

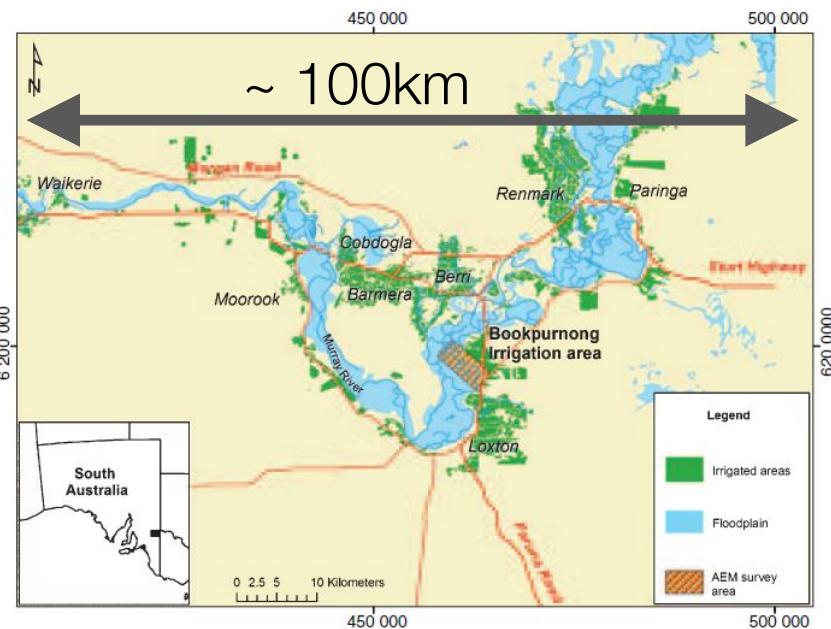
EM Fundamentals



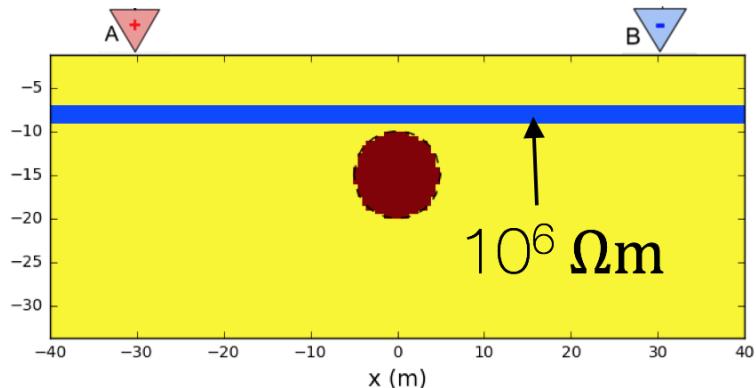
$$\nabla \times$$

Motivation: applications difficult for DC

Large areas to be covered



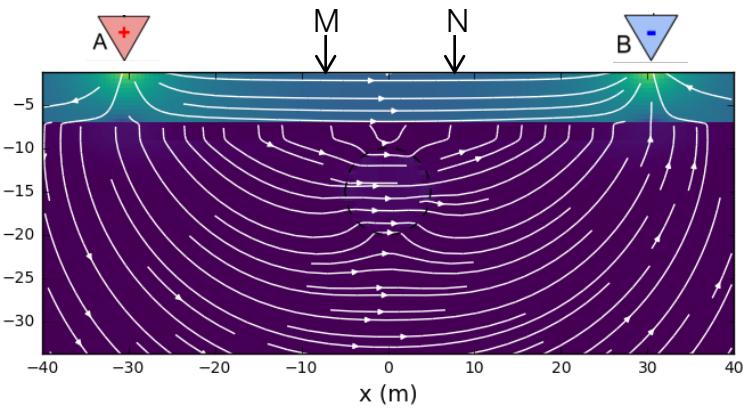
Resistive layer “shields” target



Rugged terrain



Hard to inject



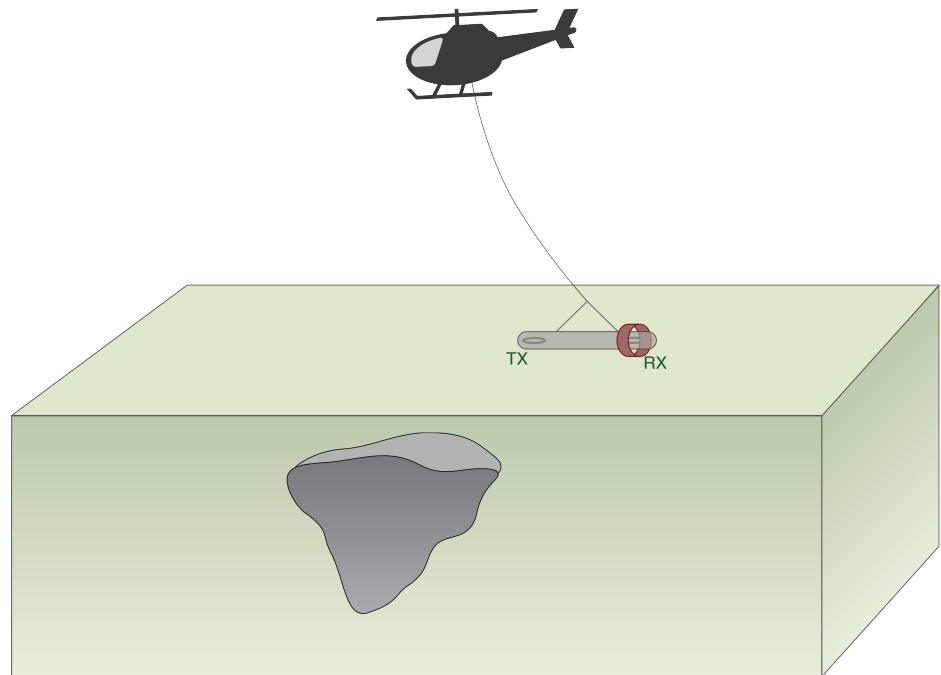
Outline

- Basic Survey
- Ampere's and Faraday's Laws (2-coil App)
- Circuit model for EM induction
- Frequency and time domain data
- Sphere in homogeneous earth
- Cyl code
- Energy losses in the ground

