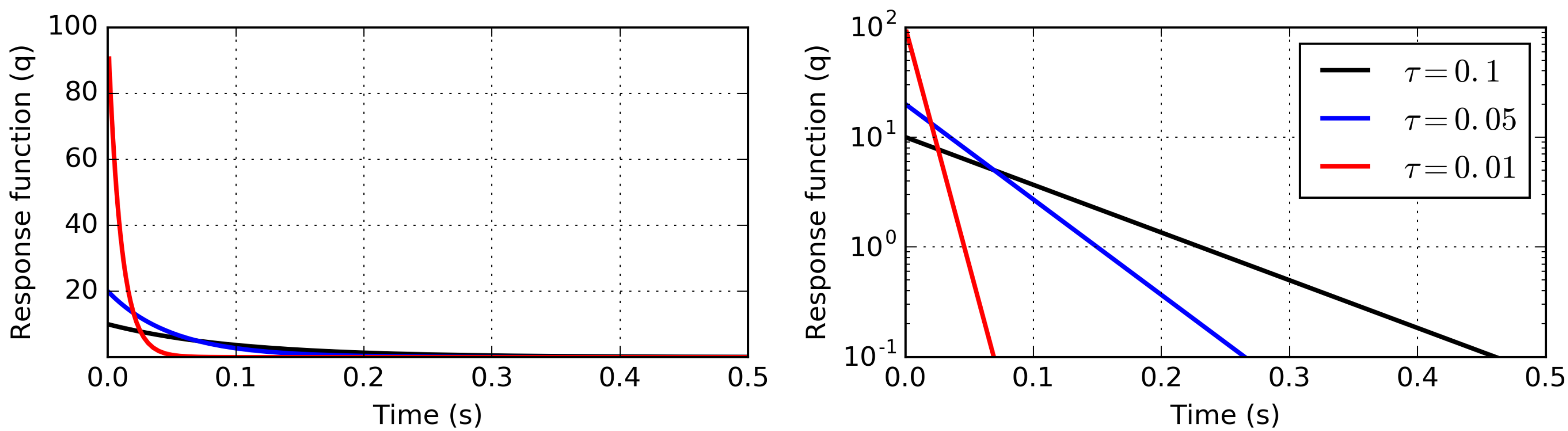


Time constant



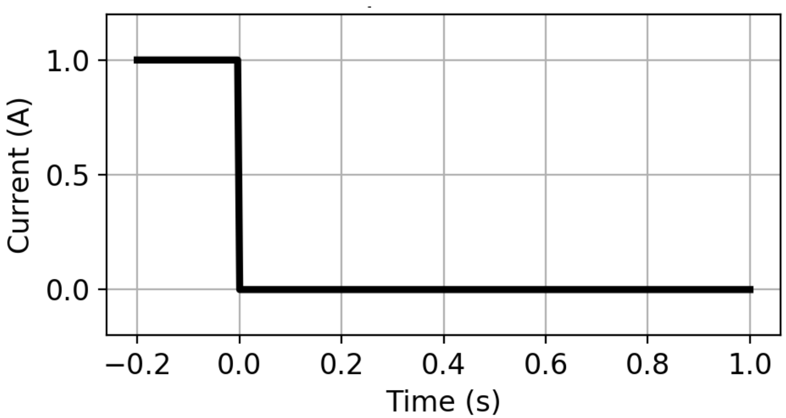
$$\tau = L/R$$

Response function: $q(t) = e^{-t/\tau}$

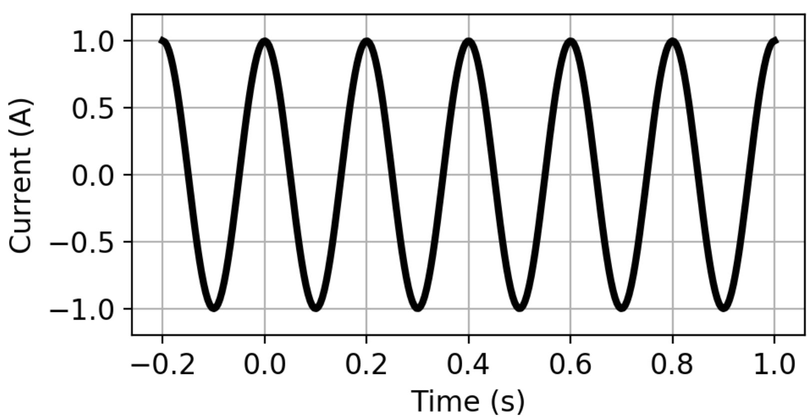


Transient and Harmonic Signals

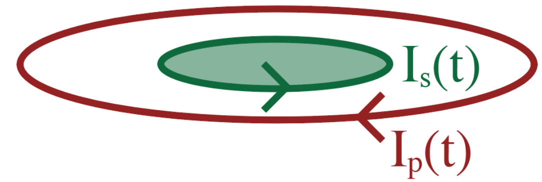
We have seen a transient pulse…



What happens when he have a harmonic?



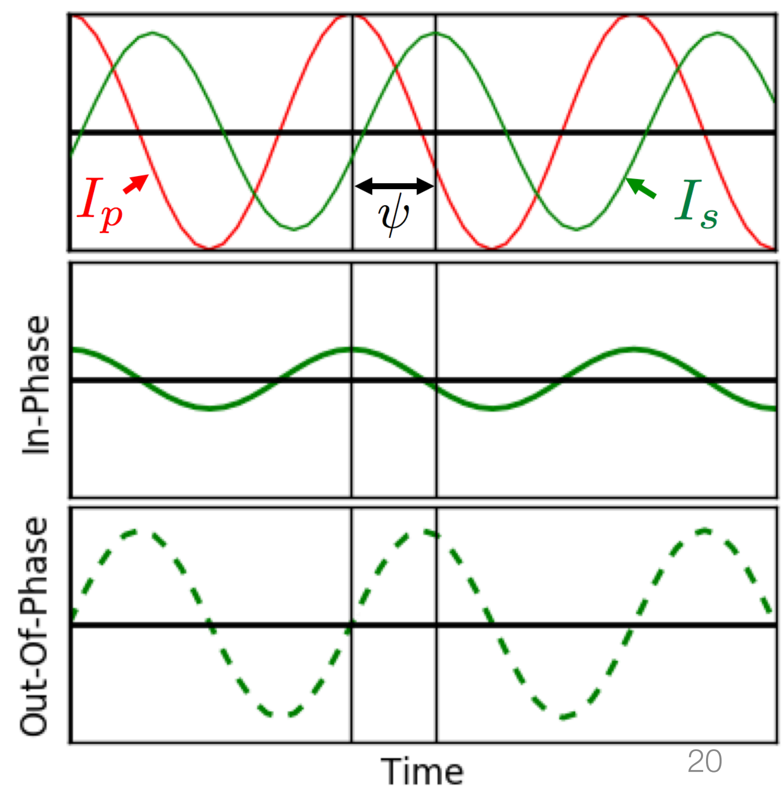
Two Coil Example: Harmonic



Induced Currents

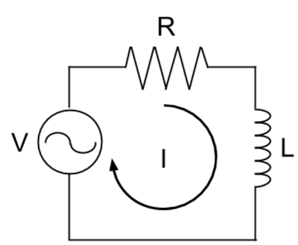
$$I_p(t) = I_p \cos \omega t$$

$$\begin{aligned} I_s(t) &= I_s \cos(\omega t - \psi) \\ &= \underbrace{I_s \cos\psi}_{\substack{\text{In-Phase} \\ \text{Real}}} \cos \omega t + \underbrace{I_s \sin\psi}_{\substack{\text{Out-of-Phase} \\ \text{Quadrature} \\ \text{Imaginary}}} \sin \omega t \end{aligned}$$



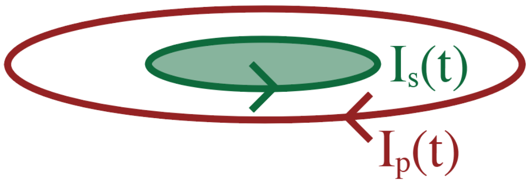
Phase Lag

$$\psi = \frac{\pi}{2} + \tan^{-1}\left(\frac{\omega L}{R}\right)$$



Two Coil Example: Harmonic

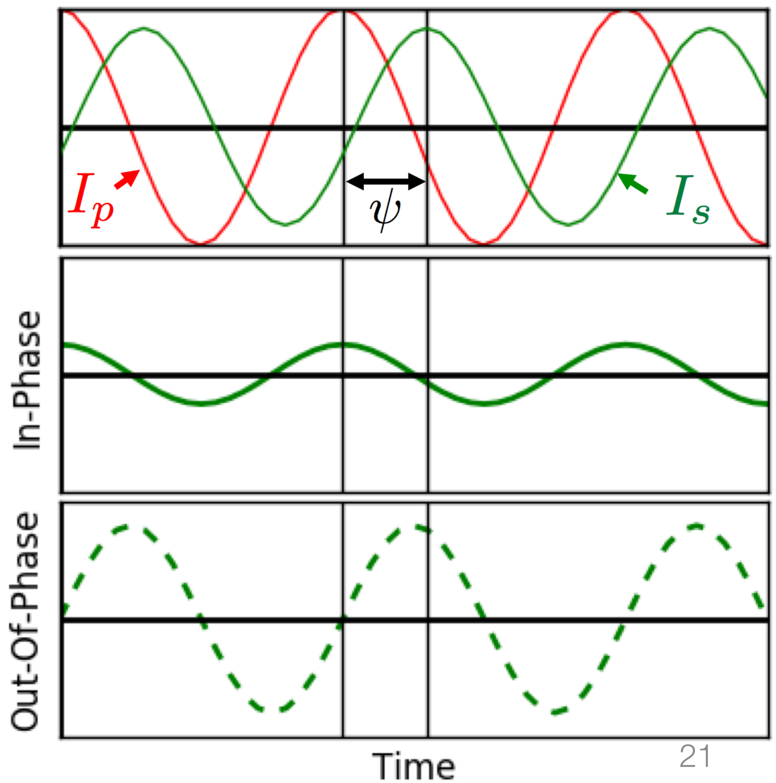
Induced Currents



$$I_p(t) = I_p \cos \omega t$$

$$I_s(t) = I_s \cos(\omega t - \psi)$$

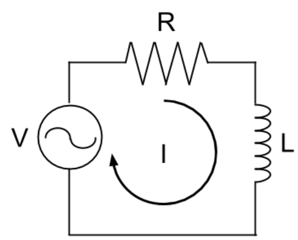
$$= \underbrace{I_s \cos \psi}_{\substack{\text{In-Phase} \\ \text{Real}}} \cos \omega t + \underbrace{I_s \sin \psi}_{\substack{\text{Out-of-Phase} \\ \text{Quadrature} \\ \text{Imaginary}}} \sin \omega t$$



Phase Lag

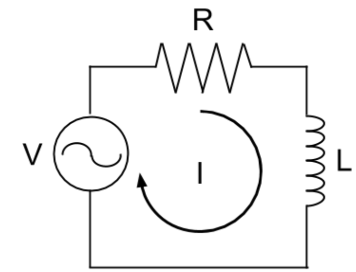
$$\psi = \frac{\pi}{2} + \tan^{-1} \underbrace{\left(\frac{\omega L}{R} \right)}_{\alpha}$$

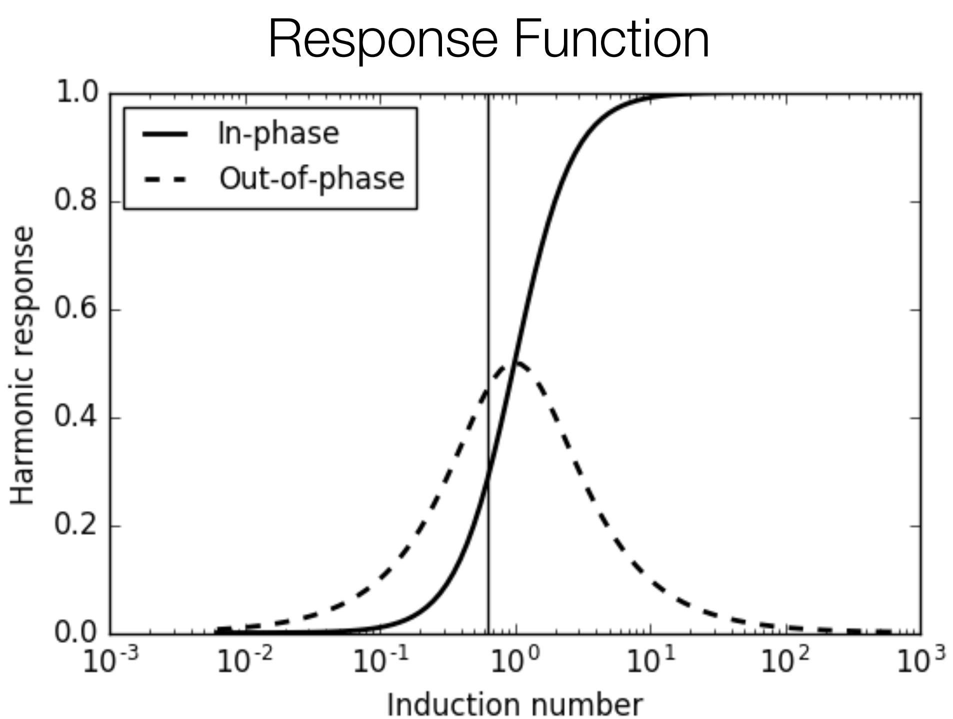
Induction number



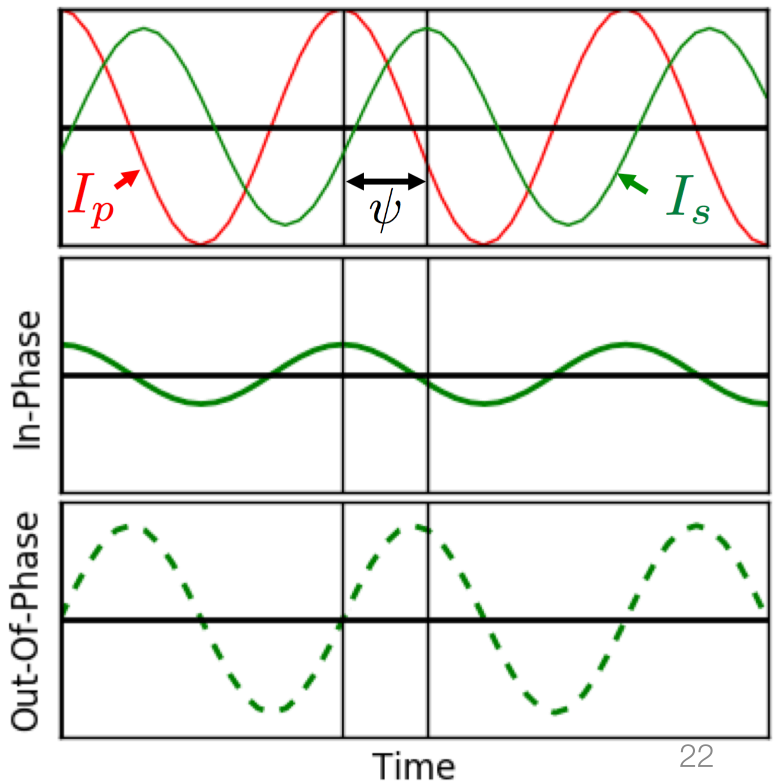
Response Function

- Quantifies how a target responds to a time varying magnetic field
- Partitions real and imaginary parts