Basic Experiment

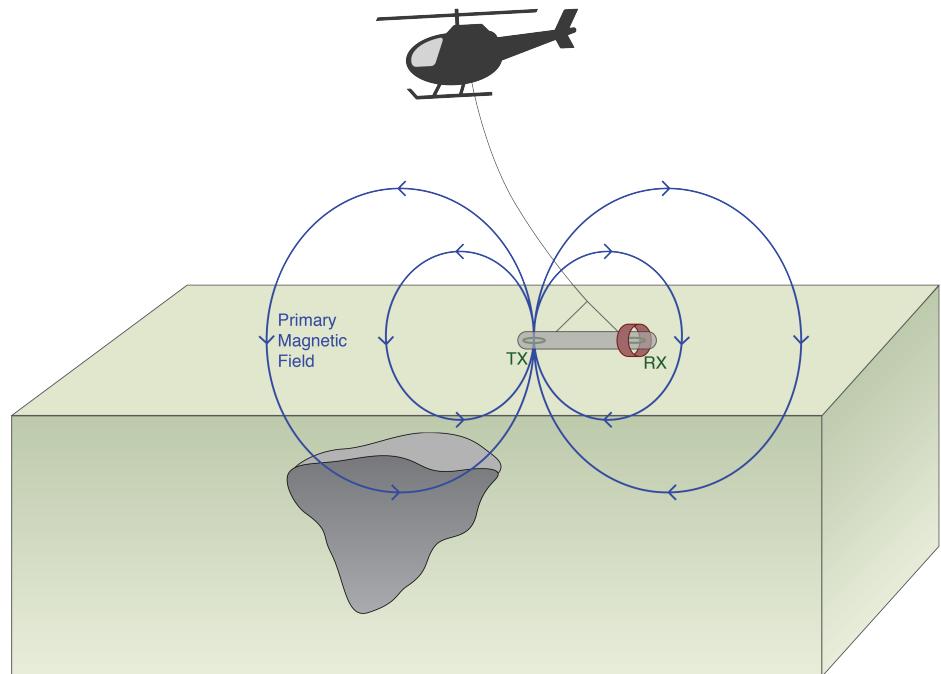
- **Setup:**

- transmitter and receiver are in a towed bird



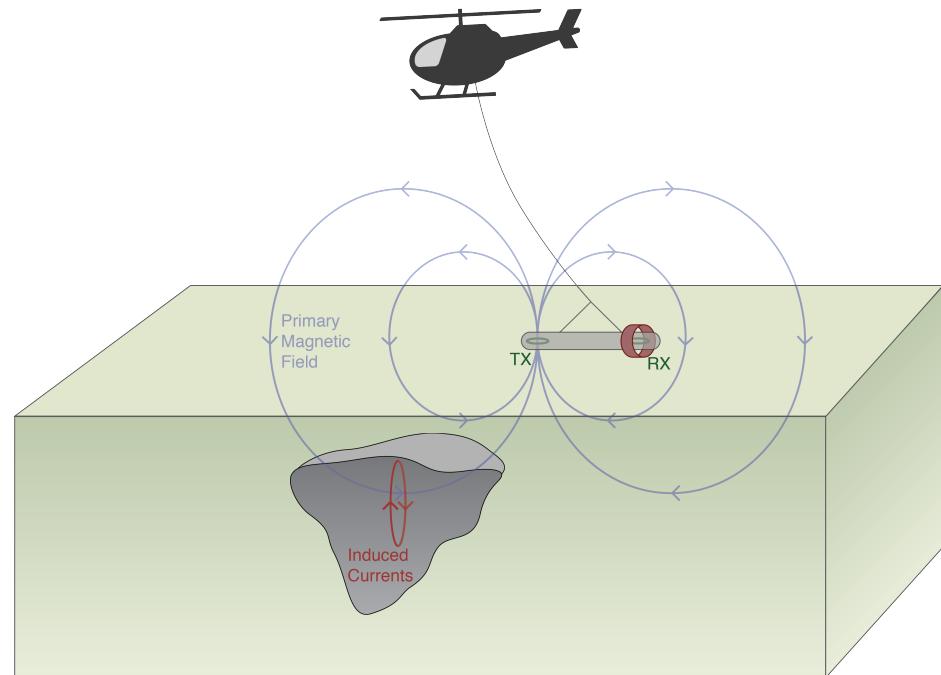
Basic Experiment

- **Setup:**
 - transmitter and receiver are in a towed bird
- **Primary:**
 - Transmitter produces a primary magnetic field



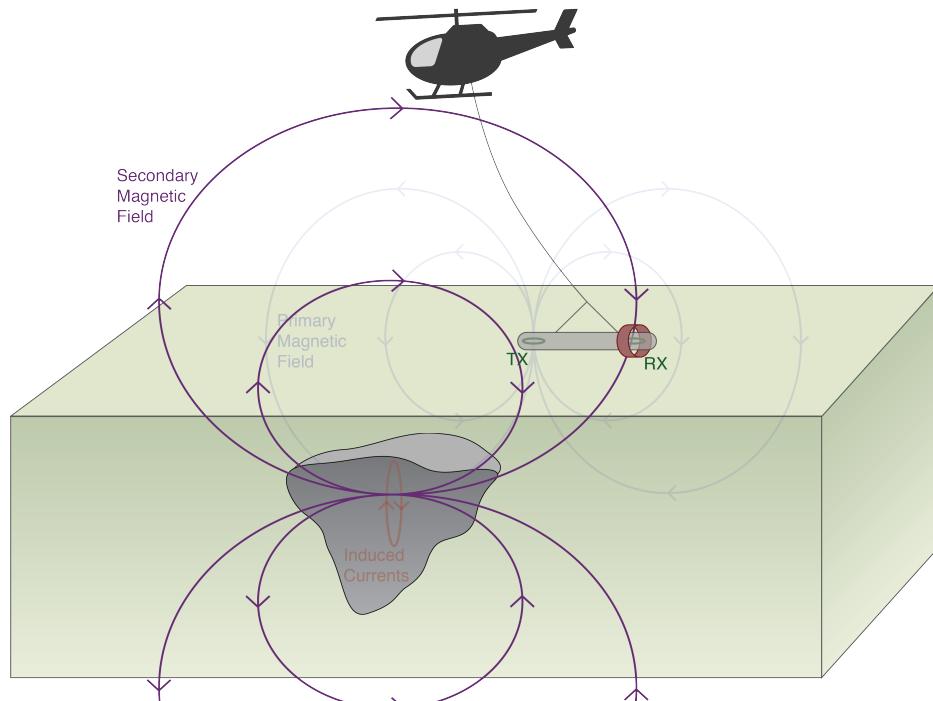
Basic Experiment

- **Setup:**
 - transmitter and receiver are in a towed bird
- **Primary:**
 - Transmitter produces a primary magnetic field
- **Induced Currents:**
 - Time varying magnetic fields generate electric fields everywhere and currents in conductors



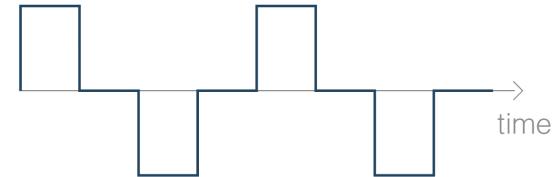
Basic Experiment

- **Setup:**
 - transmitter and receiver are in a towed bird
- **Primary:**
 - Transmitter produces a primary magnetic field
- **Induced Currents:**
 - Time varying magnetic fields generate electric fields everywhere and currents in conductors
- **Secondary Fields:**
 - The induced currents produce a secondary magnetic field.

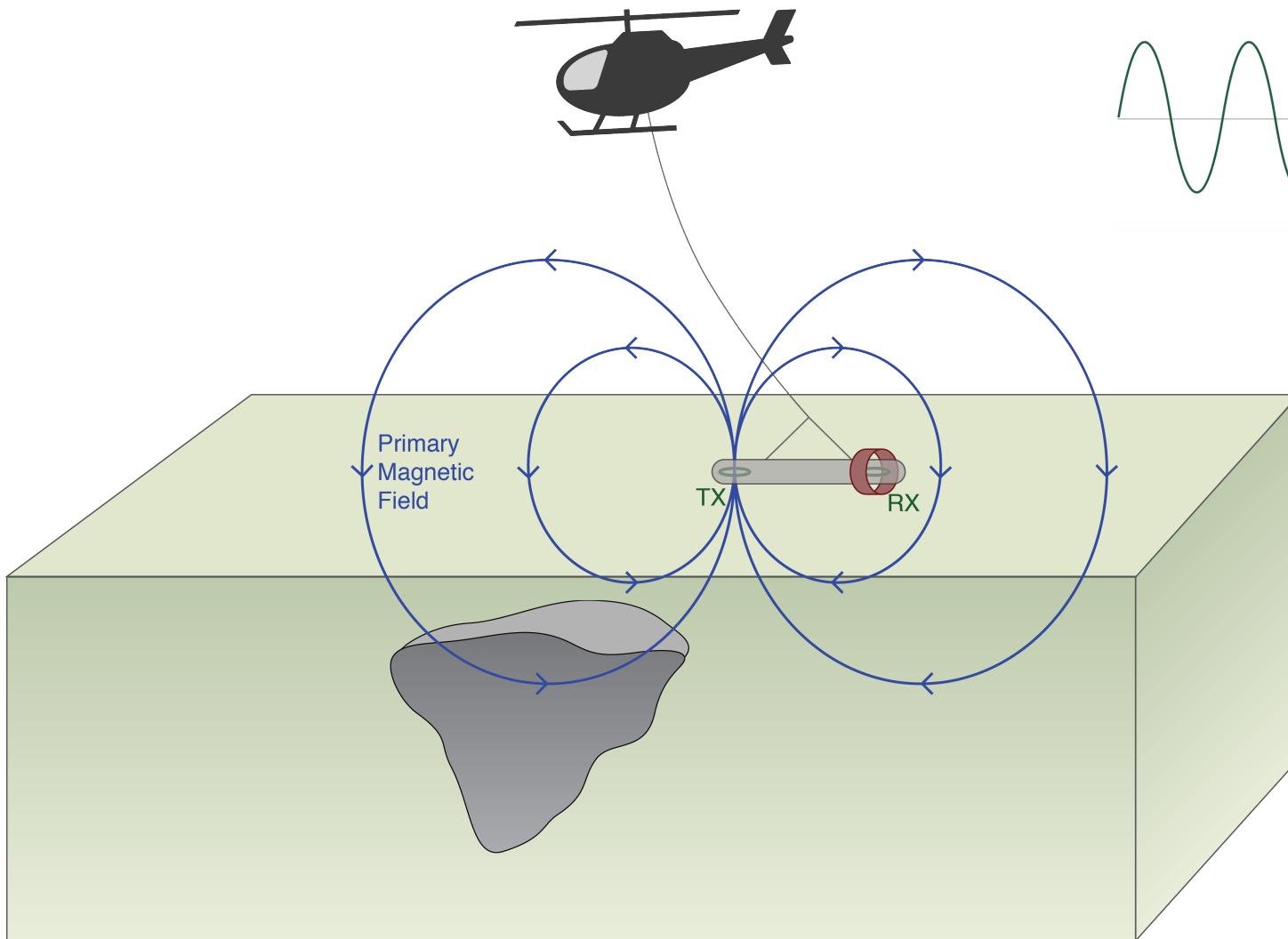
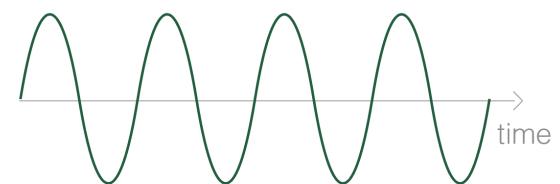


Transmitter

waveform



or



Basic Equations: Quasi-static

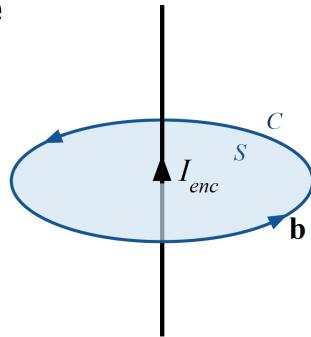
	Time	Frequency
Faraday's Law	$\nabla \times \mathbf{e} = - \frac{\partial \mathbf{b}}{\partial t}$	$\nabla \times \mathbf{E} = - i\omega \mathbf{B}$
Ampere's Law	$\nabla \times \mathbf{h} = \mathbf{j} + \frac{\partial \mathbf{d}}{\partial t}$	$\nabla \times \mathbf{H} = \mathbf{J} + i\omega \mathbf{D}$
No Magnetic Monopoles	$\nabla \cdot \mathbf{b} = 0$	$\nabla \cdot \mathbf{B} = 0$
Constitutive Relationships (non-dispersive)	$\mathbf{j} = \sigma \mathbf{e}$ $\mathbf{b} = \mu \mathbf{h}$ $\mathbf{d} = \epsilon \mathbf{e}$	$\mathbf{J} = \sigma \mathbf{E}$ $\mathbf{B} = \mu \mathbf{H}$ $\mathbf{D} = \epsilon \mathbf{E}$

* Solve with sources and boundary conditions

Ampere's Law

$$\nabla \times \mathbf{H} = \mathbf{J}$$

Wire



$$\mathbf{B} = \frac{\mu_0 I_{enc}}{2\pi r} \hat{\phi}$$

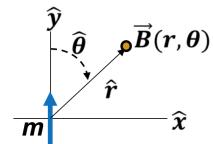
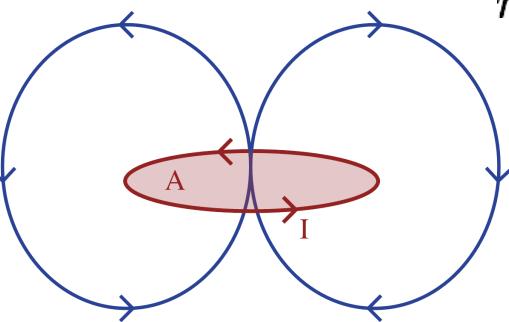
Right hand rule

Current loop

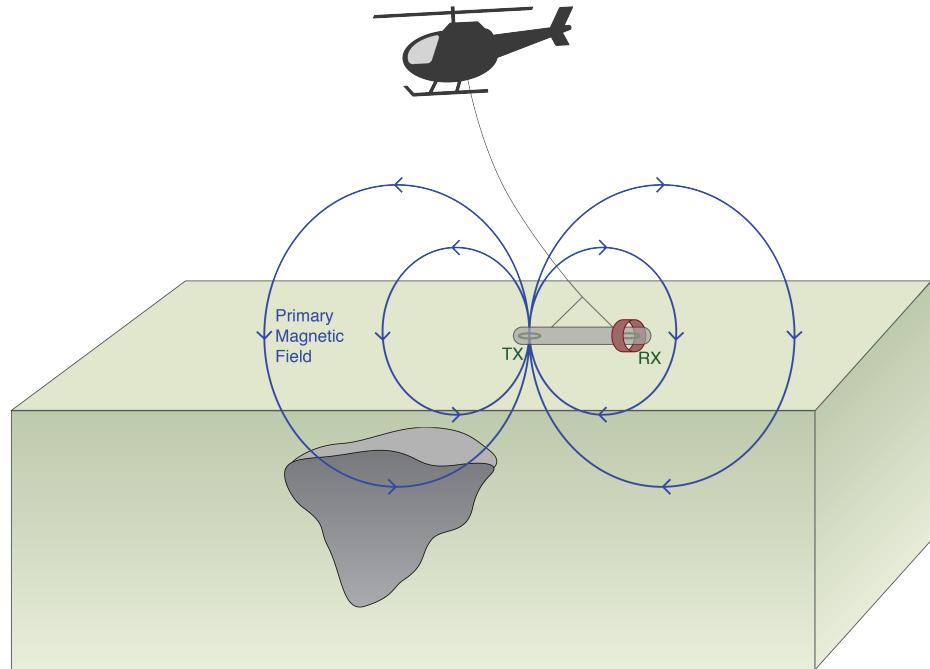
$$\mathbf{B} = \frac{\mu_0 m}{4\pi r^3} (2 \cos \theta \hat{\mathbf{r}} + \sin \theta \hat{\theta})$$

$$m = IA$$

Primary
Magnetic
Field



$$\vec{B}(r, \theta)$$



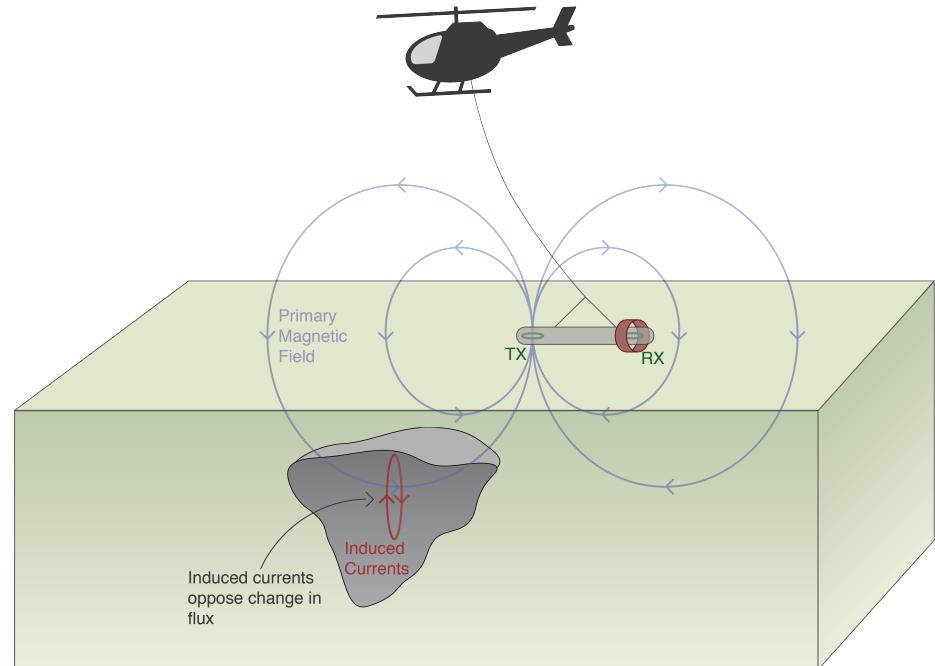
Faraday's Law and Induced Currents

$$\nabla \times \mathbf{e} = -\frac{\partial \mathbf{b}}{\partial t}$$

Lenz' Law

Ohm's Law

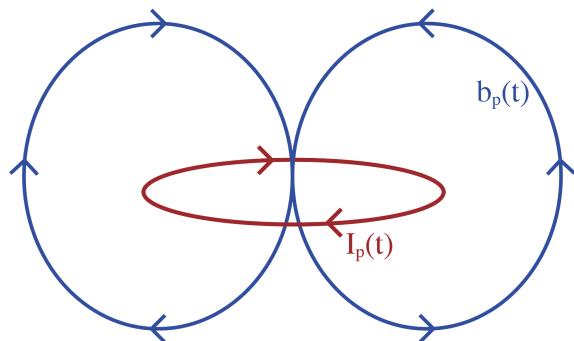
$$\mathbf{j} = \sigma \mathbf{e}$$



Two Coil Example: Harmonic

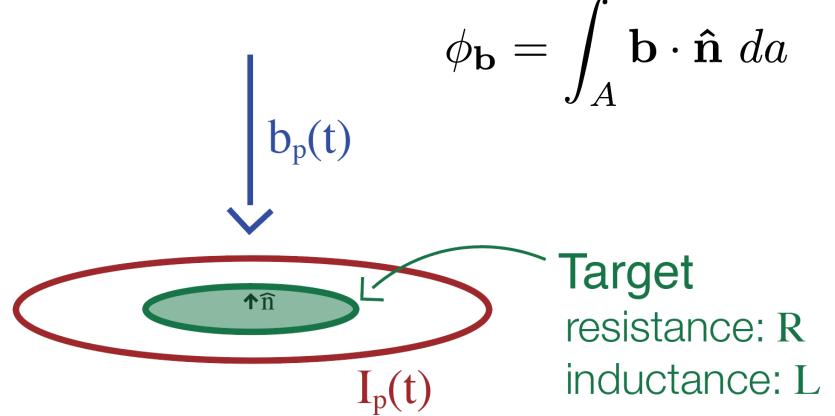
Source (red loop)