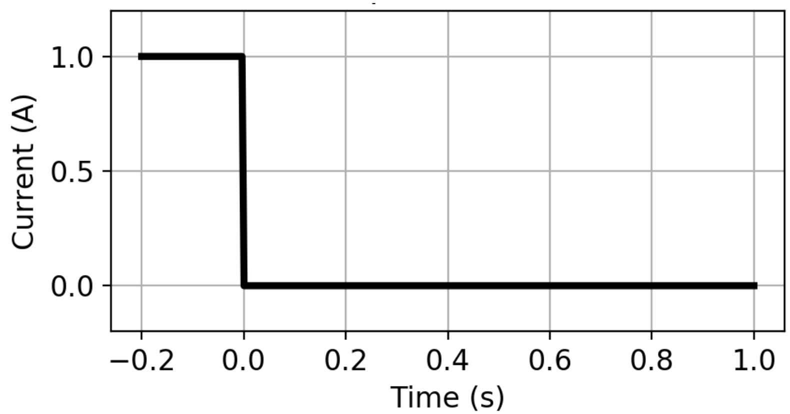


$$\alpha = \frac{\omega L}{R}$$

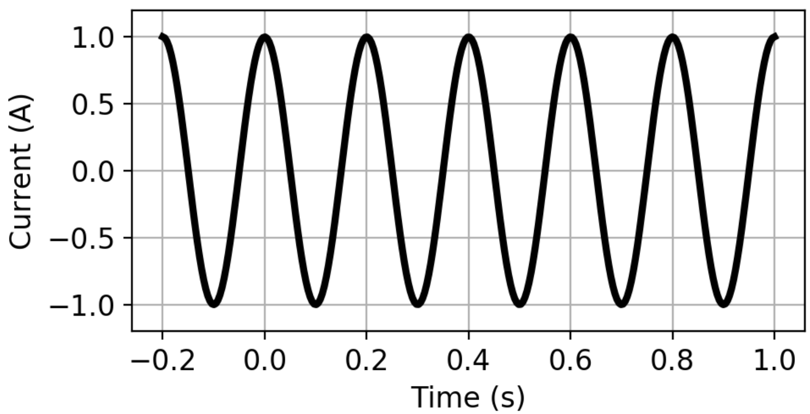


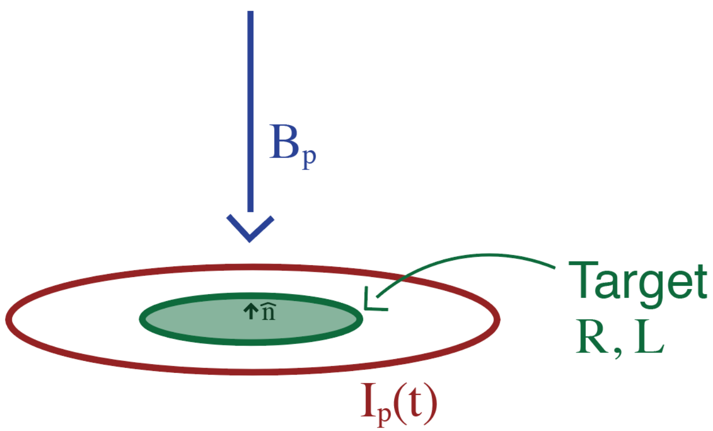
Response Functions: Summary

Step-off



Harmonic





In both:

- Induce currents

$$\nabla \times \mathbf{e} = -\frac{\partial \mathbf{b}}{\partial t}$$

- Generate secondary magnetic fields

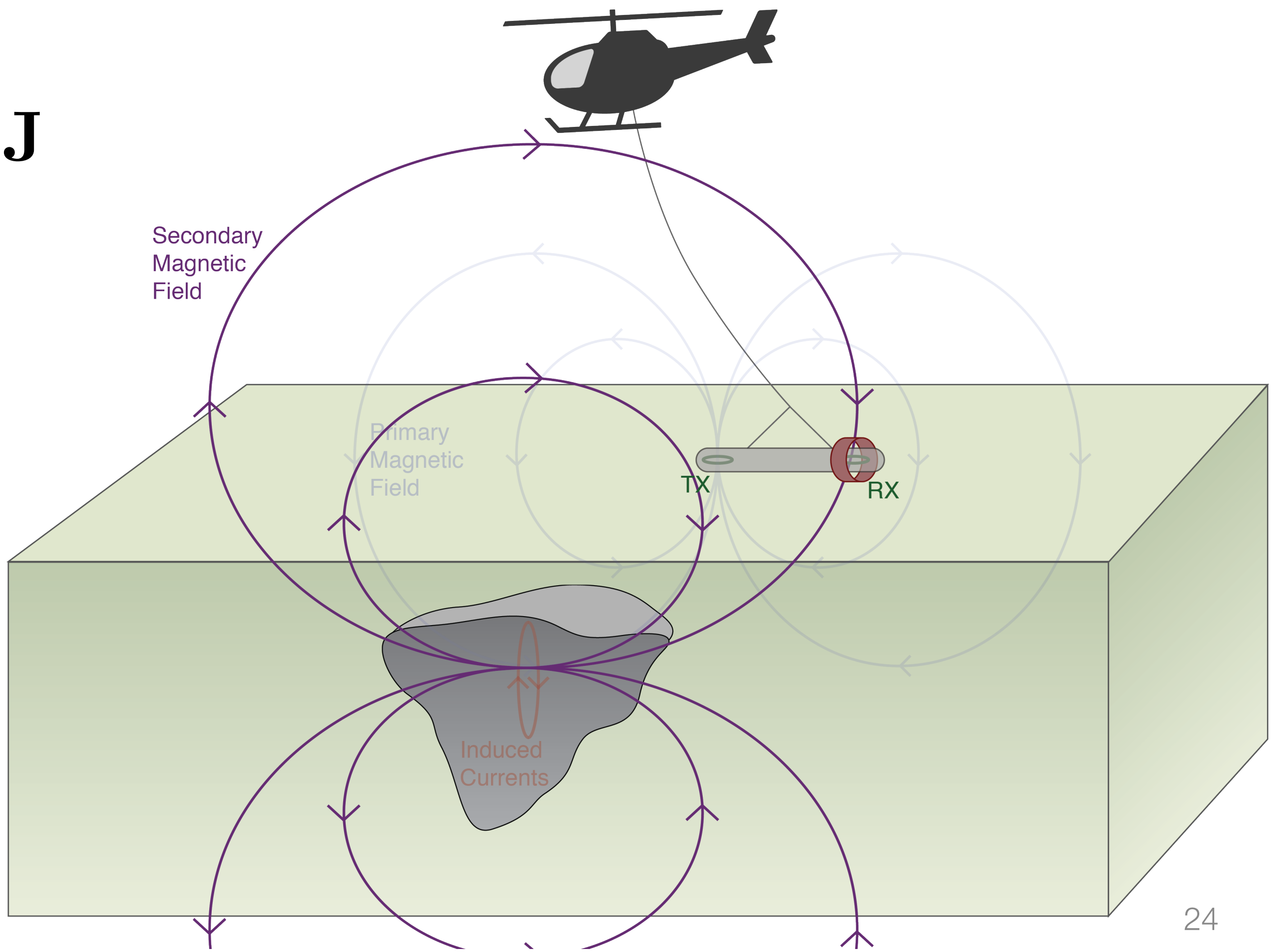
$$\nabla \times \mathbf{h} = \mathbf{j}$$

Secondary magnetic fields

Induced currents generate magnetic fields

- Ampere's Law

$$\nabla \times \mathbf{H} = \mathbf{J}$$

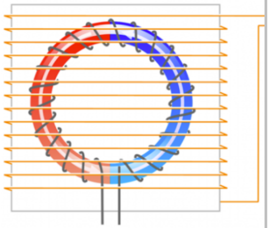


Receiver and Data

Magnetometer

$$\mathbf{b}(t)$$

- Measures:
  - Magnetic fields
  - 3 components
- eg. 3-component fluxgate

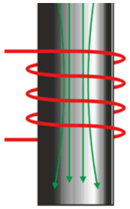


Fluxgate

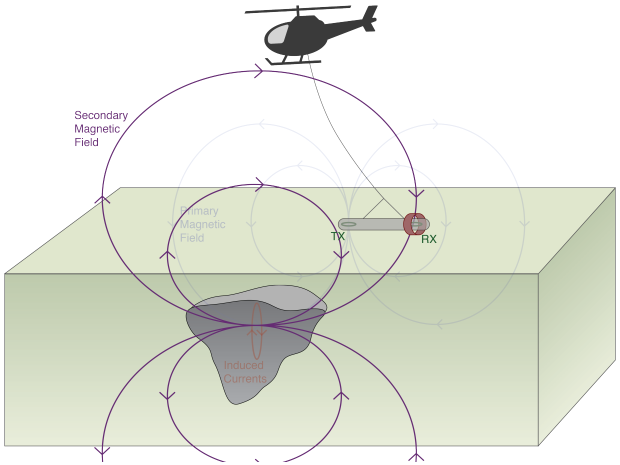
Coil

$$\frac{\partial b}{\partial t}$$

- Measures:
  - Voltage
  - Single component that depends on coil orientation
    - Coupling matters
- eg. airborne frequency domain
  - ratio of Hs/Hp is the same as Vs/Vp



Coil



Coupling

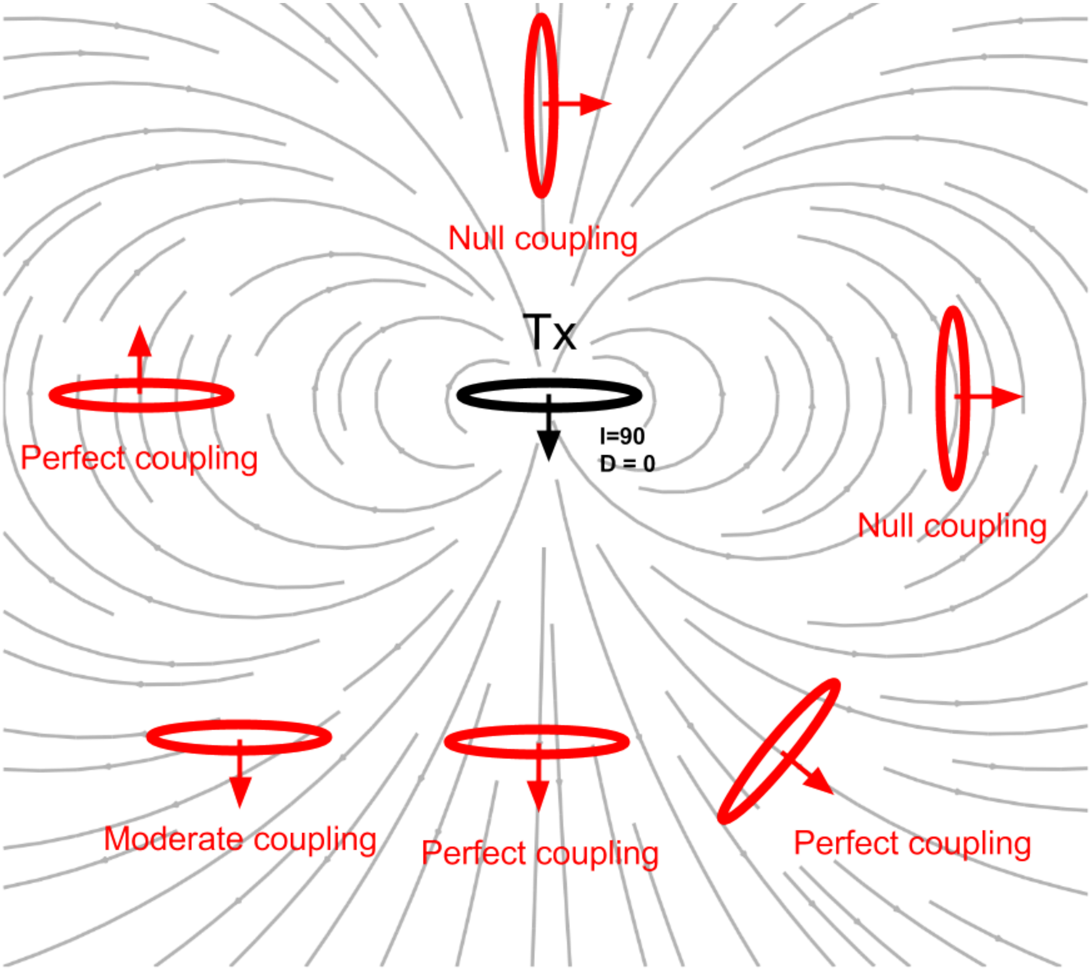
- Transmitter: Primary

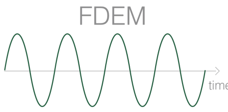
$$I_p(t) = I_p \cos(\omega t)$$

$$\mathbf{B}_p(t) \sim I_p cos(\omega t)$$

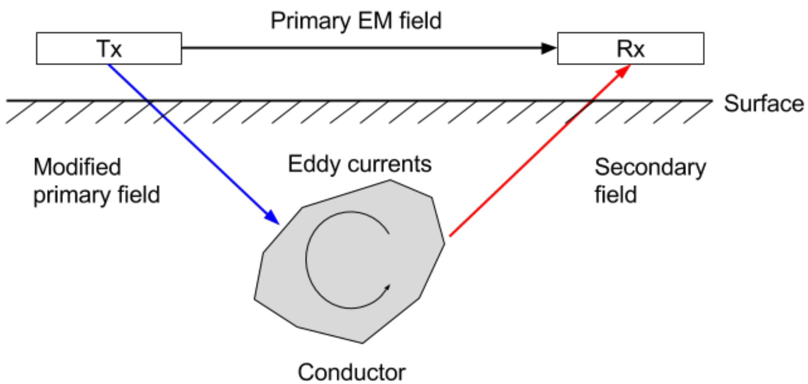
- Target: Secondary

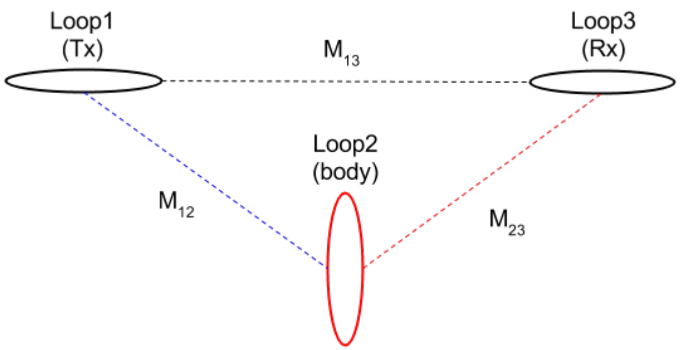
$$\begin{aligned} EMF &= -\frac{\partial \phi_{\mathbf{B}}}{\partial t} \\ &= -\frac{\partial}{\partial t} \left(\mathbf{B}_p \cdot \hat{\mathbf{n}}\right) A \end{aligned}$$





Circuit model of EM induction





Coupling coefficient

- Depends on geometry

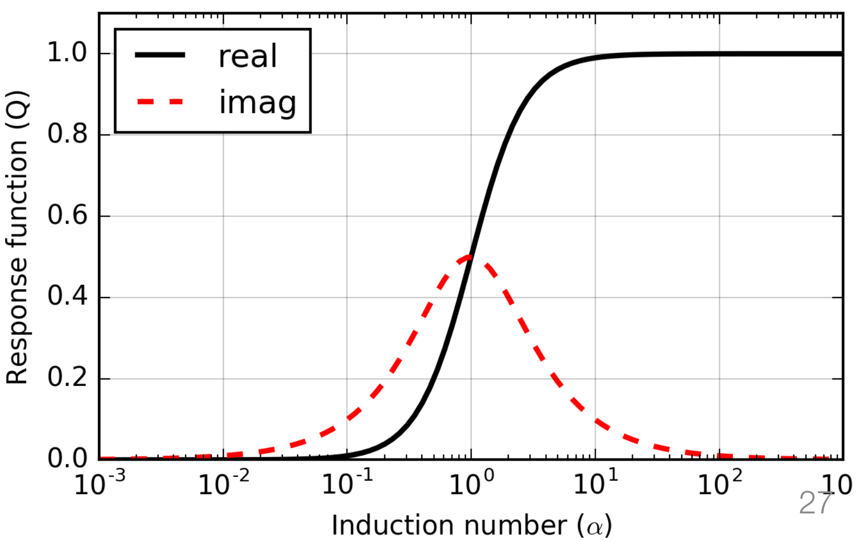
$$M_{12} = \frac{\mu_0}{4\pi} \oint \oint \frac{dl_1 \cdot dl_2}{|\mathbf{r} - \mathbf{r}'|^2}.$$

Magnetic field at the receiver

$$\frac{H^s}{H^p} = -\frac{M_{12} M_{23}}{M_{13} L} \underbrace{\left[\frac{\alpha^2 + \imath\alpha}{1 + \alpha^2} \right]}_{Q}$$

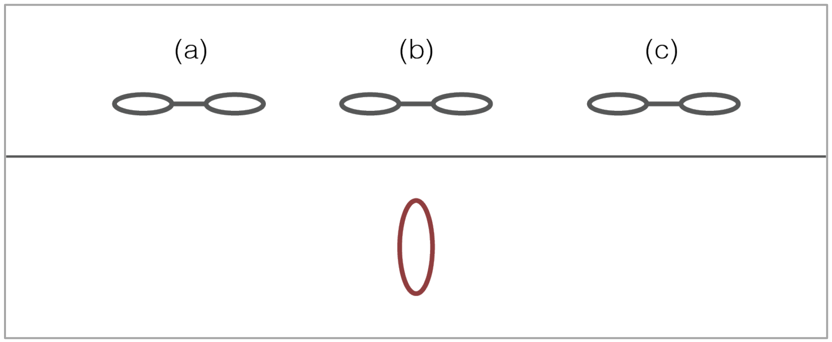
Induction Number

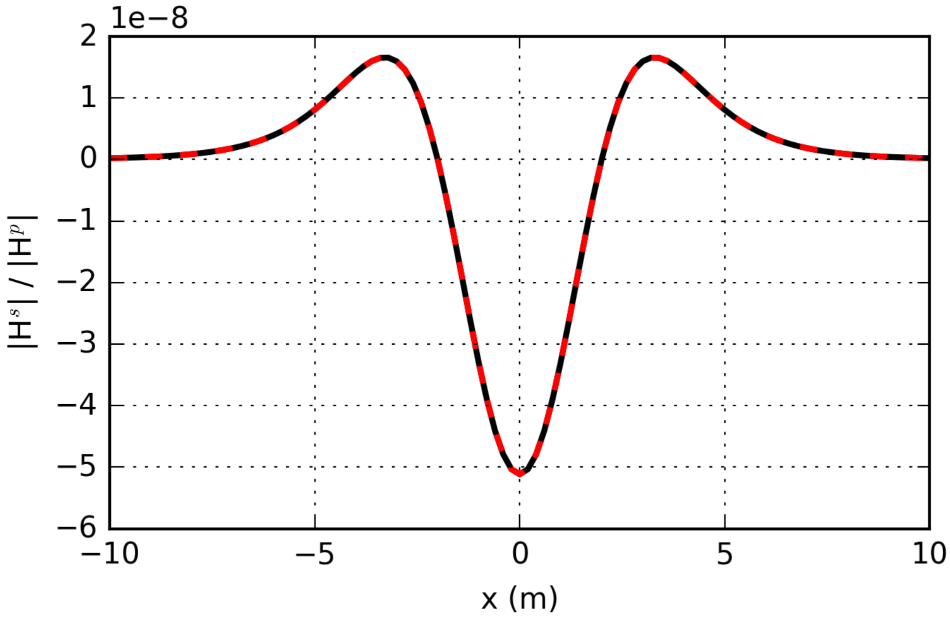
- Depends on properties of target $\alpha = \frac{\omega L}{R}$