- Time varying current \rightarrow Time varying magnetic flux



Target (green loop)

- Time varying magnetic flux



Faraday's Law

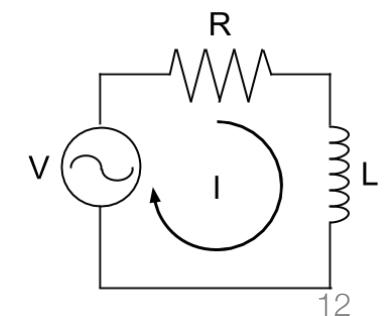
$$\nabla \times \mathbf{e} = -\frac{\partial \mathbf{b}}{\partial t}$$

Ohm's Law

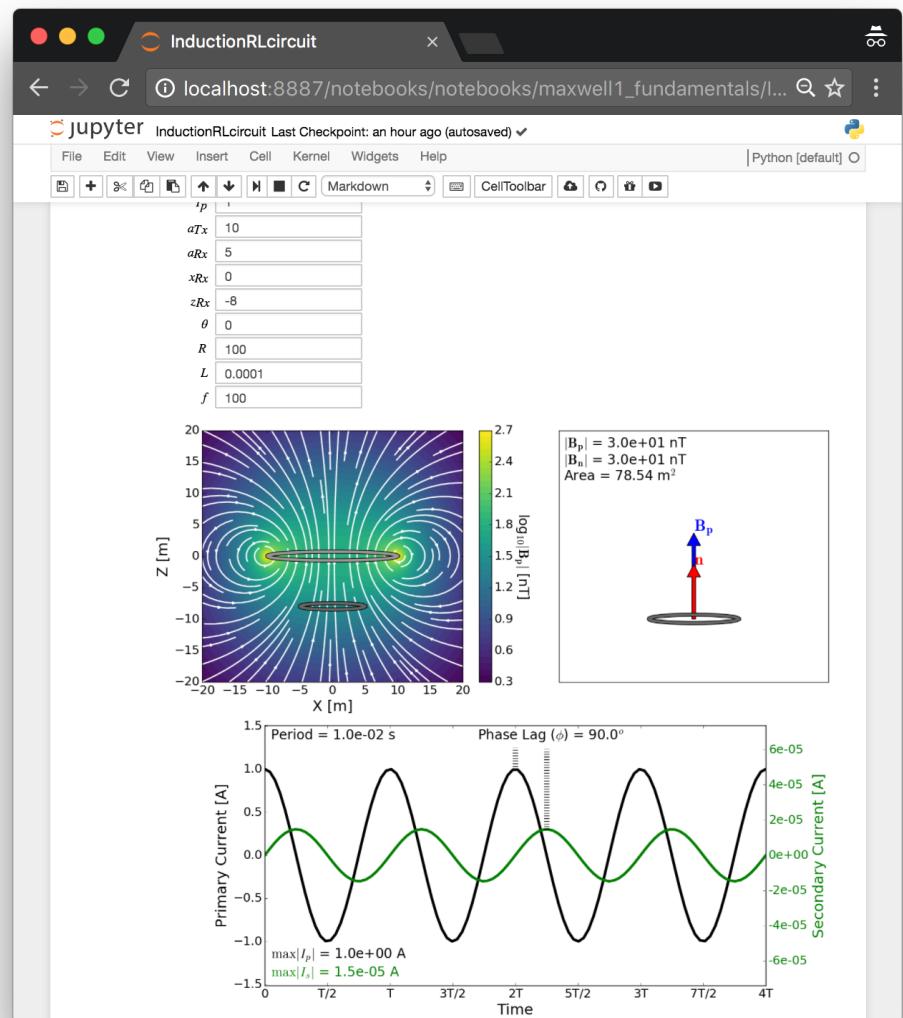
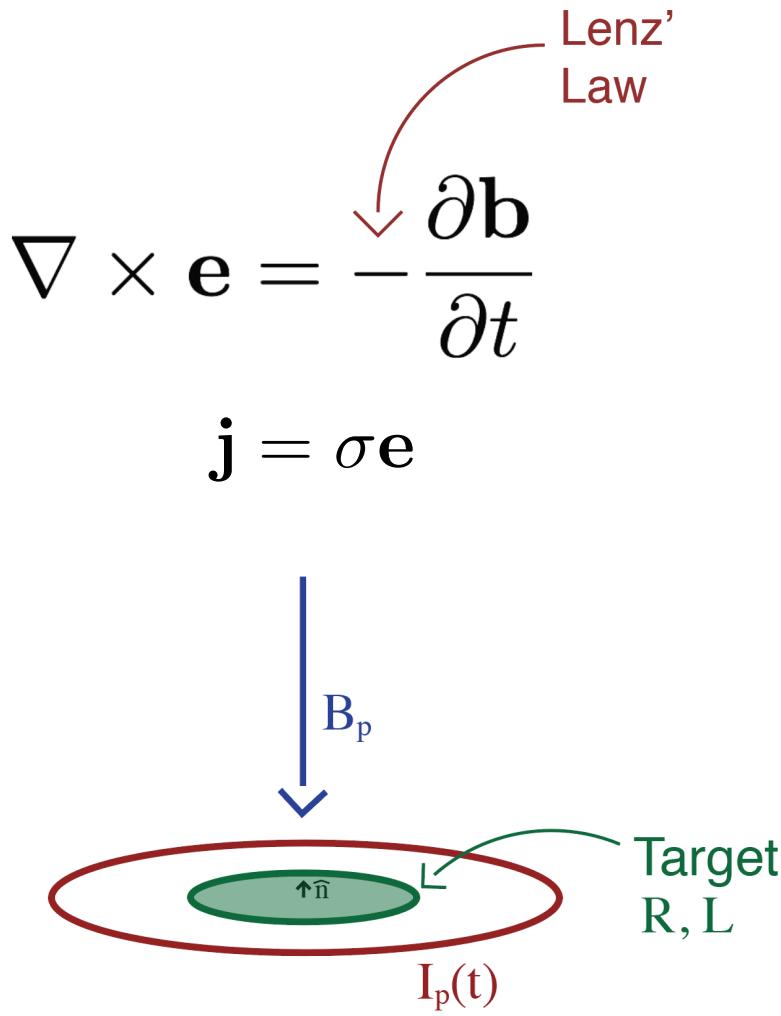
$$\mathbf{j} = \sigma \mathbf{e}$$

EMF (voltage) is related to time rate of change in flux.

$$V = EMF = -\frac{d\phi_b}{dt}$$



App for Faraday's Law



Two Coil Example: Harmonic

Induced Currents

$$I_p(t) = I_p \cos \omega t$$

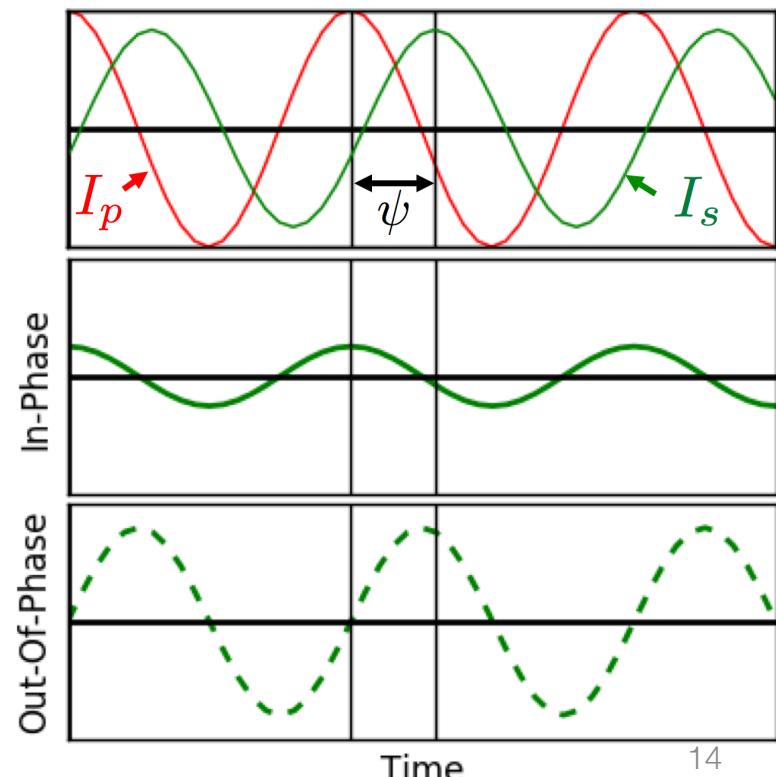
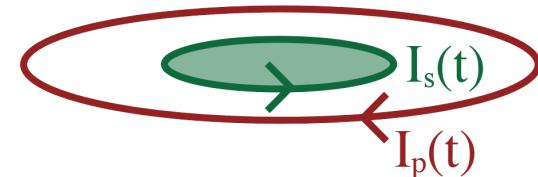
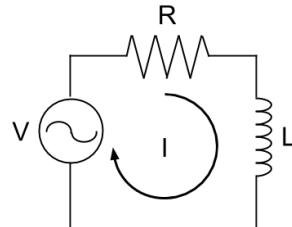
$$I_s(t) = I_s \cos(\omega t - \psi)$$

$$= I_s \cos \psi \cos \omega t + I_s \sin \psi \sin \omega t$$

$\underbrace{}_{\text{In-Phase Real}}$ $\underbrace{}_{\text{Out-of-Phase Quadrature Imaginary}}$

Phase Lag

$$\psi = \frac{\pi}{2} + \tan^{-1} \left(\frac{\omega L}{R} \right)$$



Two Coil Example: Harmonic

Induced Currents

$$I_p(t) = I_p \cos \omega t$$

$$I_s(t) = I_s \cos(\omega t - \psi)$$

$$= I_s \cos \psi \cos \omega t + I_s \sin \psi \sin \omega t$$

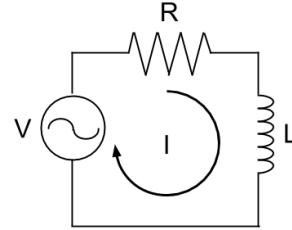
In-Phase
Real

Out-of-Phase
Quadrature
Imaginary

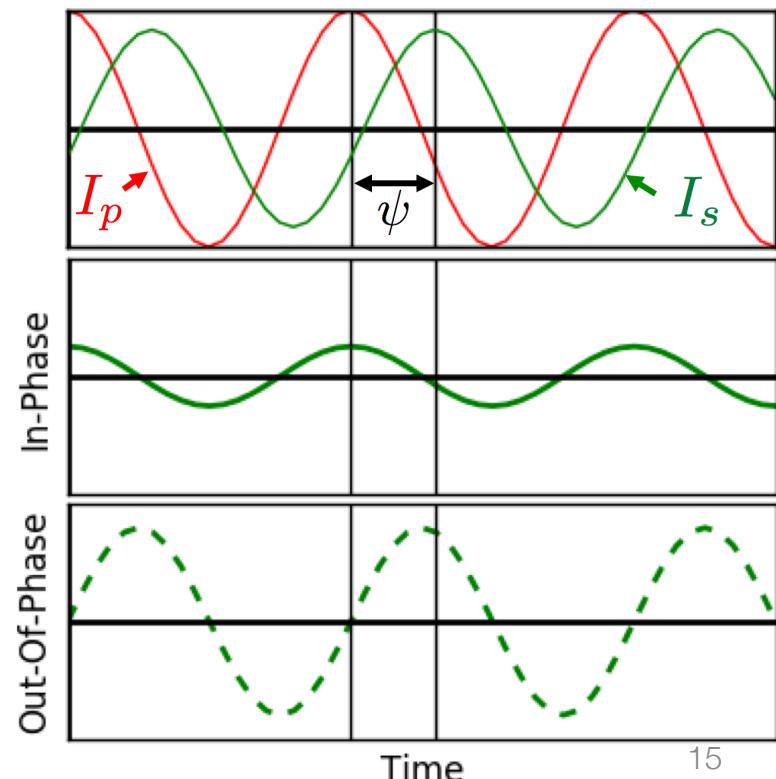
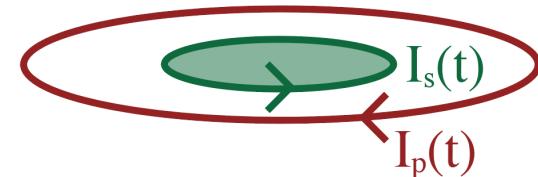
Phase Lag

$$\psi = \frac{\pi}{2} + \tan^{-1} \left(\frac{\omega L}{R} \right)$$

α

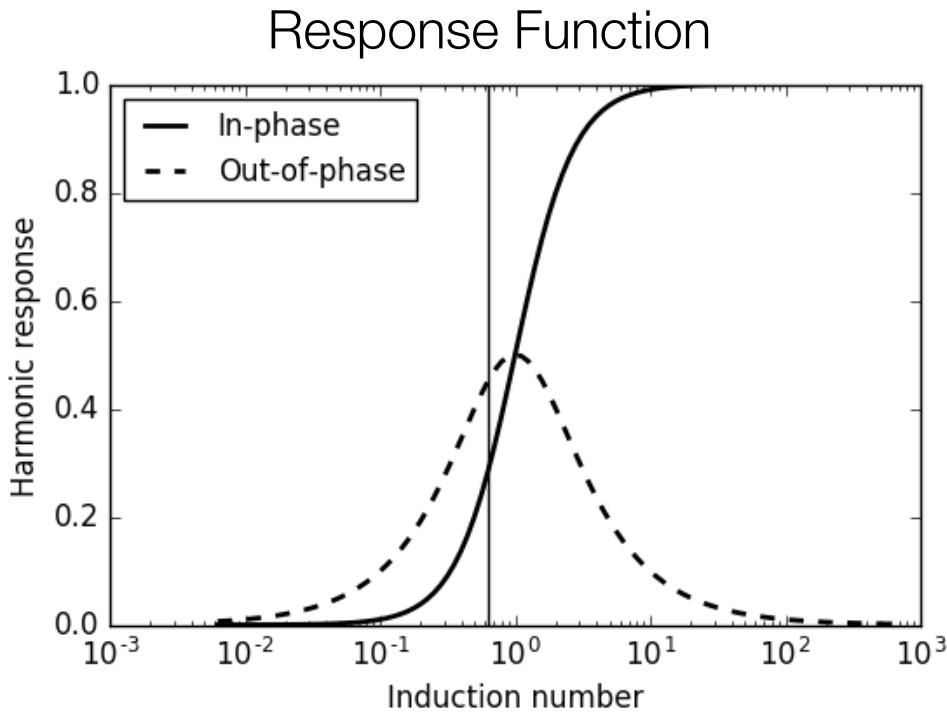
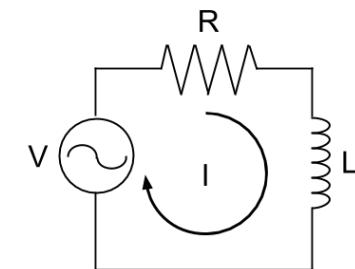


Induction number

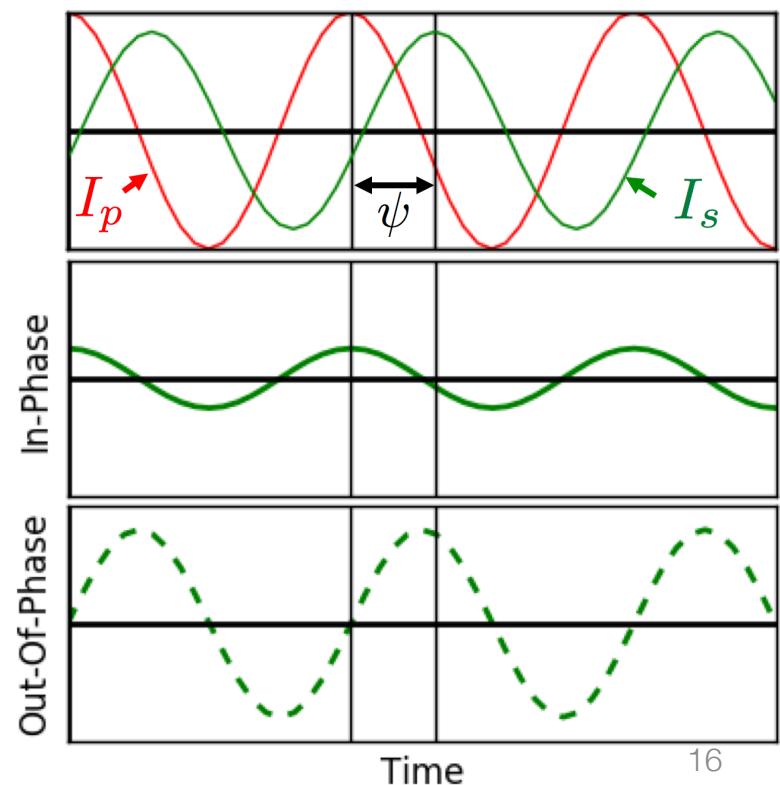


Response Function

- Quantifies how a target responds to a time varying magnetic field
- Partitions real and imaginary parts