Conductor in a resistive earth: Frequency

Profile over the loop



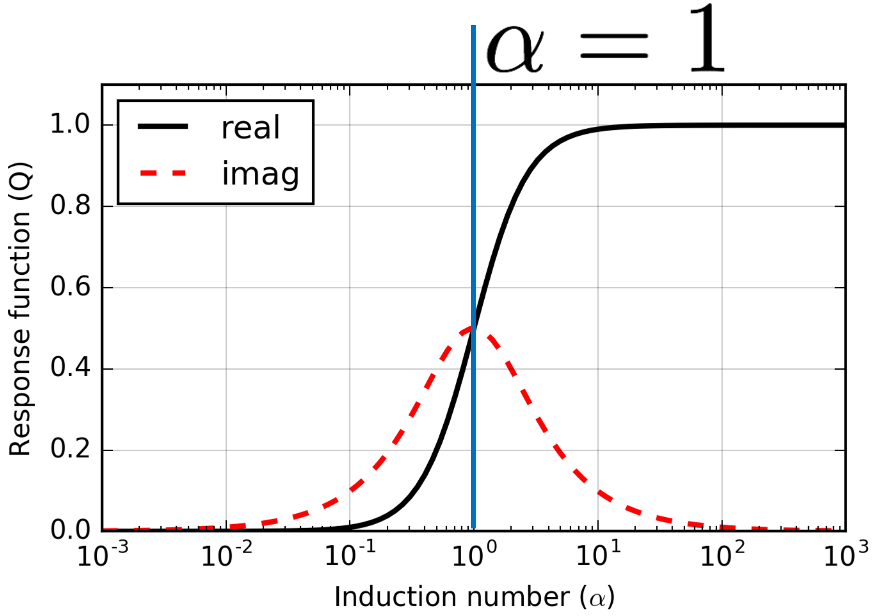


- Induction number

$$\alpha = \frac{\omega L}{R}$$

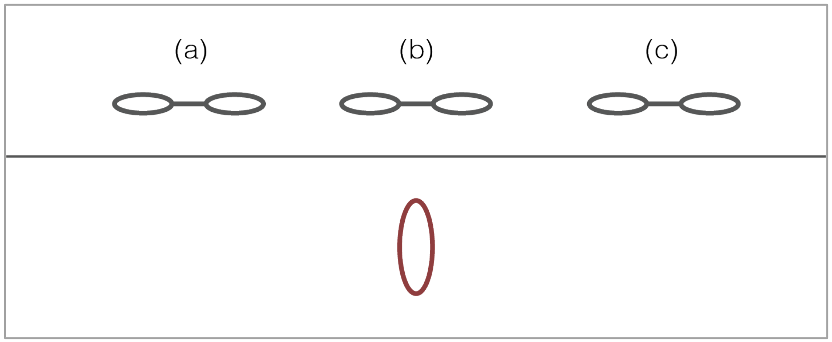
- When $\alpha = 1$
  - Real = Imag

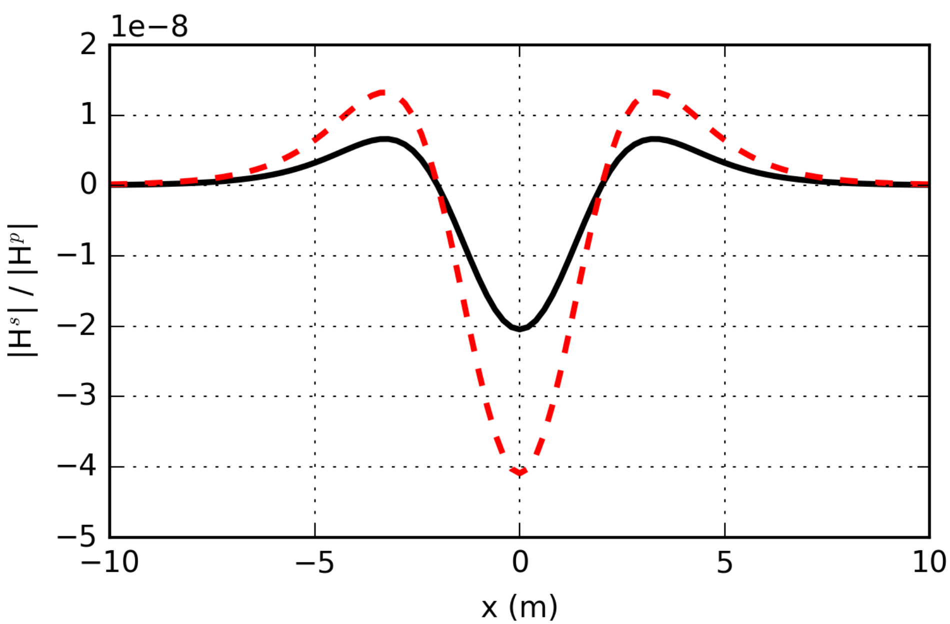
$\alpha = 1$



Conductor in a resistive earth: Frequency

Profile over the loop

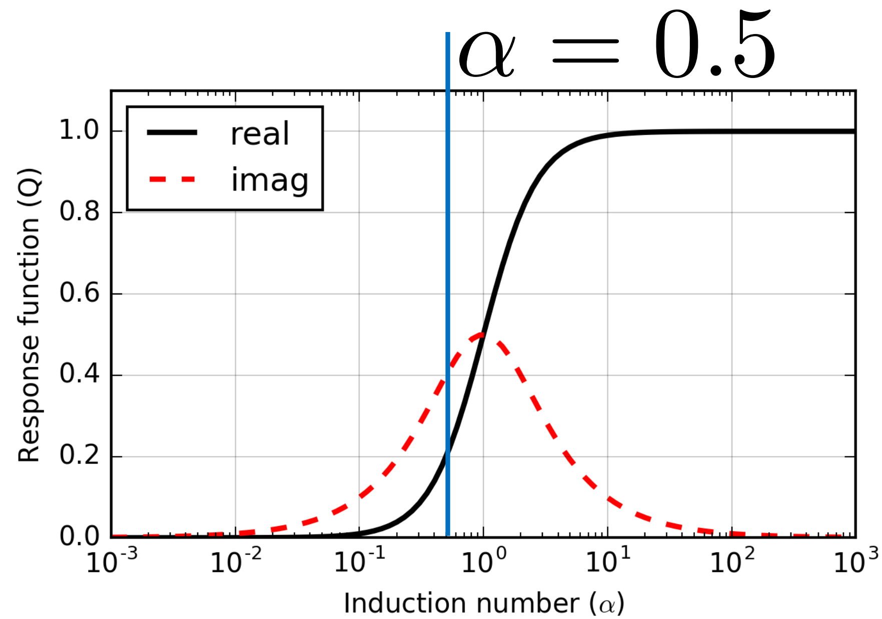




- Induction number

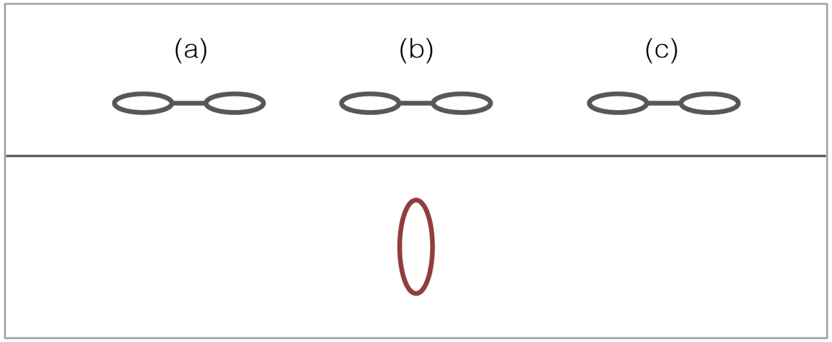
$$\alpha = \frac{\omega L}{R}$$

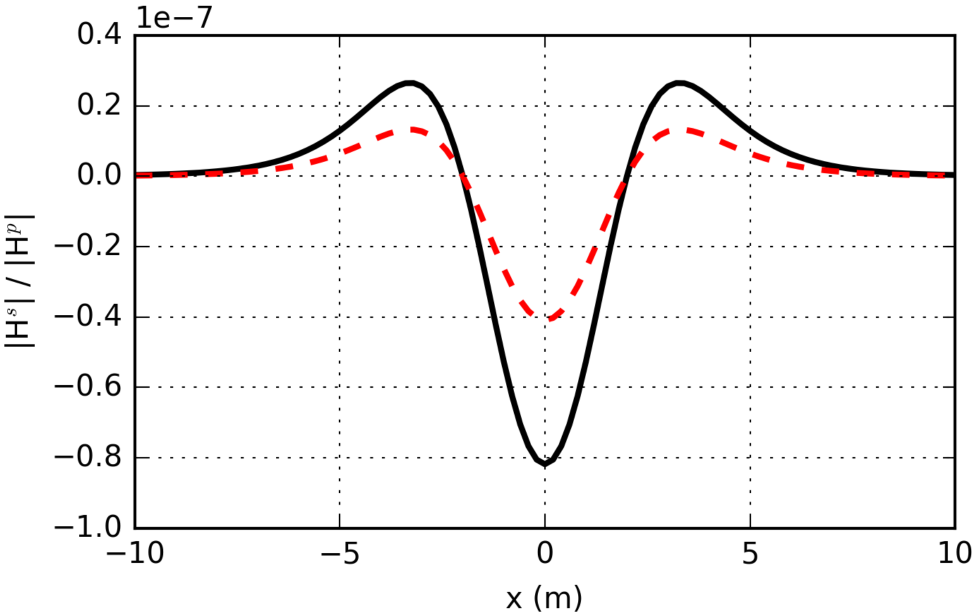
- When $\alpha < 1$
  - Real < Imag



Conductor in a resistive earth: Frequency

Profile over the loop

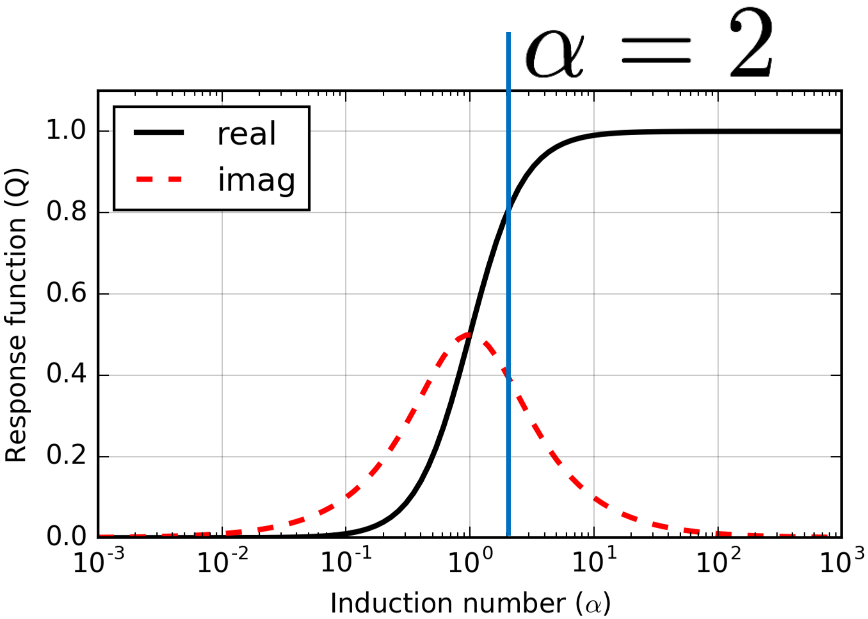




- Induction number

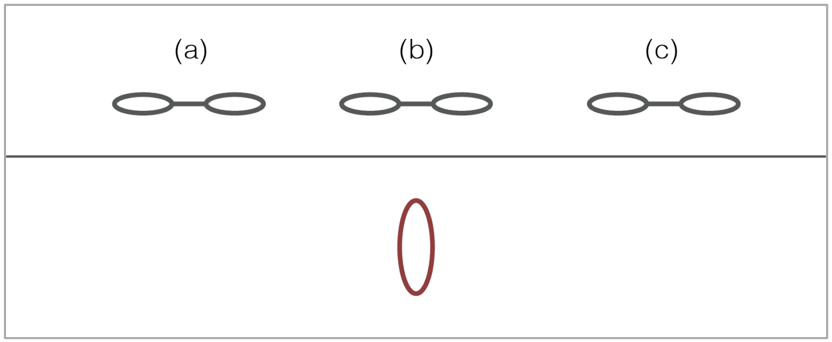
$$\alpha = \frac{\omega L}{R}$$

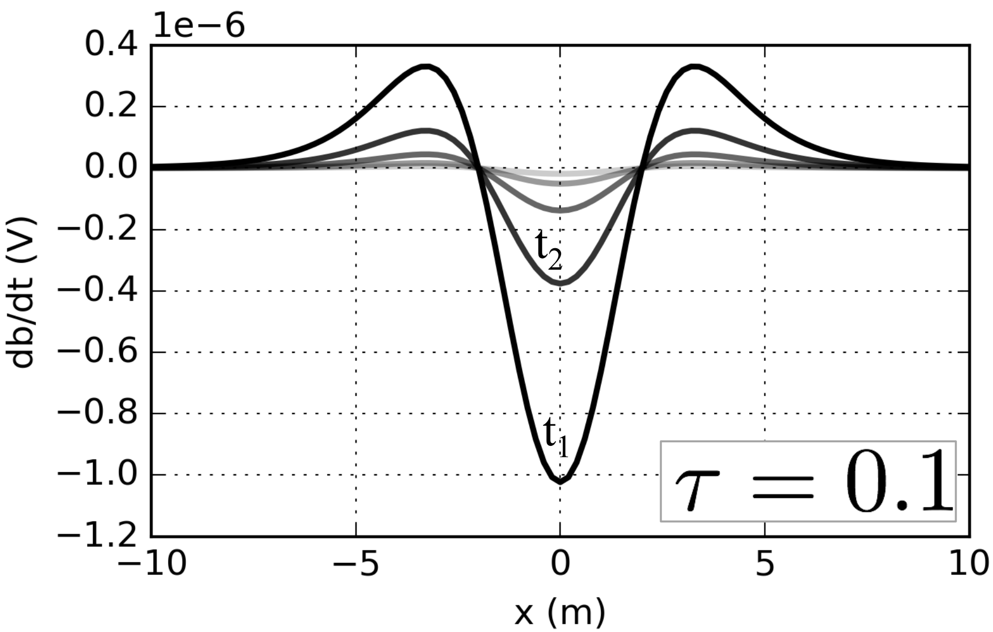
- When $\alpha > 1$
  - Real > Imag



Conductor in a resistive earth: Transient

Profile over the loop

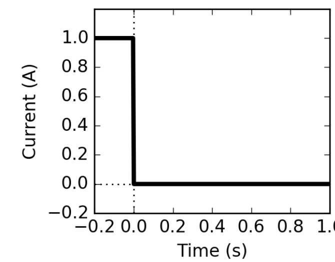




- Time constant

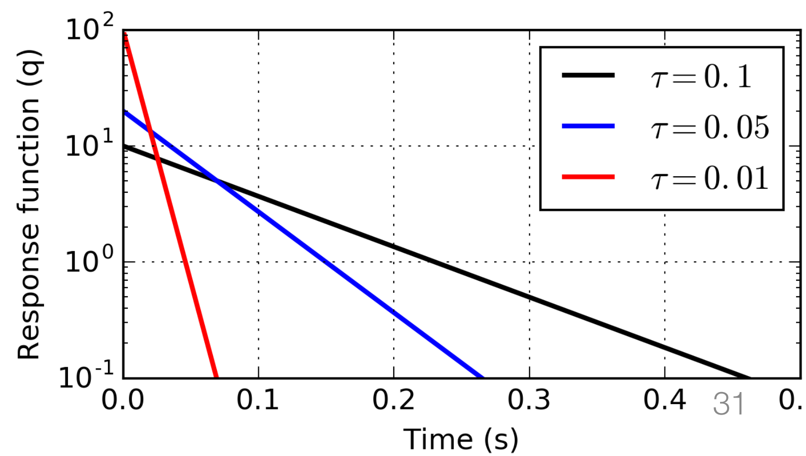
$$\tau = L/R$$

- Step-off current in Tx



- Response function depends on time, τ

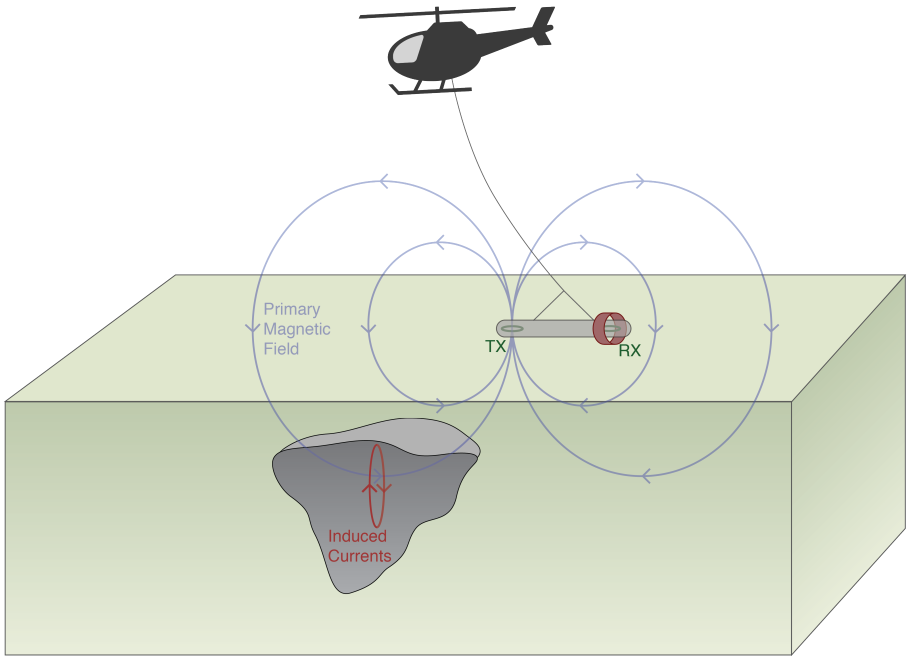
$$q(t) = e^{-t/\tau}$$



Recap: what have we learned?