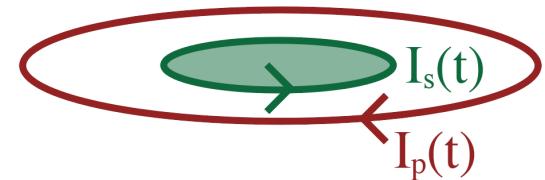
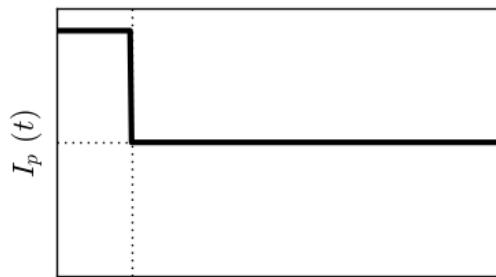


$$\alpha = \frac{\omega L}{R}$$

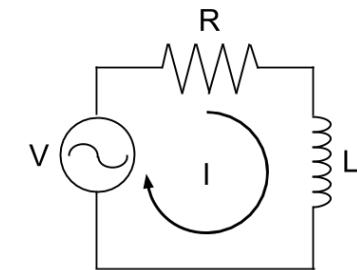
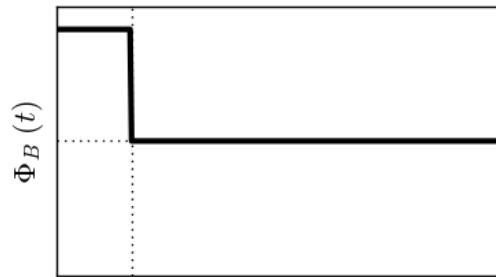


Two Coil Example: Transient

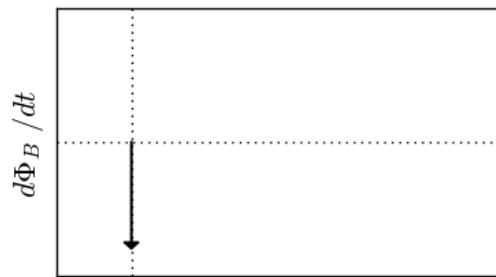
Primary currents



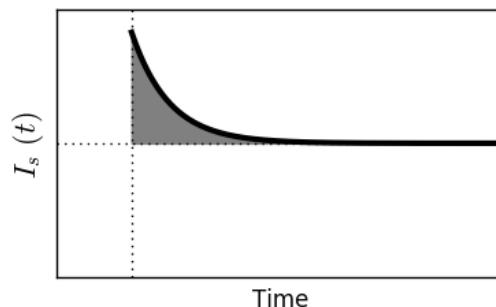
Magnetic flux



Time-variation of magnetic flux



Secondary currents



$$I_s(t) = I_s e^{-t/\tau}$$

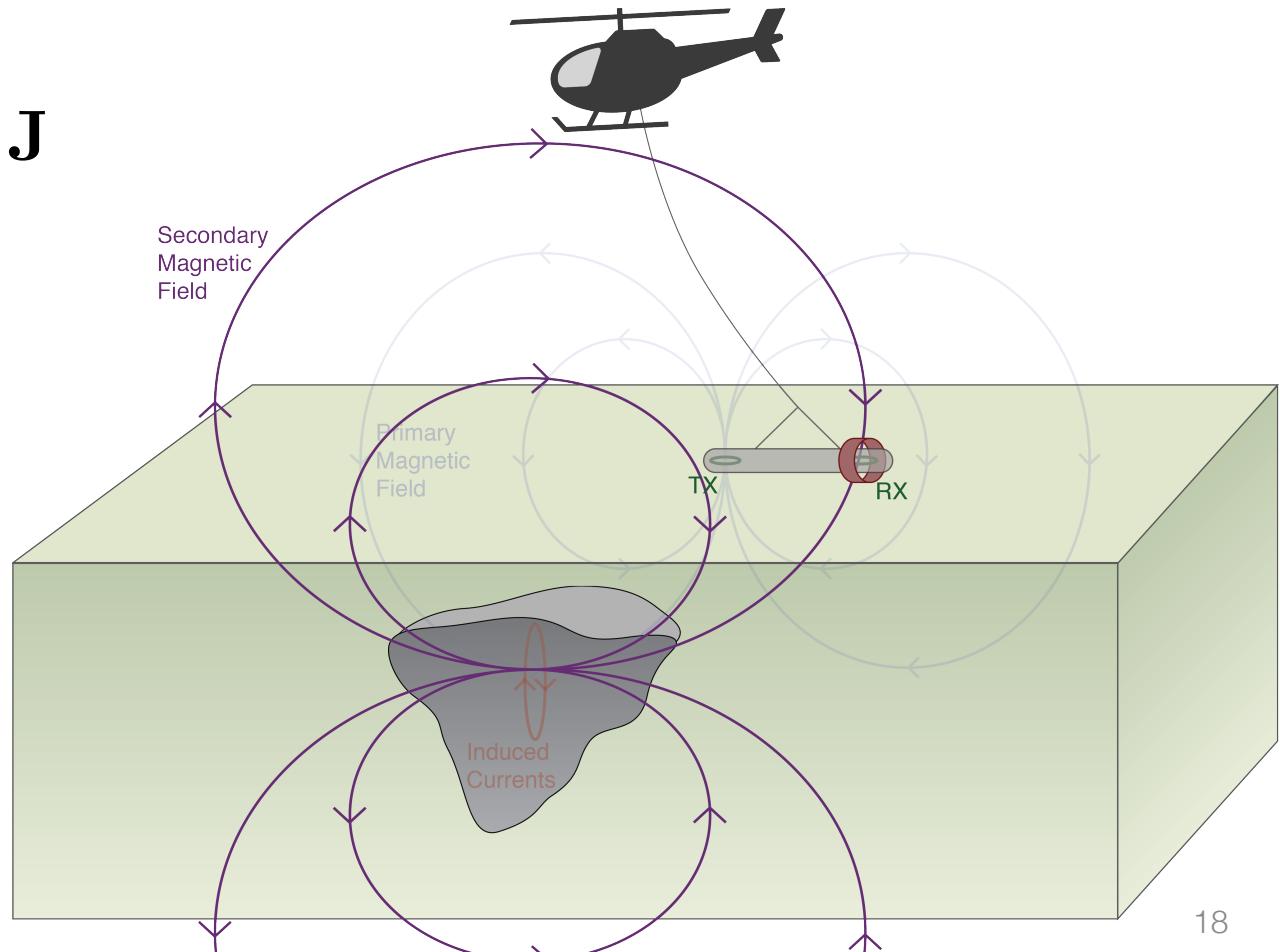
$$\tau = L/R$$

Secondary magnetic fields

Induced currents generate magnetic fields

- Ampere's Law

$$\nabla \times \mathbf{H} = \mathbf{J}$$

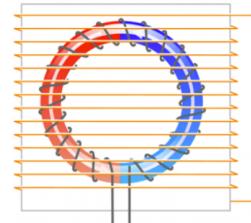


Receiver and Data

Magnetometer

- Measures:
 - Magnetic fields
 - 3 components
- eg. 3-component fluxgate

$$\mathbf{b}(t)$$

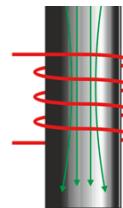


Fluxgate

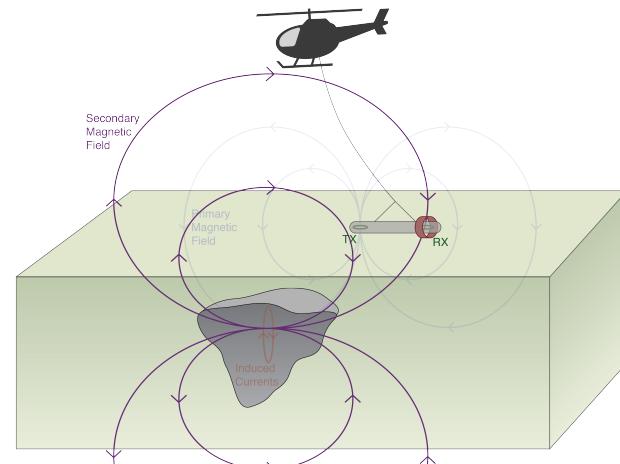
Coil

- Measures:
 - Voltage
 - Single component that depends on coil orientation
 - Coupling matters
- eg. airborne frequency domain
 - ratio of H_s/H_p is the same as V_s/V_p

$$\frac{\partial \mathbf{b}}{\partial t}$$



Coil



Coupling

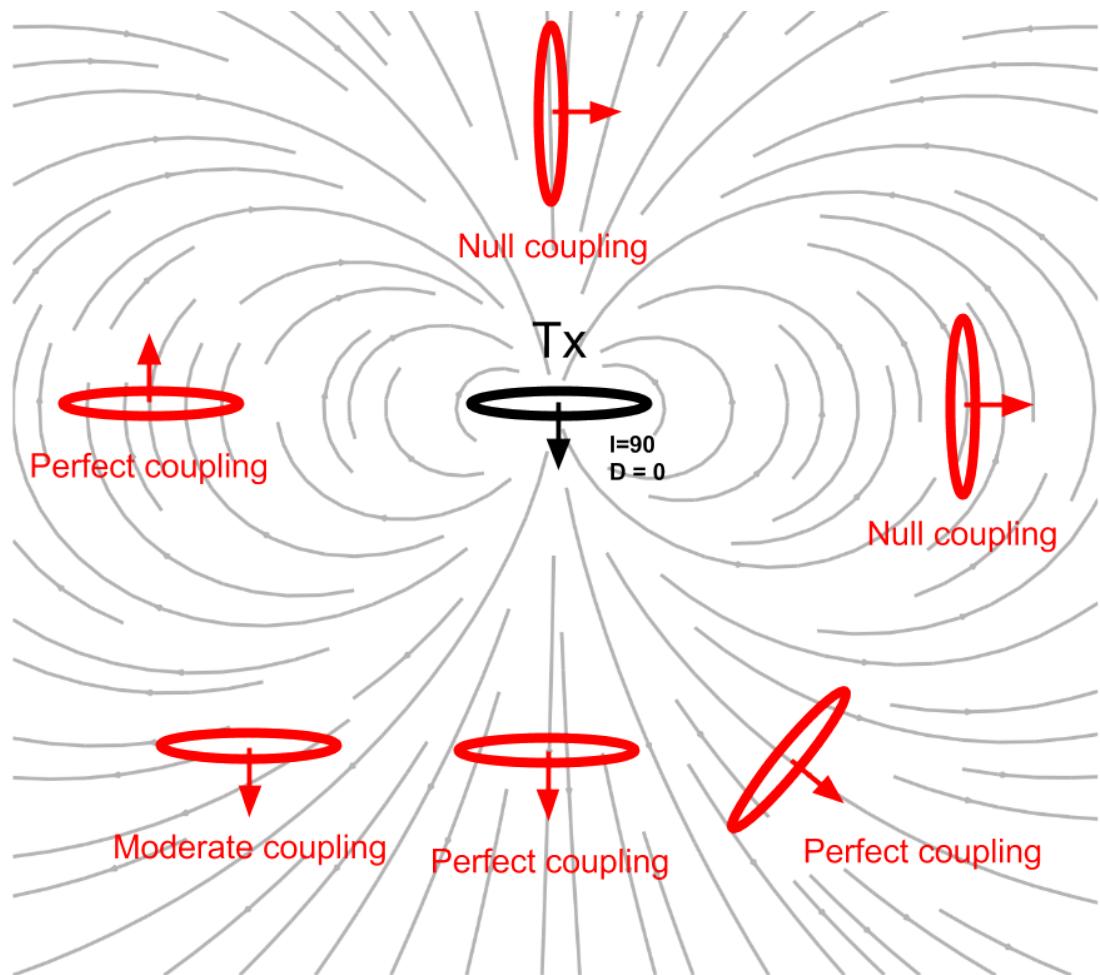
- Transmitter: Primary

$$I_p(t) = I_p \cos(\omega t)$$

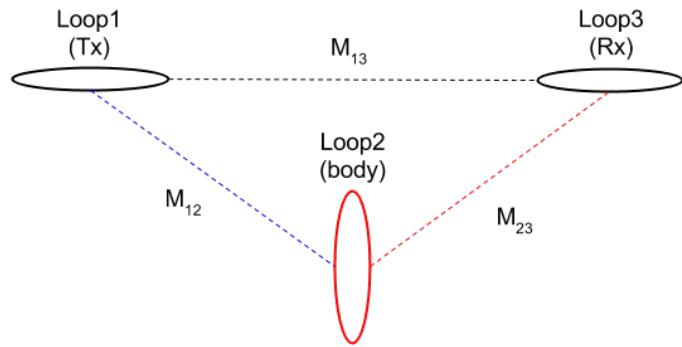
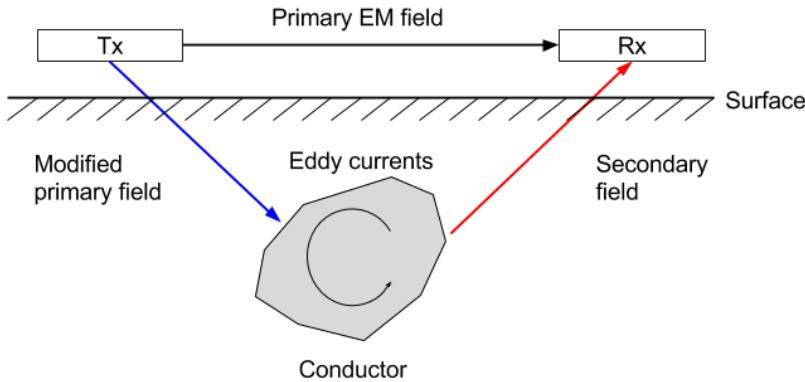
$$\mathbf{B}_p(t) \sim I_p \cos(\omega t)$$

- Target: Secondary

$$\begin{aligned} EMF &= -\frac{\partial \phi_{\mathbf{B}}}{\partial t} \\ &= -\frac{\partial}{\partial t} (\mathbf{B}_p \cdot \hat{\mathbf{n}}) A \end{aligned}$$



Circuit model of EM induction



Coupling coefficient

- Depends on geometry

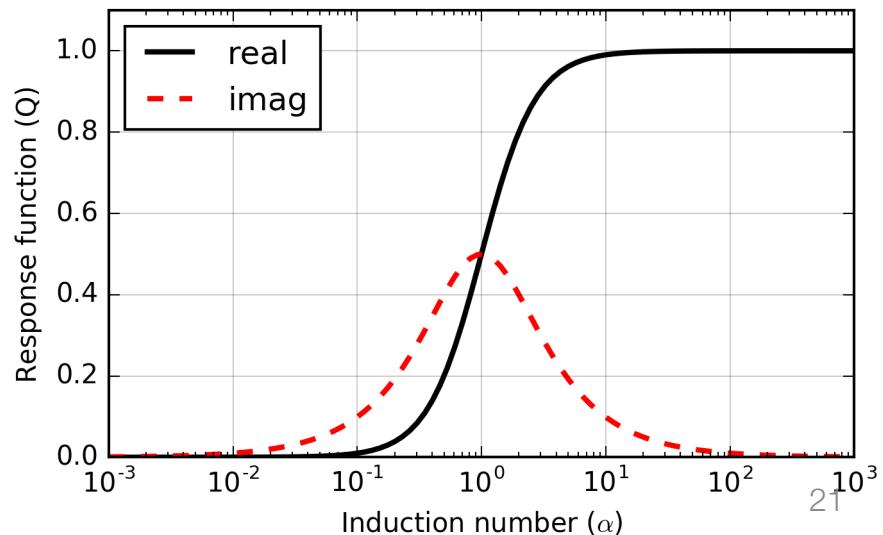
$$M_{12} = \frac{\mu_0}{4\pi} \oint \oint \frac{dl_1 \cdot dl_2}{|\mathbf{r} - \mathbf{r}'|^2}.$$

Magnetic field at the receiver

$$\frac{H^s}{H^p} = - \frac{M_{12}M_{23}}{M_{13}L} \underbrace{\left[\frac{\alpha^2 + i\alpha}{1 + \alpha^2} \right]}_Q$$