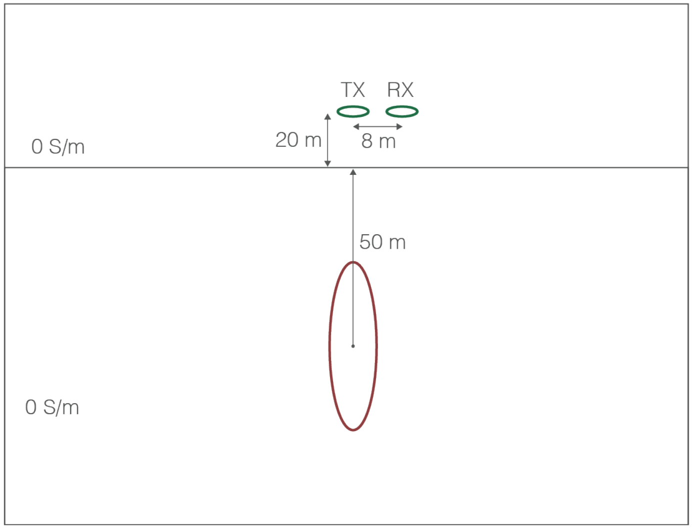
- Basics of EM induction
- Response functions
- Mutual coupling
- Data for frequency or time domain systems

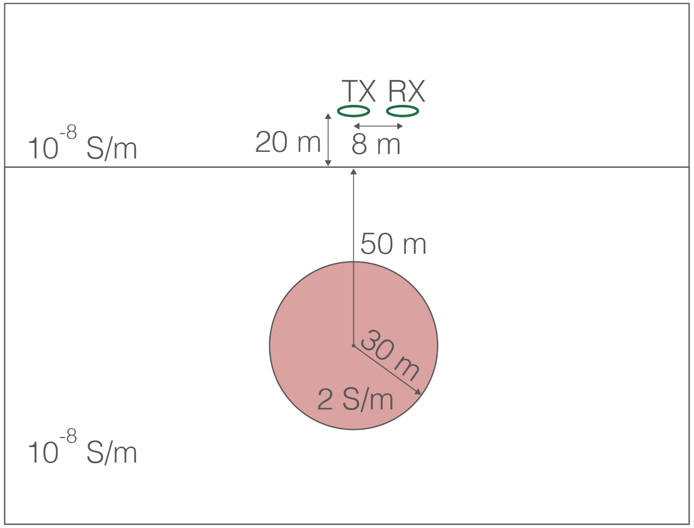
- Circuit model provides representative results
  - Applicable to geologic targets?

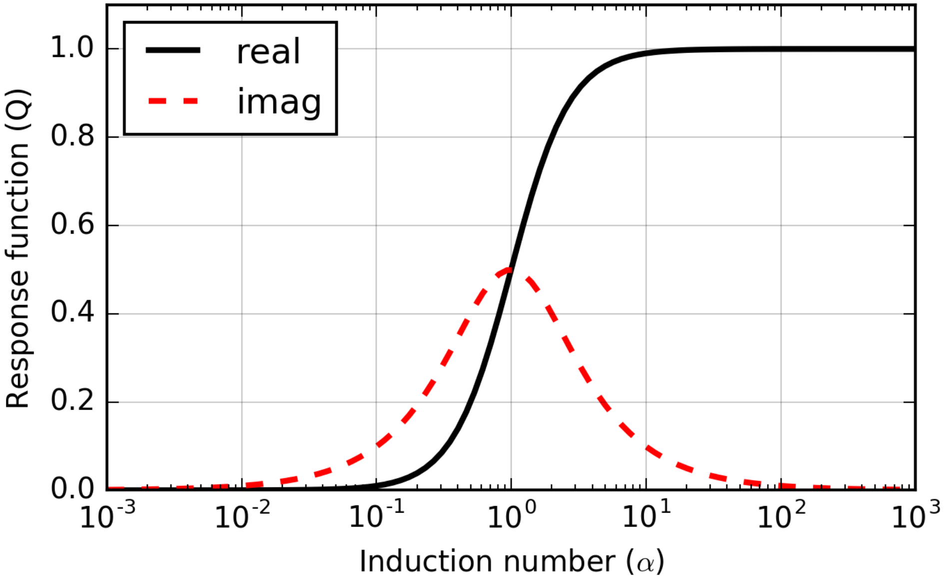


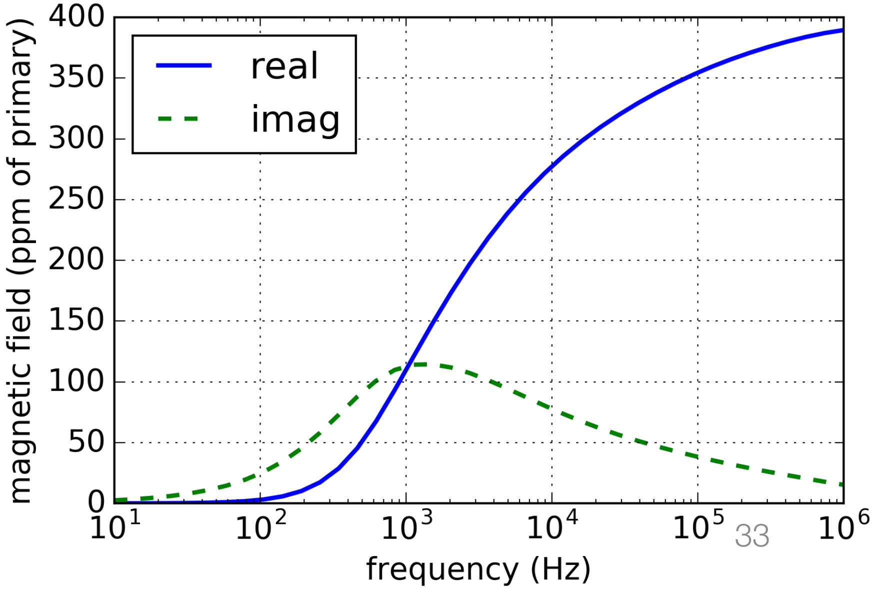
Sphere in a resistive background

How representative is a circuit model?



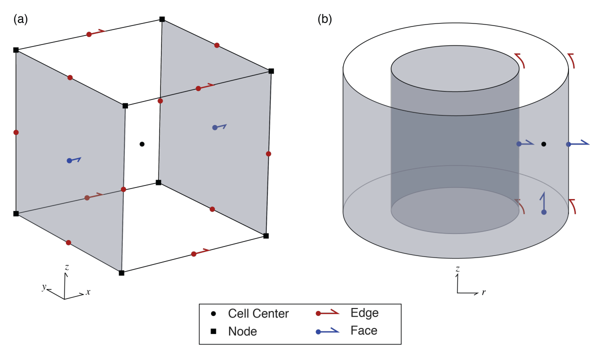






Cyl Code

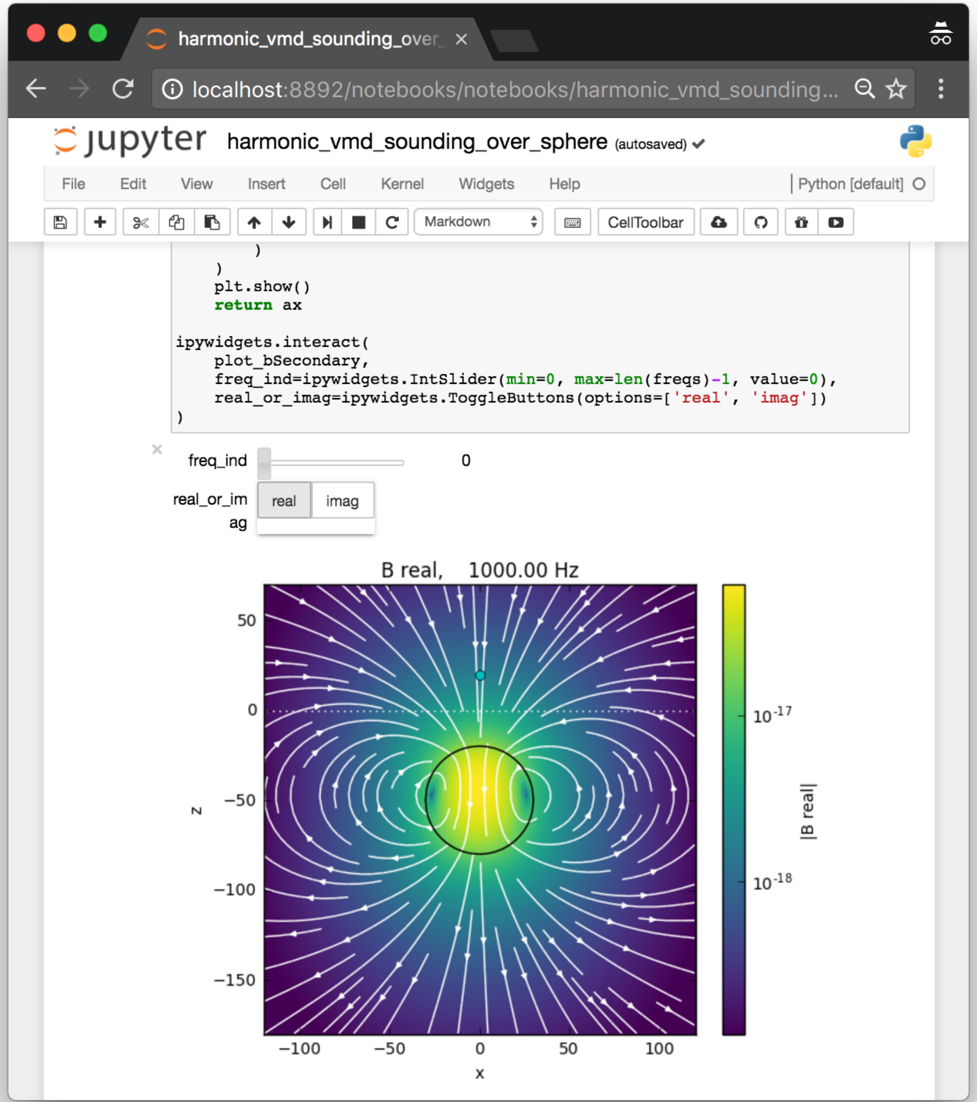
- Finite Volume EM
  - Frequency and Time



- Built on SimPEG

- Open source, available at:
http://em.geosci.xyz/apps.html

- Papers
  - [Cockett et al, 2015]()
  - [Heagy et al, 2017]()



Recap: what have we learned?

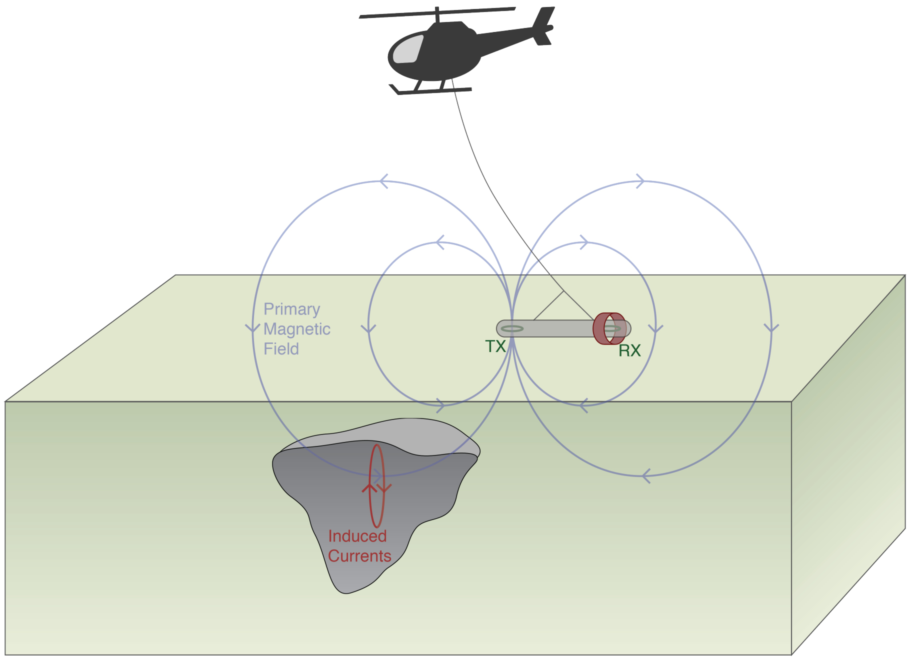
- Basics of EM induction
- Response functions
- Mutual coupling
- Data for frequency or time domain systems
- Circuit model is a good proxy

2-Coil Apps

- Frequency domain
- Time domain

Major item not yet accounted for…

- Propagation of energy from
  - Transmitter to target
  - Target to receiver



How do EM fields and fluxes behave in a conductive background?

Revisit Maxwell's equations

First order equations

$$\nabla \times \mathbf{e} = -\frac{\partial \mathbf{b}}{\partial t}$$

$$\nabla \times \mathbf{h} = \mathbf{j} + \frac{\partial \mathbf{d}}{\partial t}$$

$$\mathbf{j} = \sigma \mathbf{e}$$

$$\mathbf{b} = \mu \mathbf{h}$$

$$\mathbf{d} = \varepsilon \mathbf{e}$$

Second order equations

$$\nabla^2 \mathbf{h} - \underbrace{\mu\sigma \frac{\partial \mathbf{h}}{\partial t}}_{\text{diffusion}} - \underbrace{\mu\epsilon \frac{\partial^2 \mathbf{h}}{\partial t^2}}_{\text{wave propagation}} = 0$$

In frequency

$$\nabla^2 \mathbf{H} + k^2 \mathbf{H} = 0$$

$$k^2 = \omega^2 \mu\varepsilon - i\omega\mu\sigma$$

* Same equation holds for E

Plane waves in a homogeneous media

In frequency

$$\nabla^2 \mathbf{H} + k^2 \mathbf{H} = 0$$

$$k^2 = \omega^2 \mu \varepsilon - i\omega\mu\sigma$$

Quasi-static

$$\frac{\omega \varepsilon}{\sigma} \ll 1$$

even if…

$$\sigma = 10^{-4} S/m$$

$$f = 10^4 Hz$$

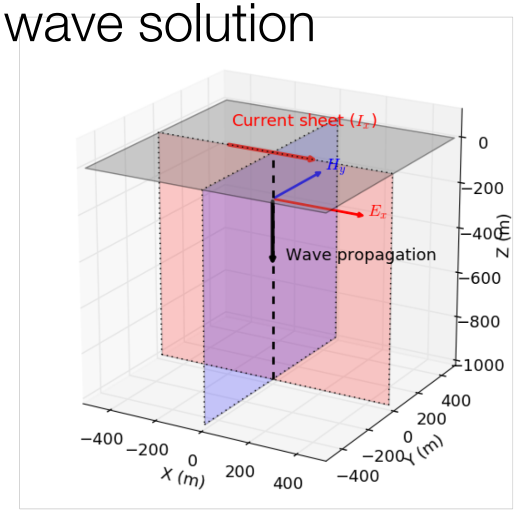
then

$$\frac{\omega \varepsilon}{\sigma} \sim 0.005$$

$$k = \sqrt{-i\omega\mu\sigma} = (1-i)\sqrt{\frac{\omega\mu\sigma}{2}}$$

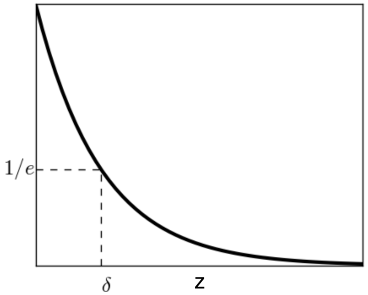
$$\equiv \alpha - i\beta$$

Plane wave solution



$$\mathbf{H} = \mathbf{H_0} \underbrace{e^{-\alpha z}}_{\text{attenuation}} \underbrace{e^{-i(\beta z - \omega t)}}_{\text{phase}}$$

Skin depth



$$\delta : \text{skin depth}$$

$$\delta = \sqrt{\frac{2}{\omega\mu\sigma}} = 503\sqrt{\frac{1}{\sigma f}}$$

Plane waves in a homogeneous media

In time

$$\nabla^2 \mathbf{h} - \mu\epsilon \frac{\partial^2 \mathbf{h}}{\partial t^2} - \mu\sigma \frac{\partial \mathbf{h}}{\partial t} = 0$$

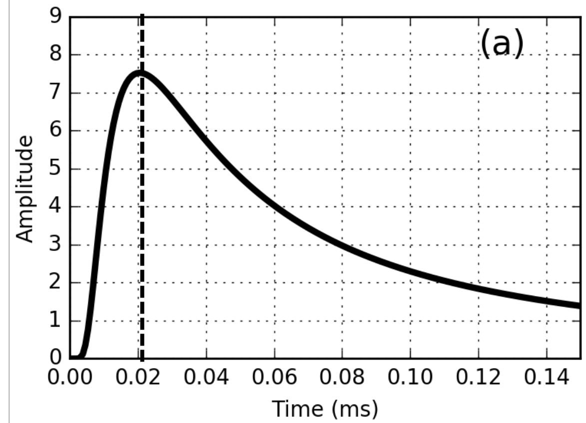
$$\mathbf{h}(t=0) = \mathbf{h}_0 \delta(t)$$

Solution for quasi-static

$$\mathbf{h}(t) = -\frac{(\mu\sigma)^{1/2} z}{2\pi^{1/2} t^{3/2}} e^{-\mu\sigma z^2/(4t)}$$

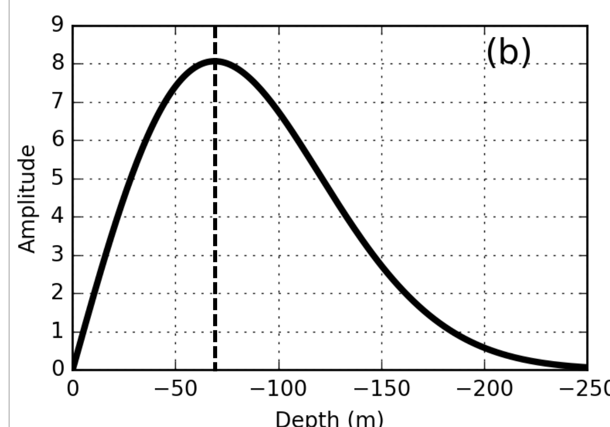
z: depth (m)

Peak time:



$$t_{max} = \frac{\mu\sigma z^2}{6}$$

Diffusion distance



$$d = \sqrt{\frac{2t}{\mu\sigma}}$$

$$\approx 1260\sqrt{\frac{t}{\sigma}}$$

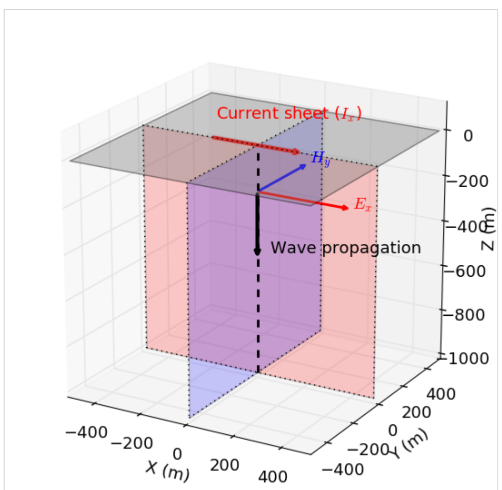
Plane Wave apps

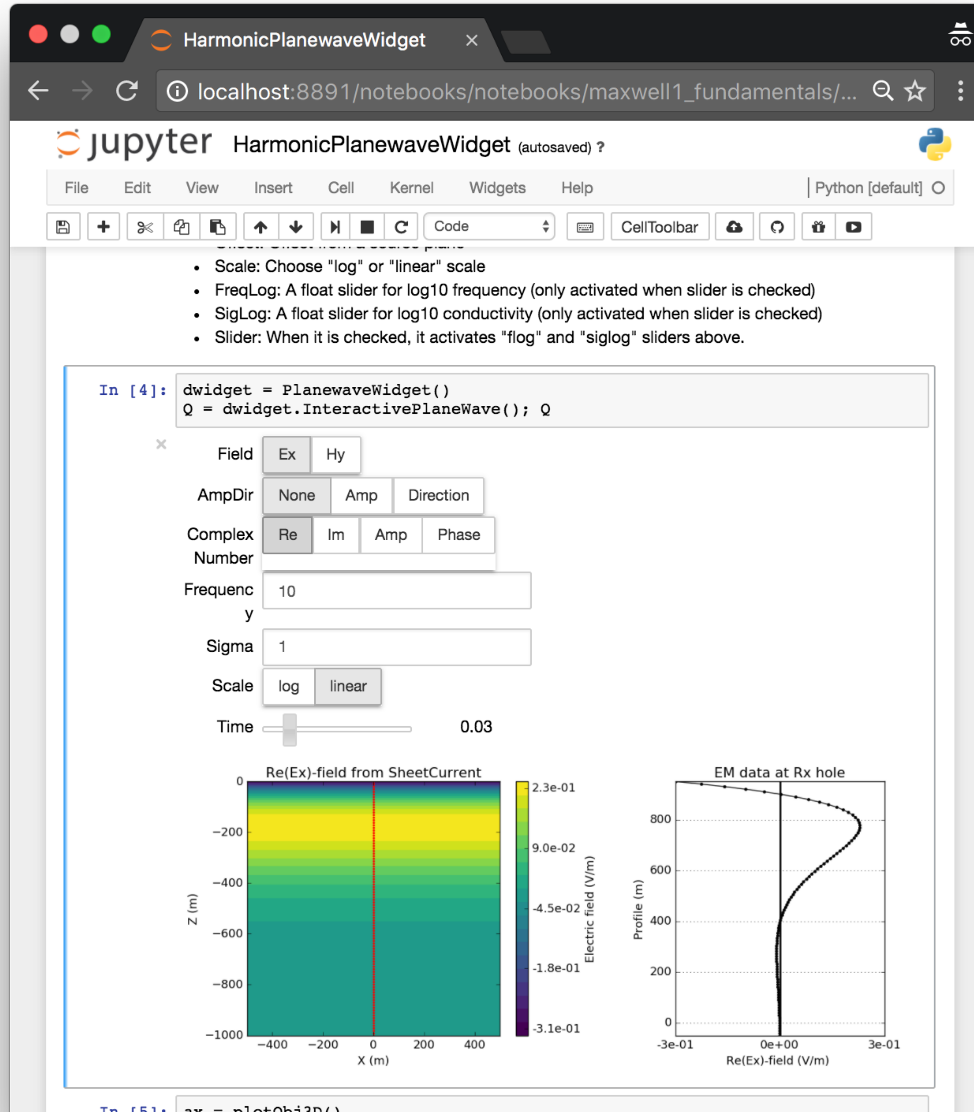
- 2 apps:
  - Transient

$$\mathbf{h}(t) = -\frac{(\mu\sigma)^{1/2} z}{2\pi^{1/2} t^{3/2}} e^{-\mu\sigma z^2/(4t)}$$

  - Harmonic

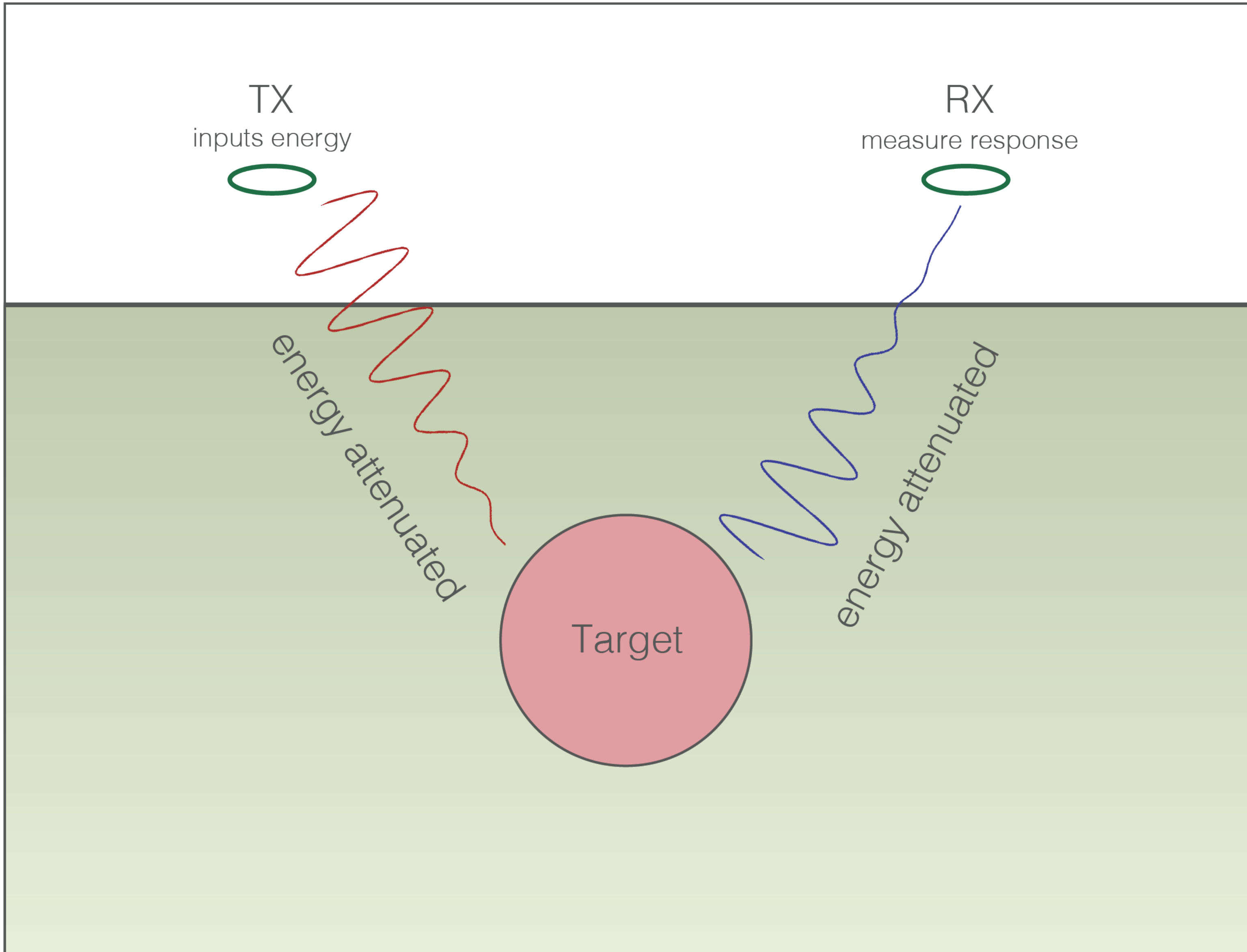
$$\mathbf{H} = \mathbf{H_0} \underbrace{e^{-\alpha z}}_{\text{attenuation}} \underbrace{e^{-i(\beta z - \omega t)}}_{\text{phase}}$$





http://em.geosci.xyz/apps.html

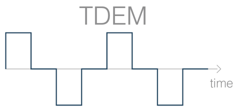
Effects of background resistivity



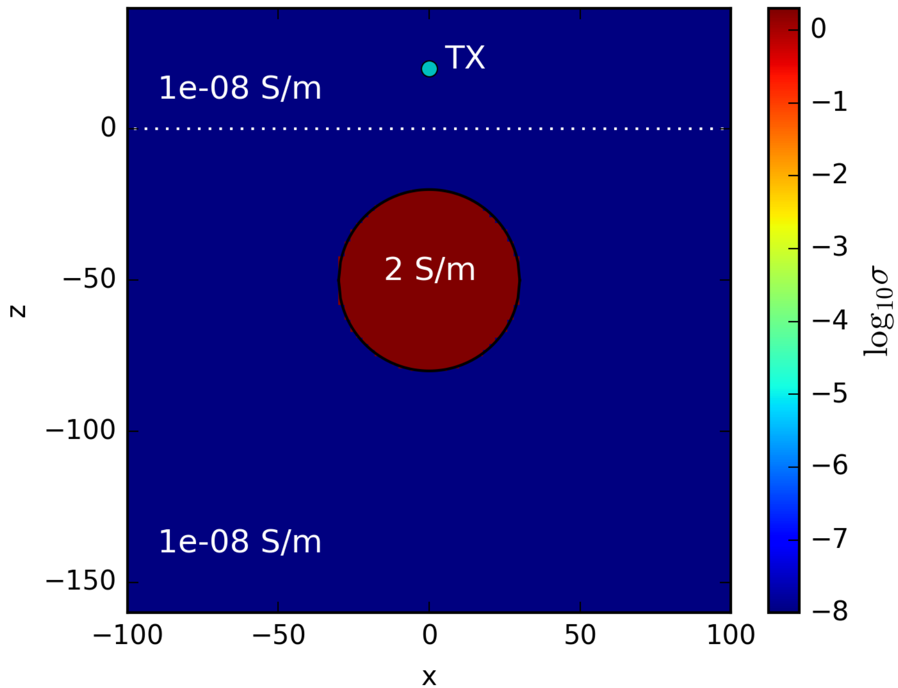
Effects of background resistivity: Time

- Buried, conductive sphere
- Vary background conductivity

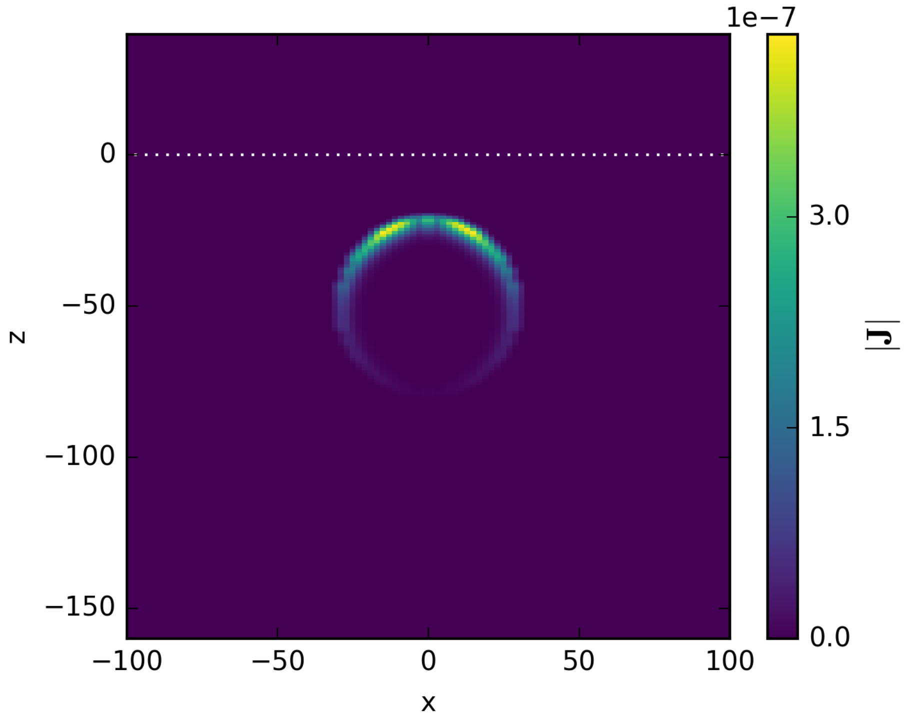
- Time: 10^{-5} s



10^{-8} S/m background



Current Density



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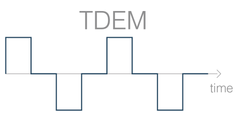
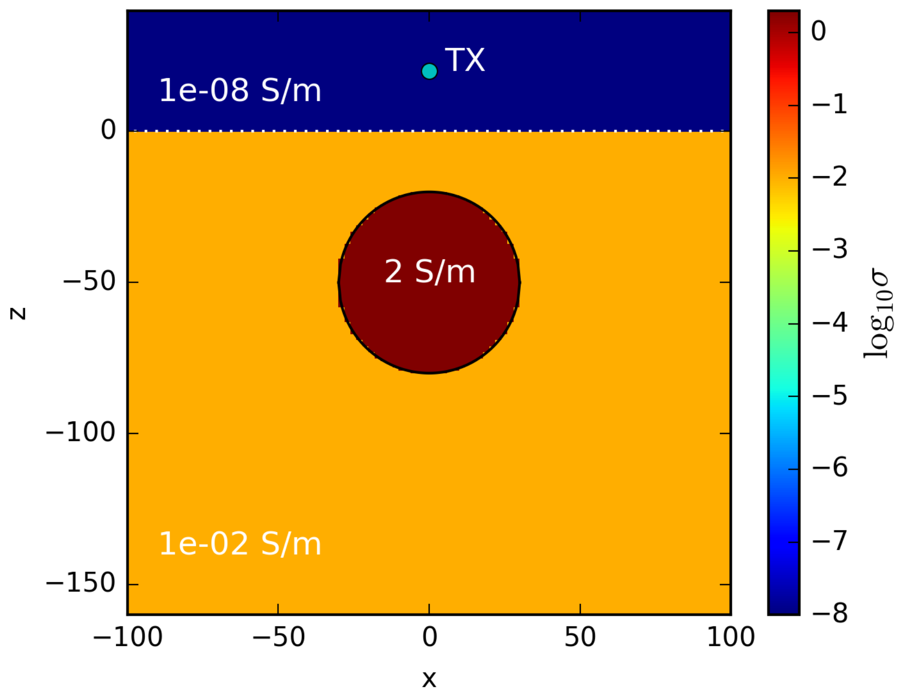
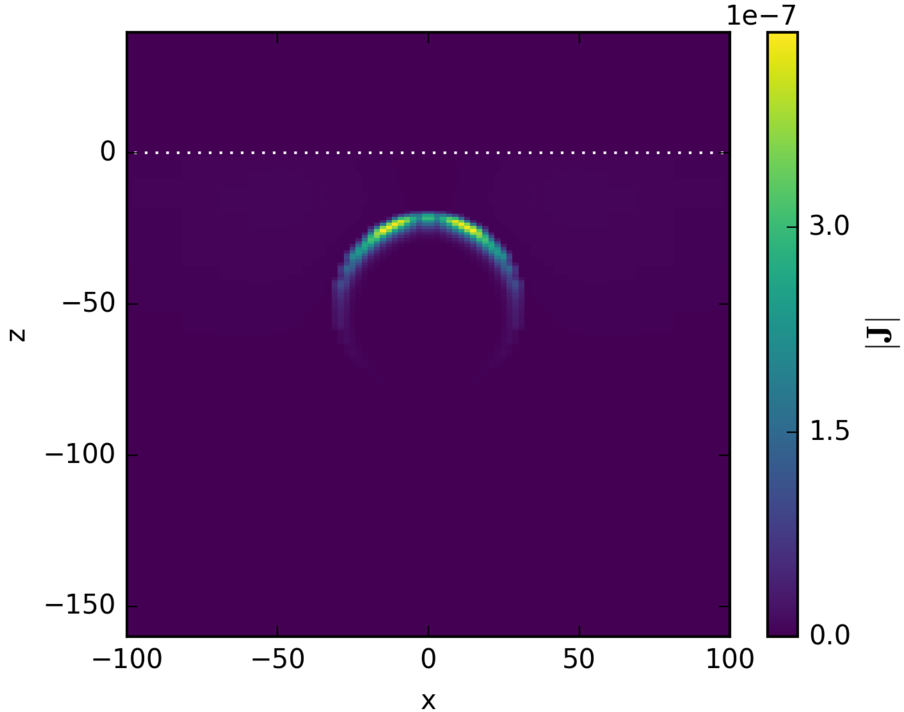


10^{-2} S/m background



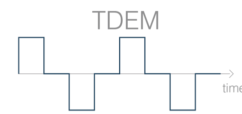
Current Density



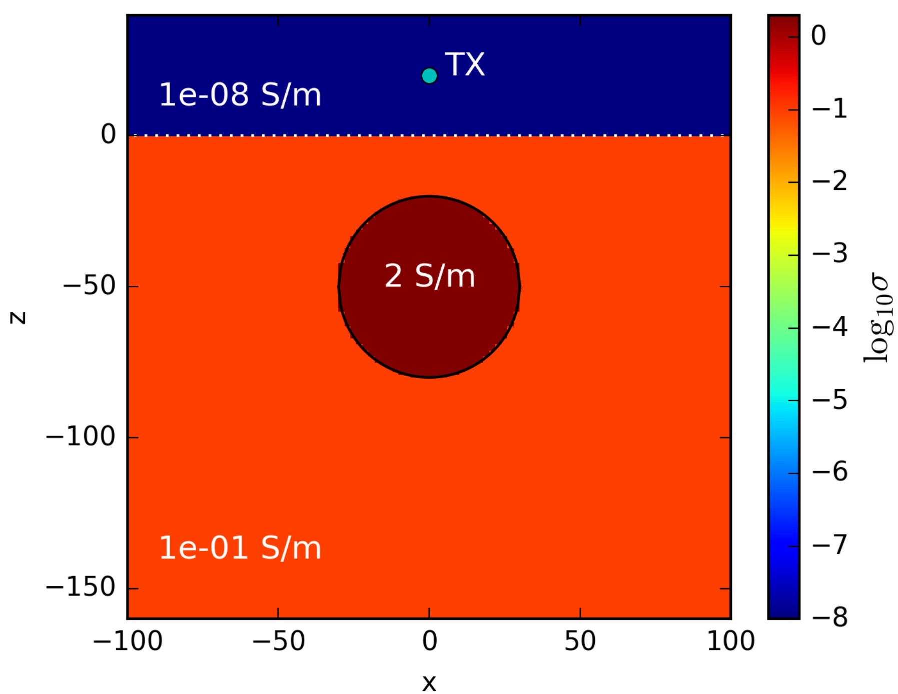
Effects of background resistivity: Time

- Buried, conductive sphere
- Vary background conductivity

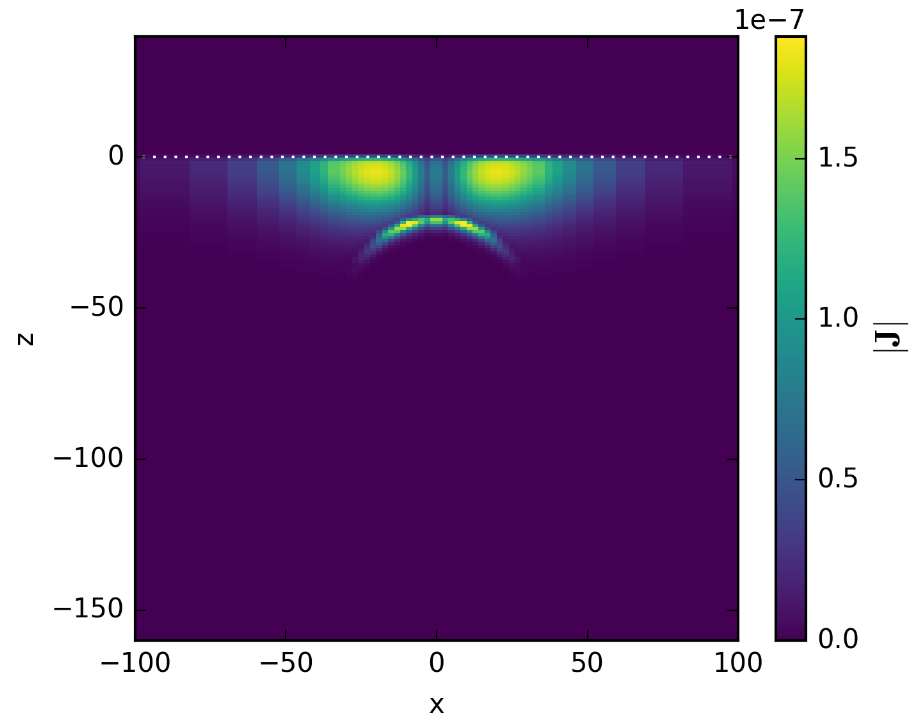
- Time: 10^{-5} s



10^{-1} S/m background

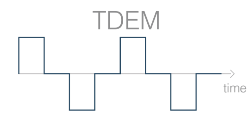


Current Density

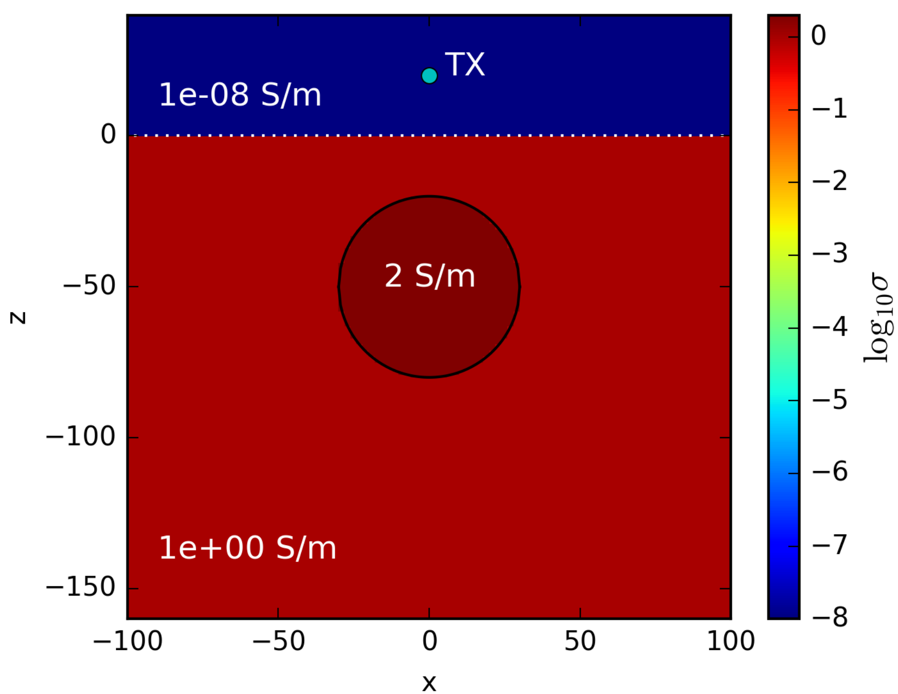


Effects of background resistivity: Time

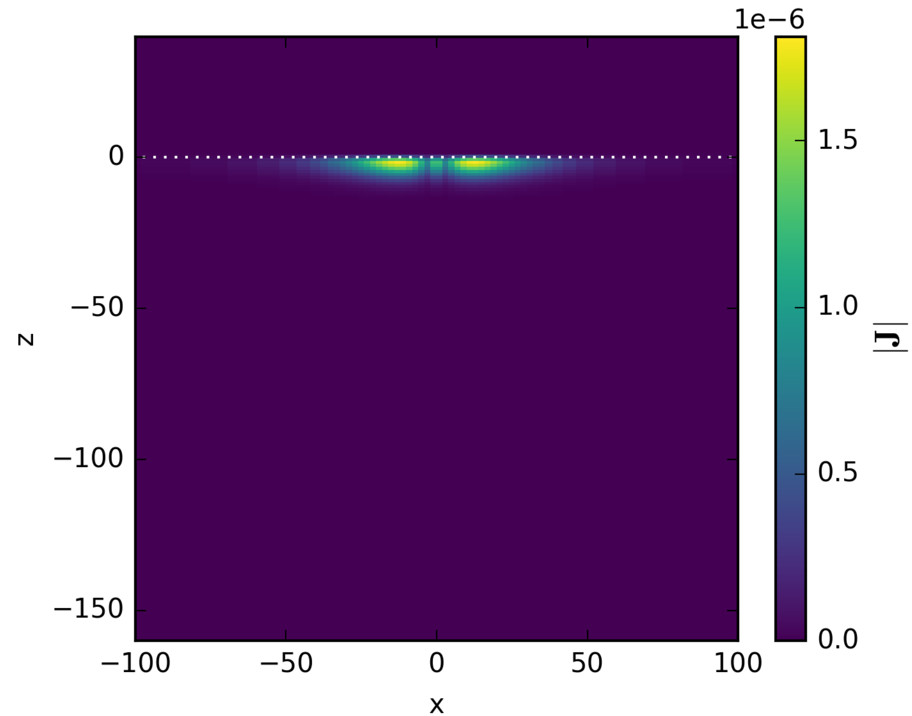
- Buried, conductive sphere
- Vary background conductivity
- Time: 10^{-5} s



1 S/m background

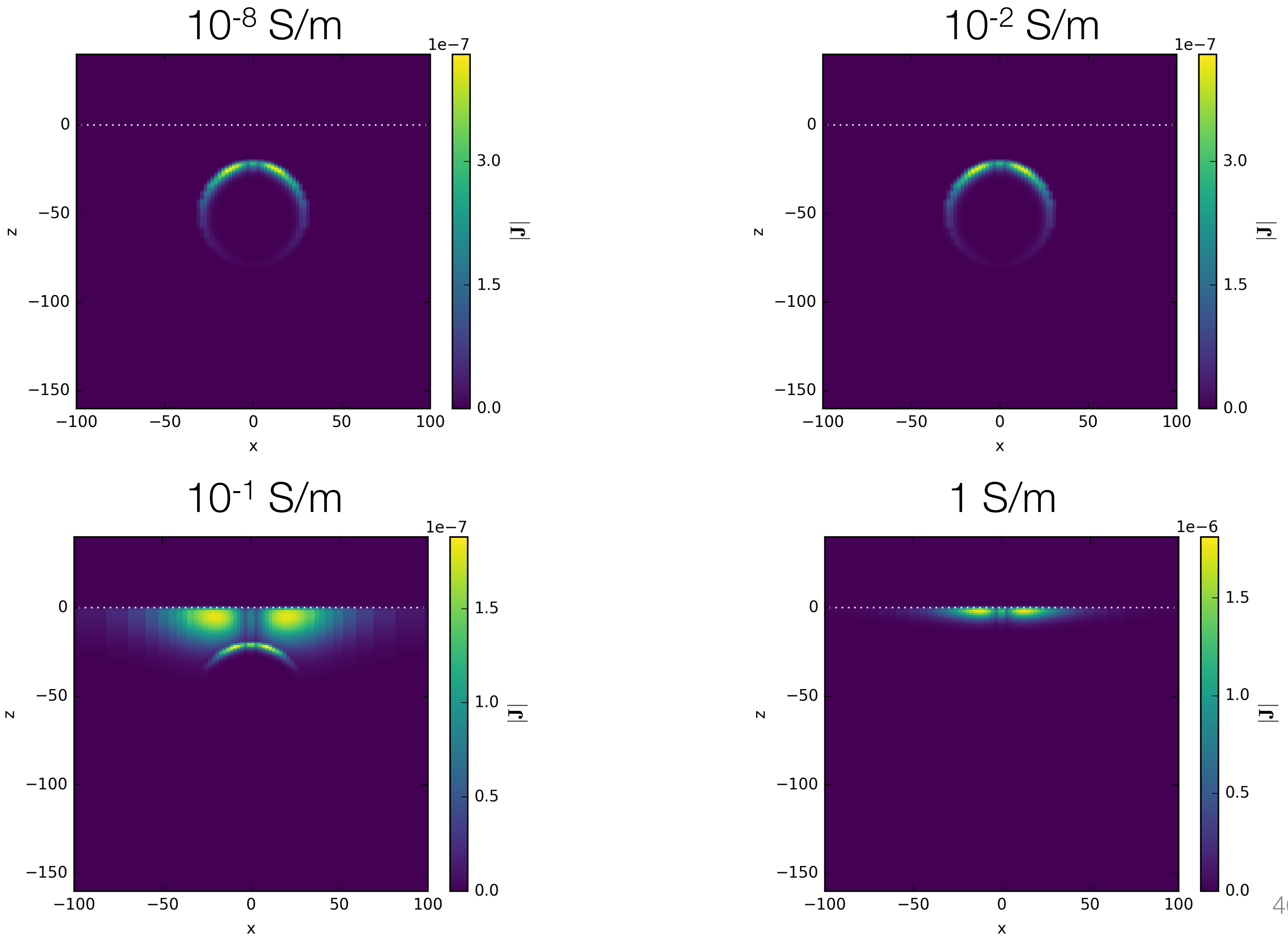


Current Density

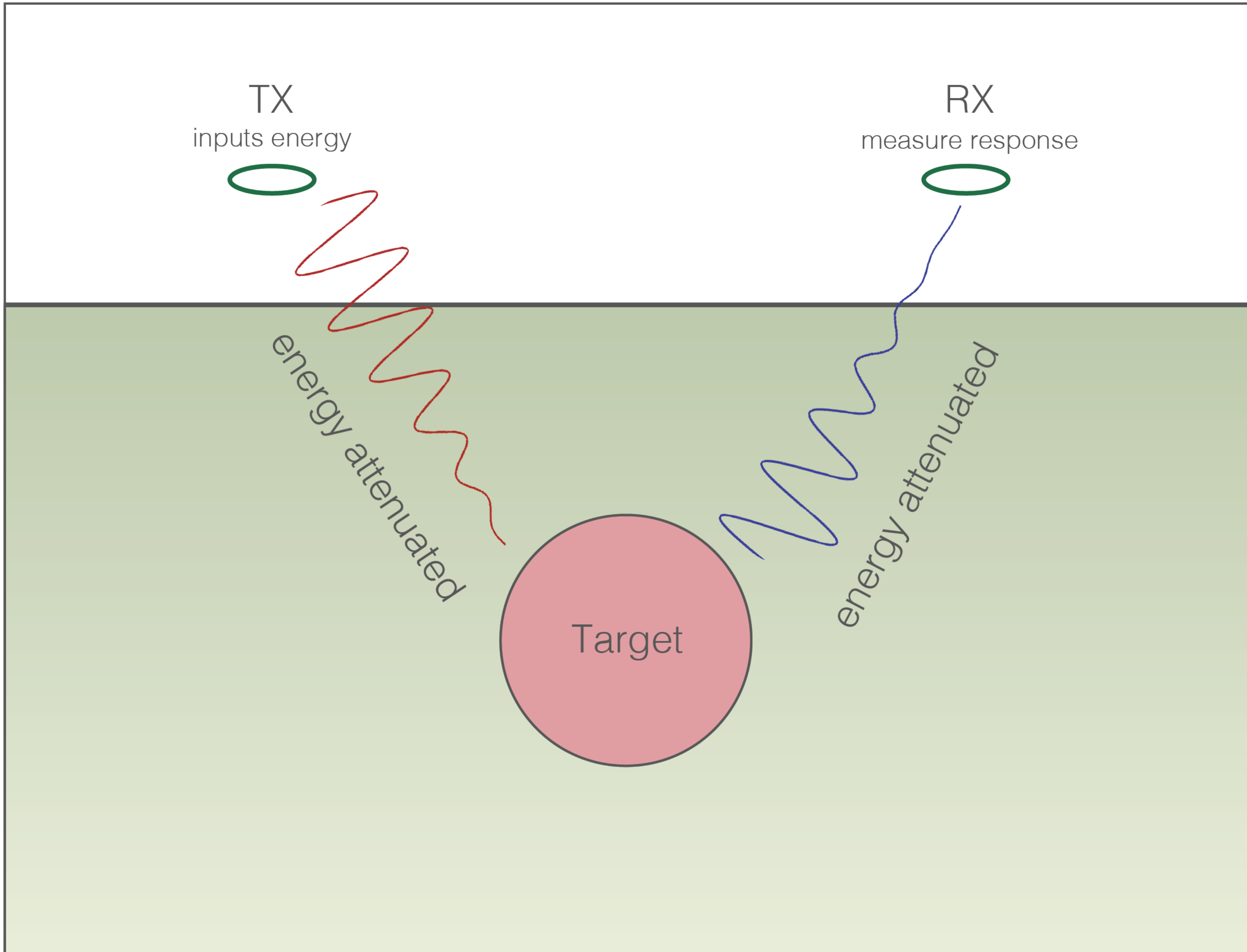


10^{-5} s

Effects of background resistivity: Time

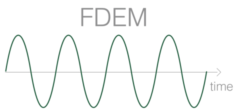


Effects of background resistivity

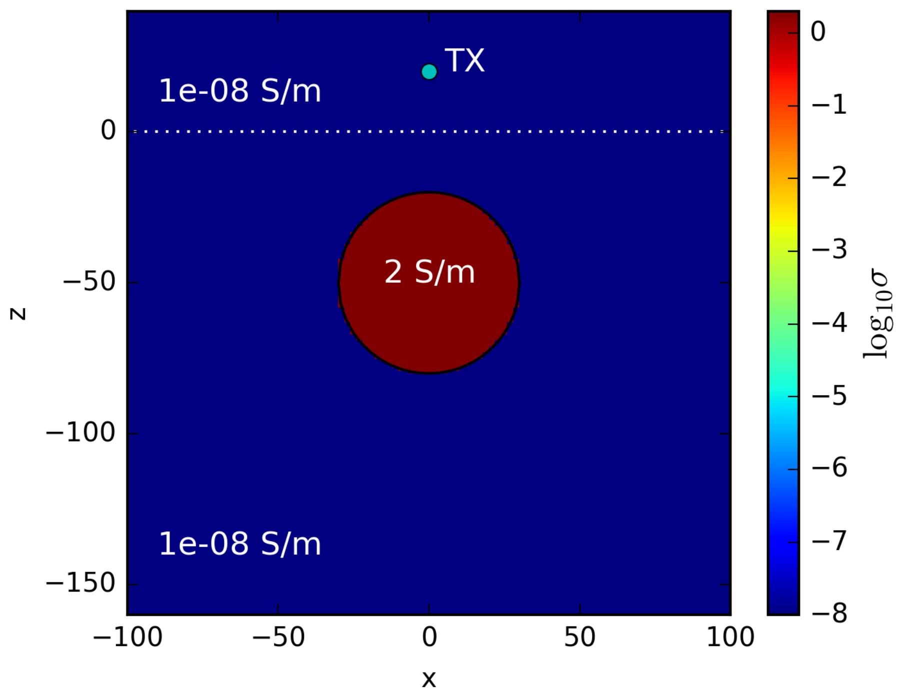


Effects of background resistivity: Frequency

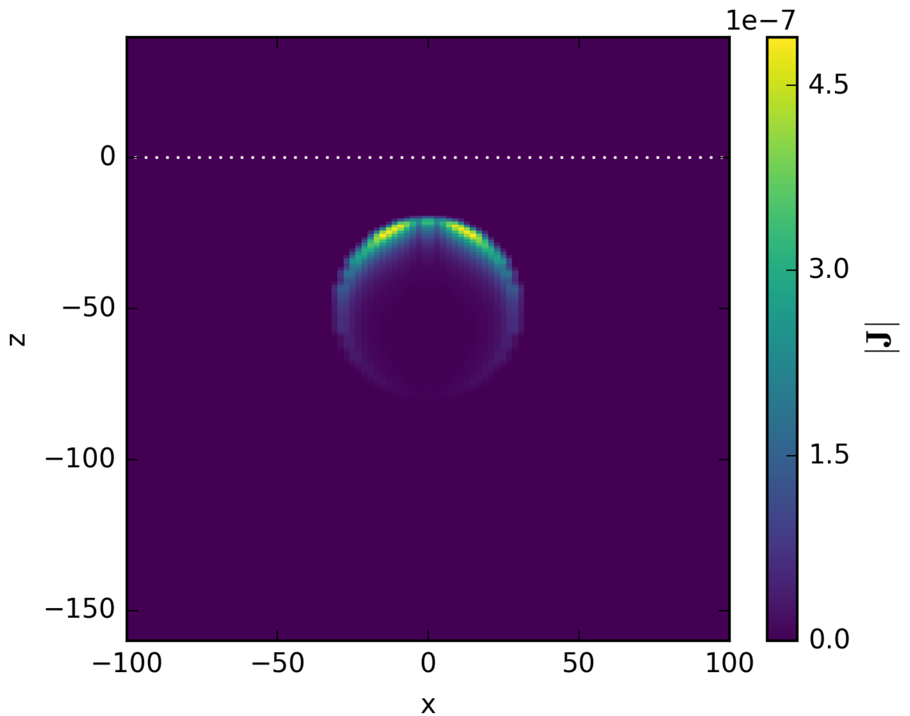
- Buried, conductive sphere
- Vary background conductivity
- Frequency: 10^4 Hz



10^{-8} S/m background



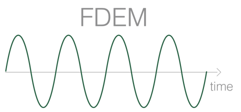
Current Density

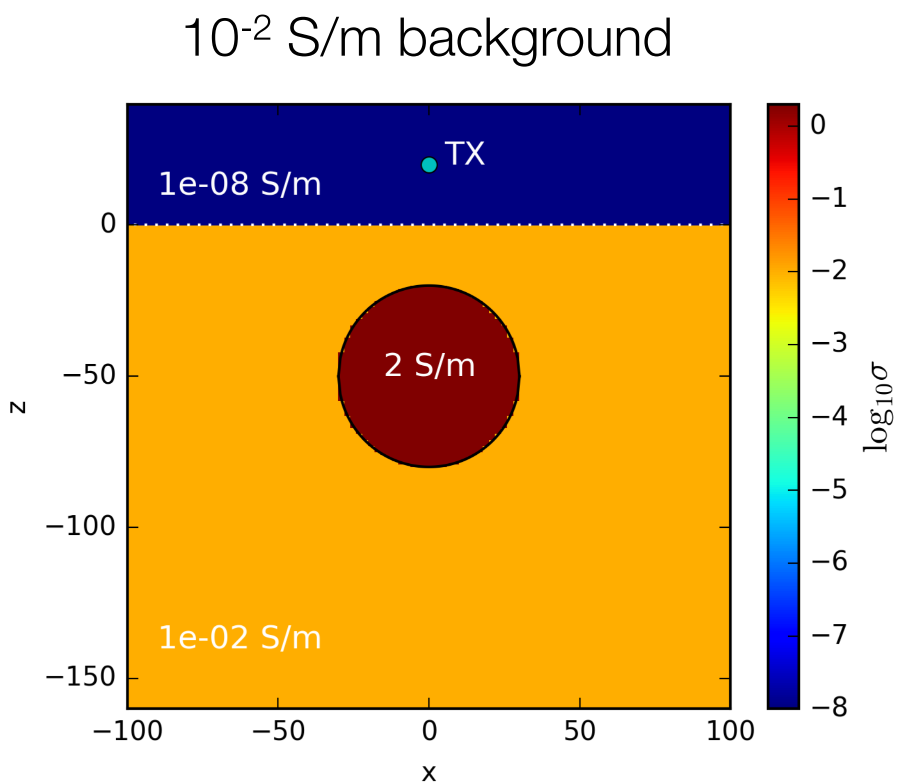


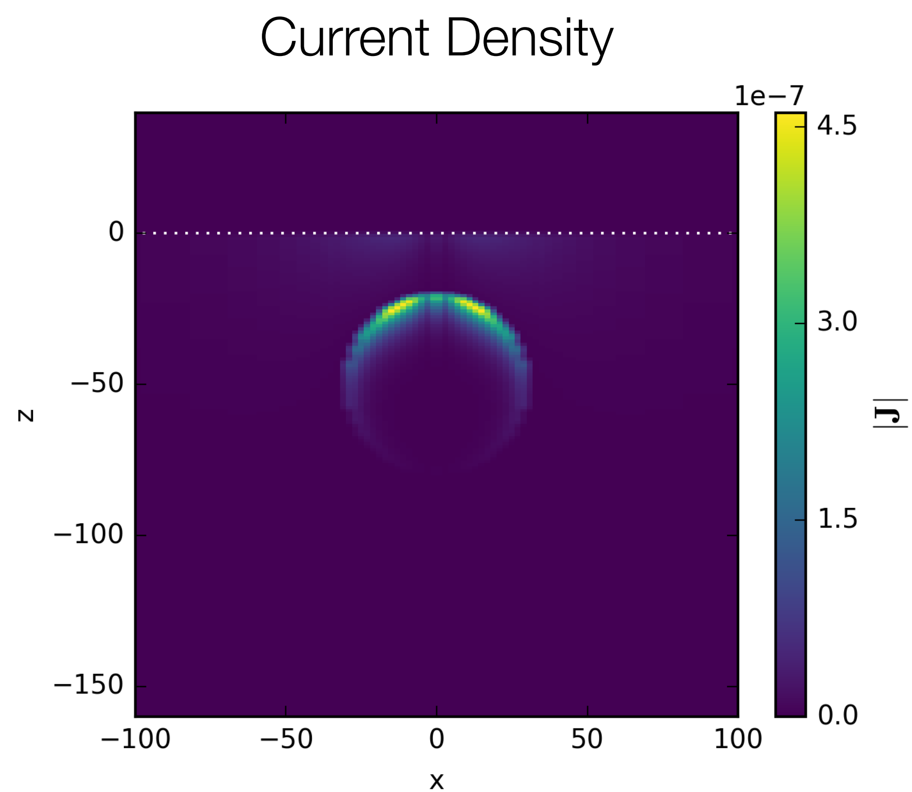
Effects of background resistivity: Frequency

- Buried, conductive sphere
- Vary background conductivity

- Frequency: 10^4 Hz







Effects of background resistivity: Frequency

- Buried, conductive sphere
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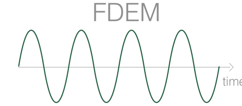
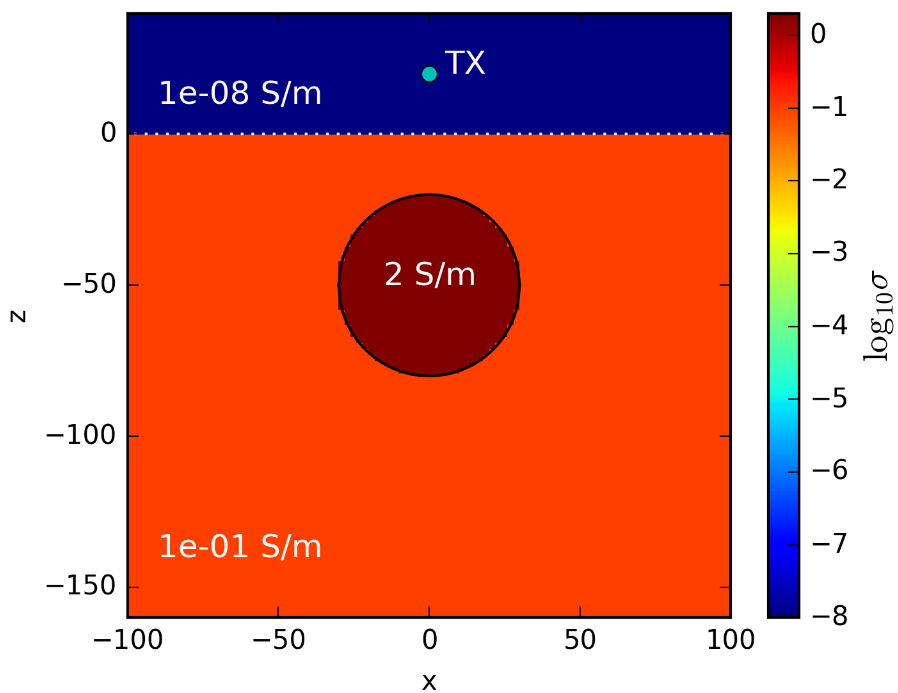
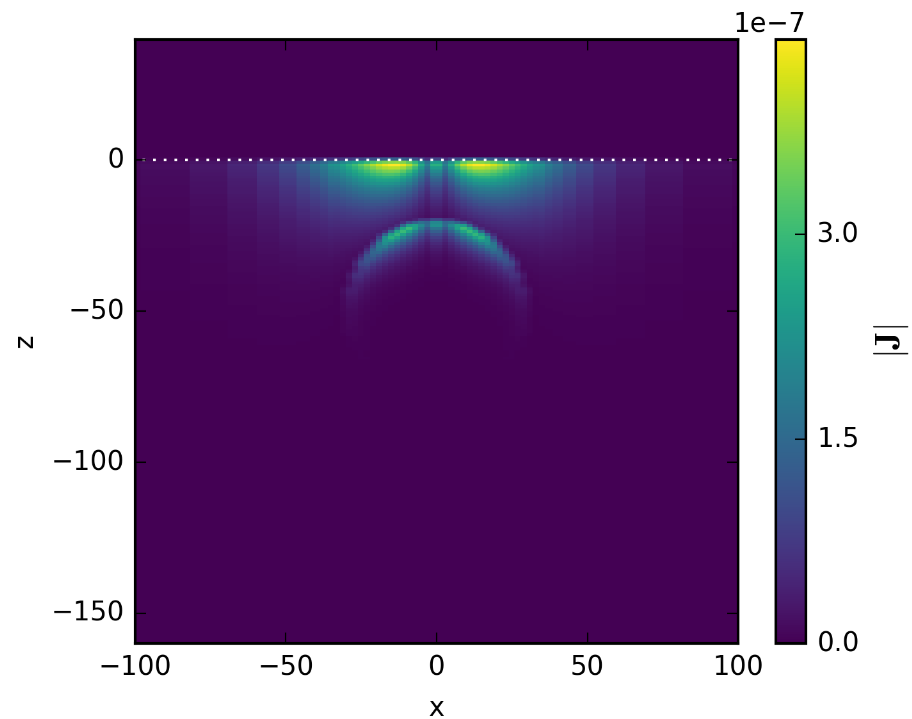
- Frequency: 10^4 Hz



10^{-1} S/m background



Current Density



Effects of background resistivity: Frequency

- Buried, conductive sphere
- Vary background conductivity
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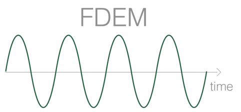
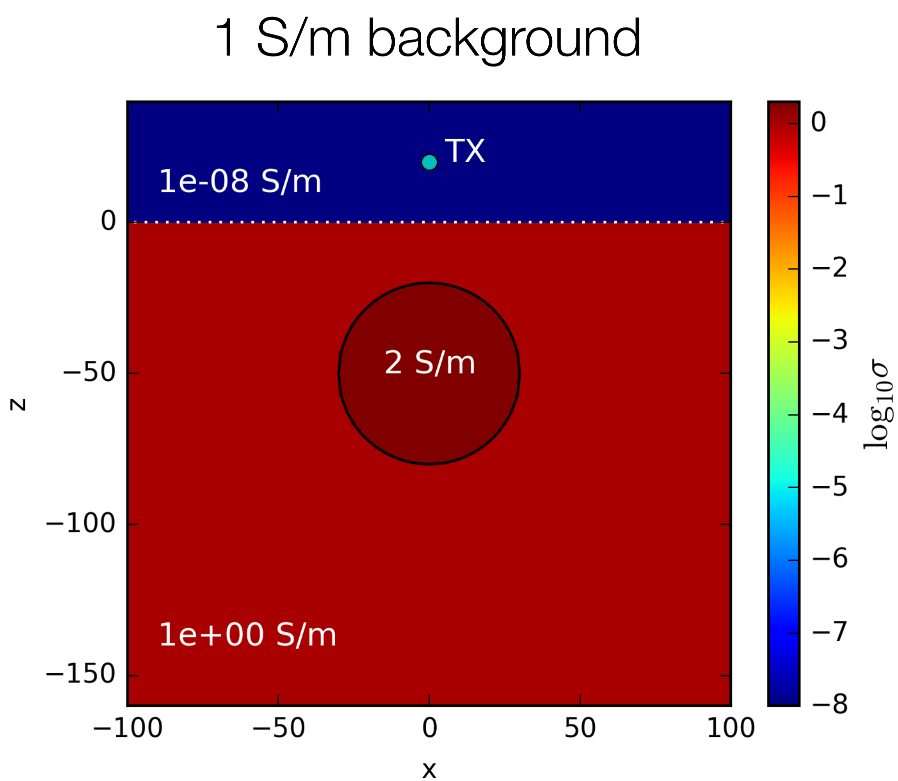
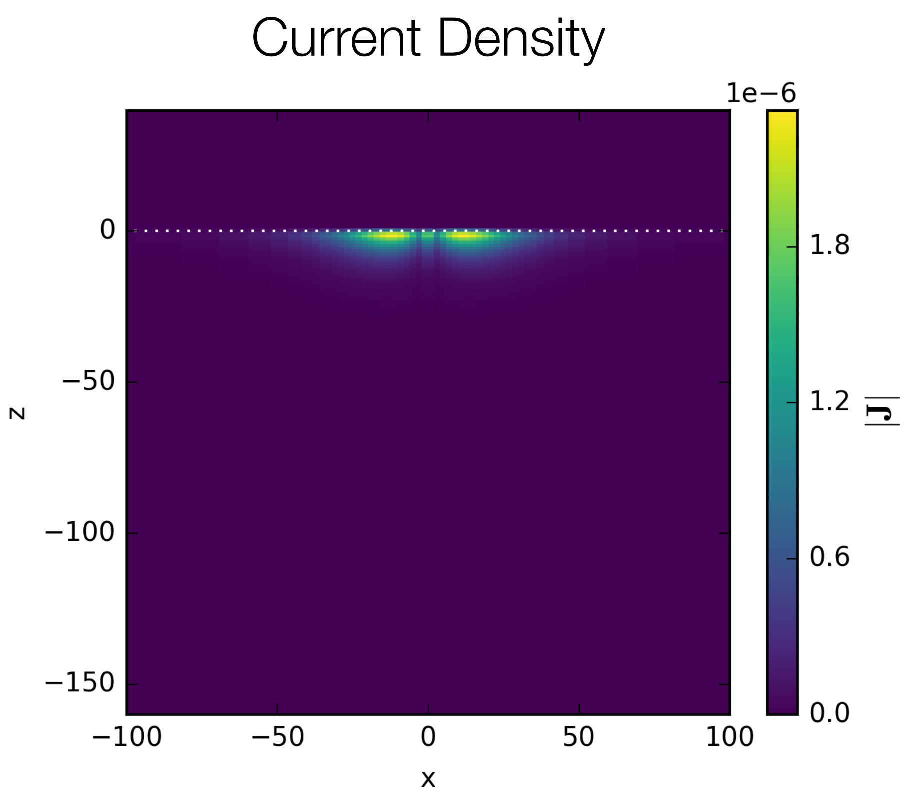
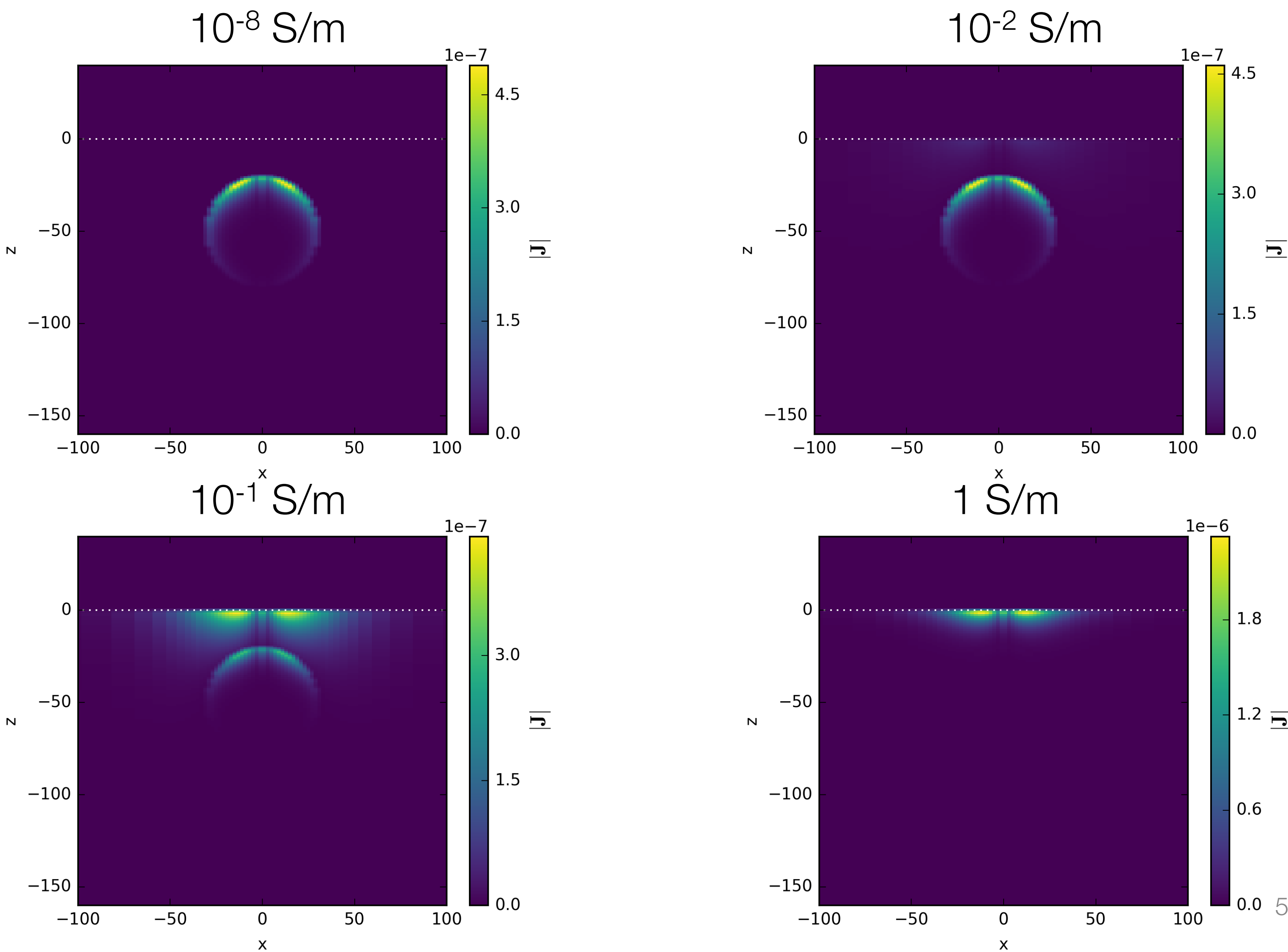






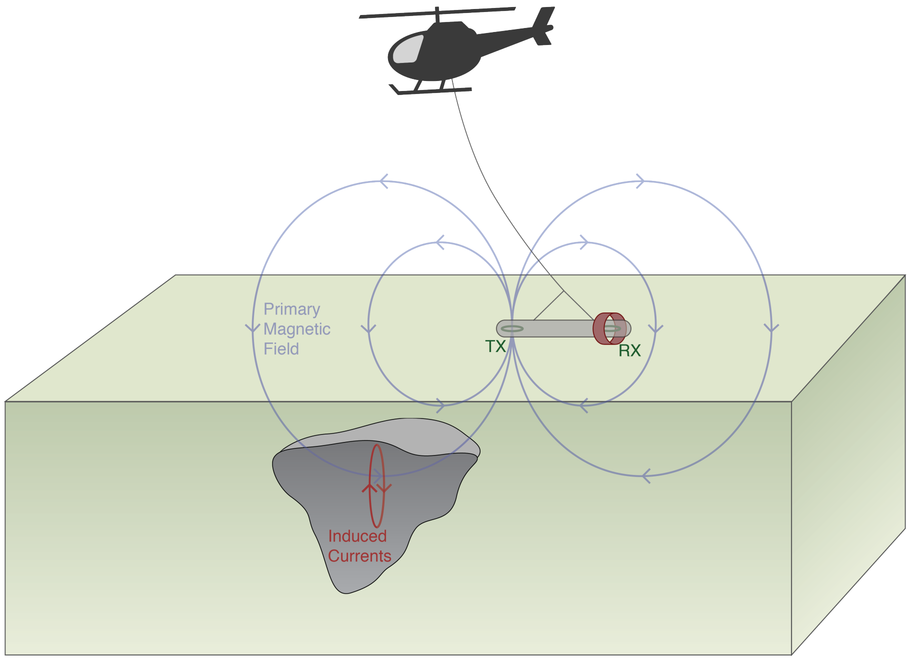
Effects of background resistivity: Frequency



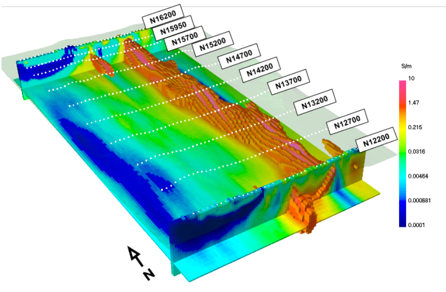
Recap: what have we learned?

- Basics of EM induction
- Response functions
- Mutual coupling
- Data for frequency or time domain systems
- Circuit model is a good proxy
- Need to account for energy losses

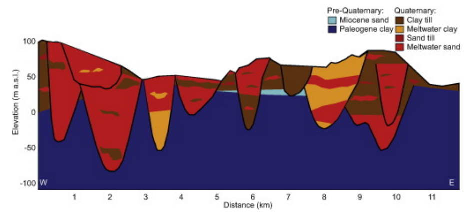
- Ready to look at some field examples



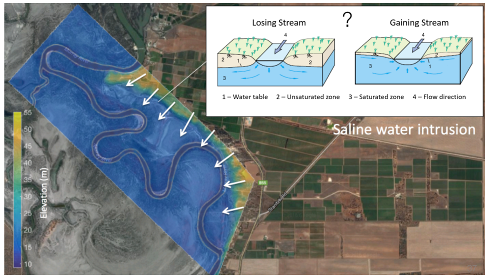
Today’s Case Histories



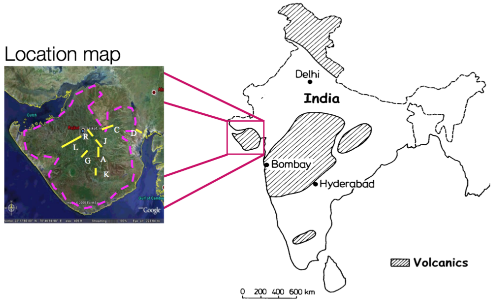
Mt. Isa, Australia: Mineral Exploration



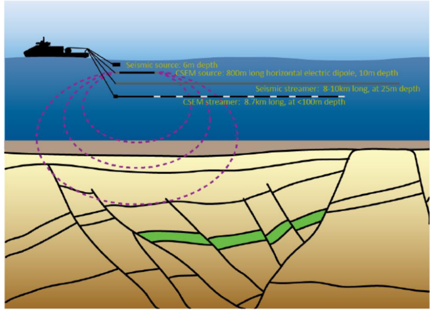
Kasted, Denmark: mapping paleochannels



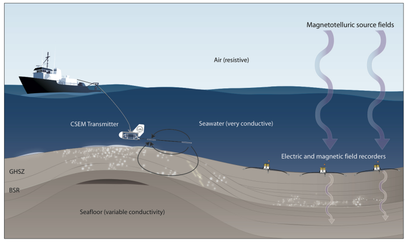
Bookpurnong, Australia: diagnosing river salinization



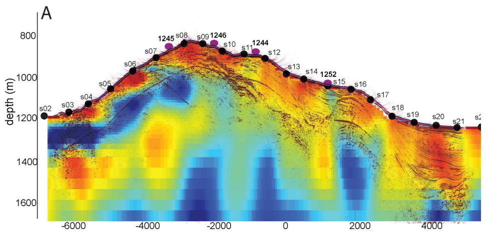
Deccan Traps, India: mapping sediment beneath basalt



Barents Sea: Hydrocarbon de-risking

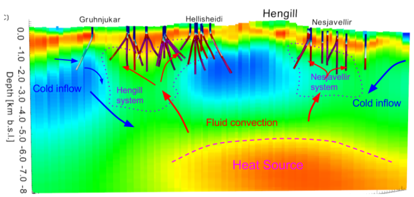


Marine CSEM

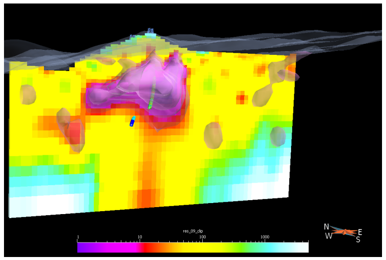


Oregon, USA: Methane Hydrates

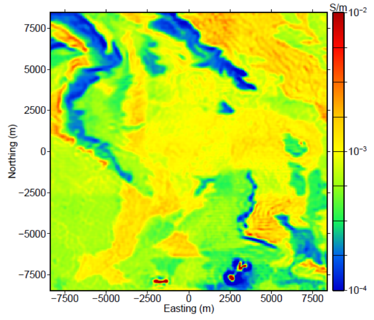
Today's Case Histories



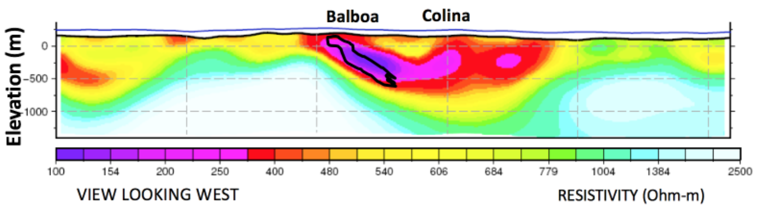
Iceland: characterizing geothermal systems



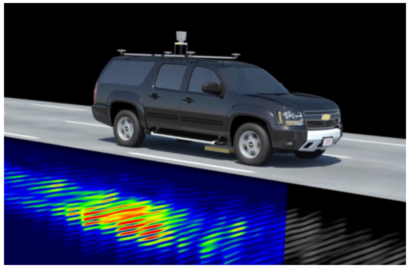
Santa Cecilia, Chile: Mineral Exploration



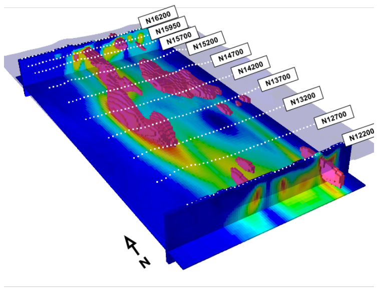
Noranda, Canada: Geologic Mapping



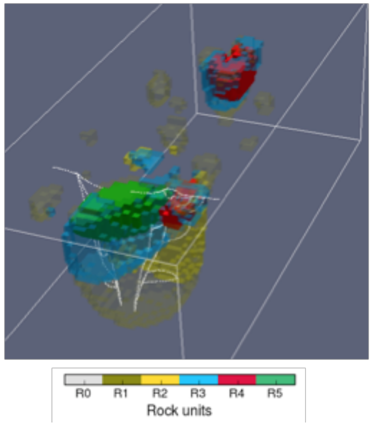
Balboa, Panama: Mineral Exploration



USA: Self-driving vehicles



Mt. Isa, Australia: Mineral Exploration



TKC, Canada: Mineral Exploration

End of EM Fundamentals