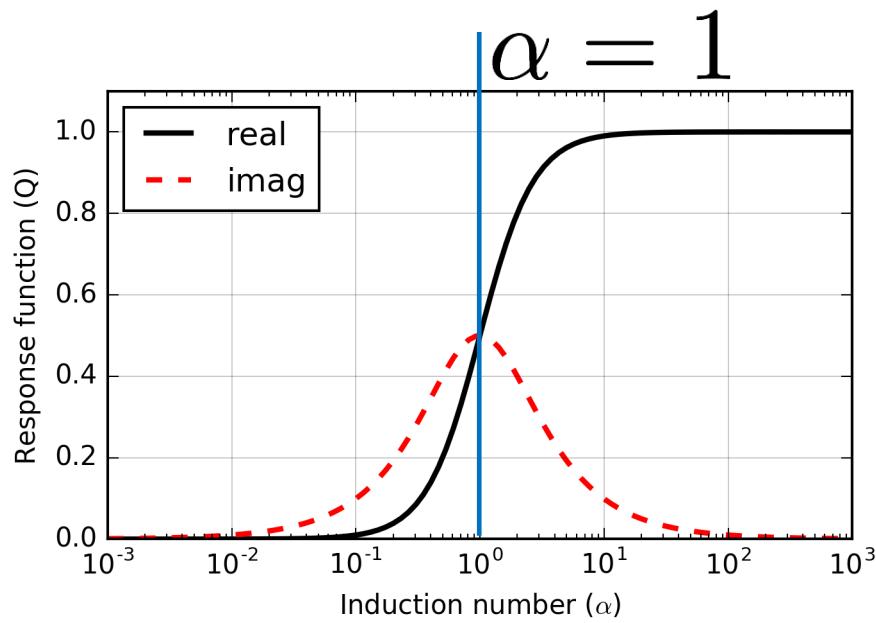
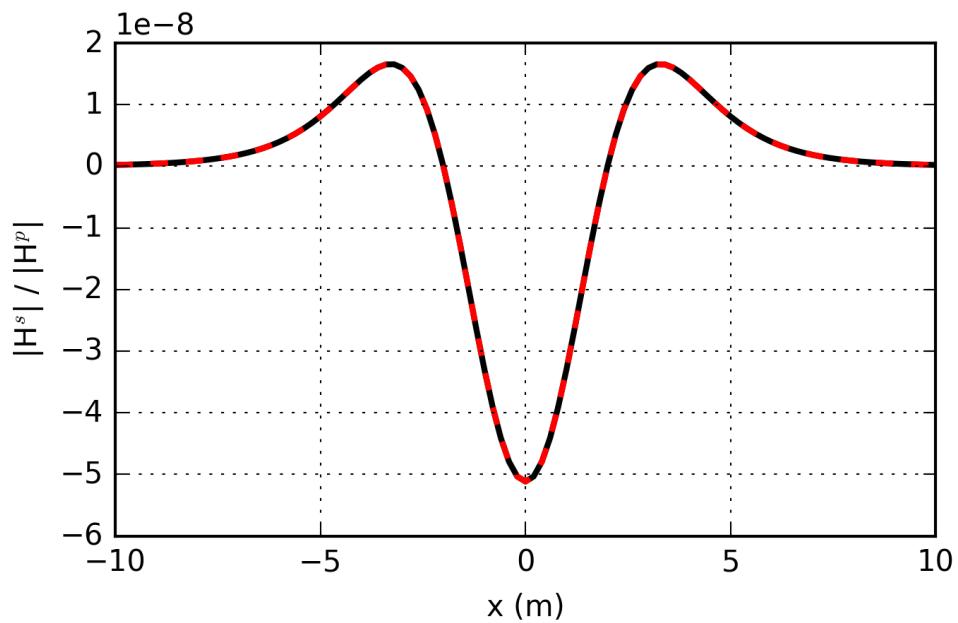
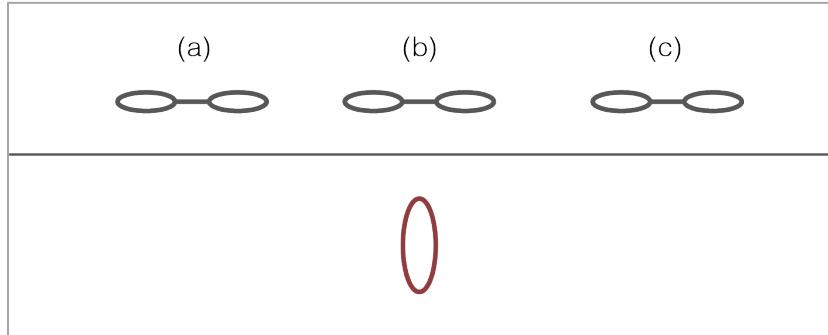
Induction Number

- Depends on properties of target $\alpha = \frac{\omega L}{R}$



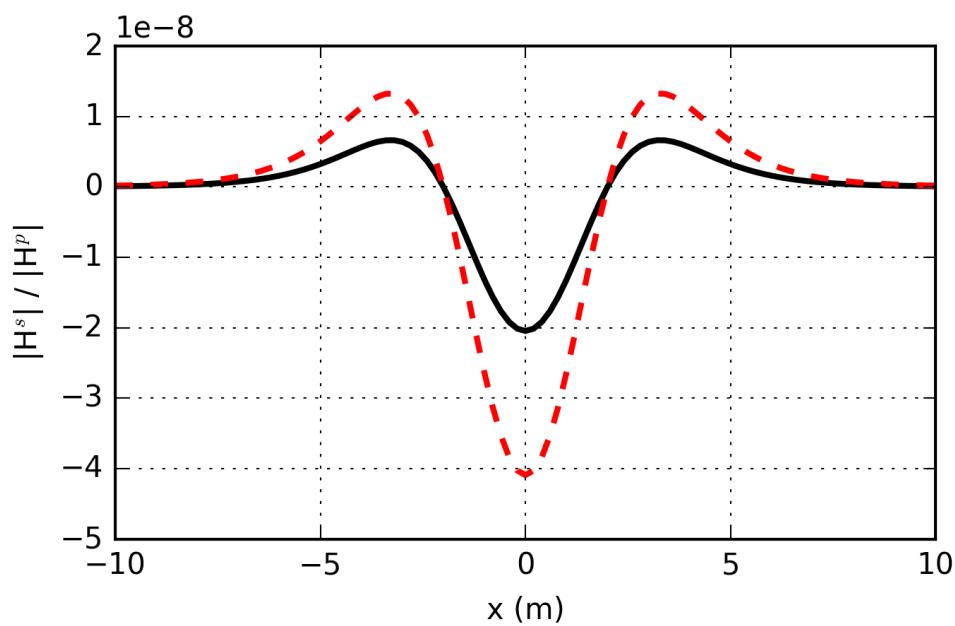
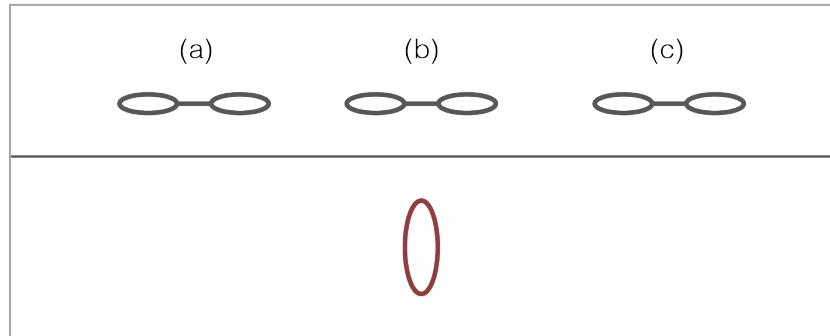
Conductor in a resistive earth: Frequency

Profile over the loop

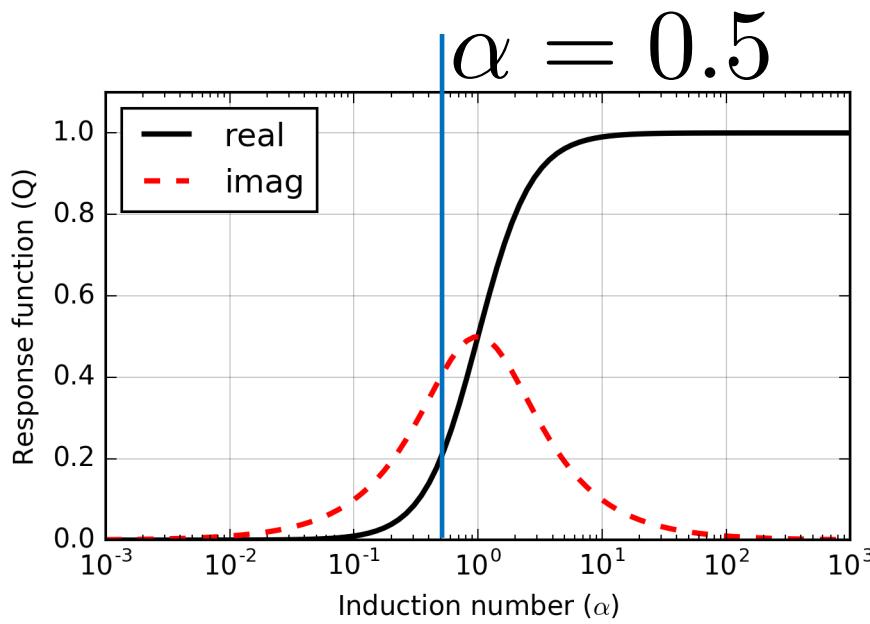


Conductor in a resistive earth: Frequency

Profile over the loop

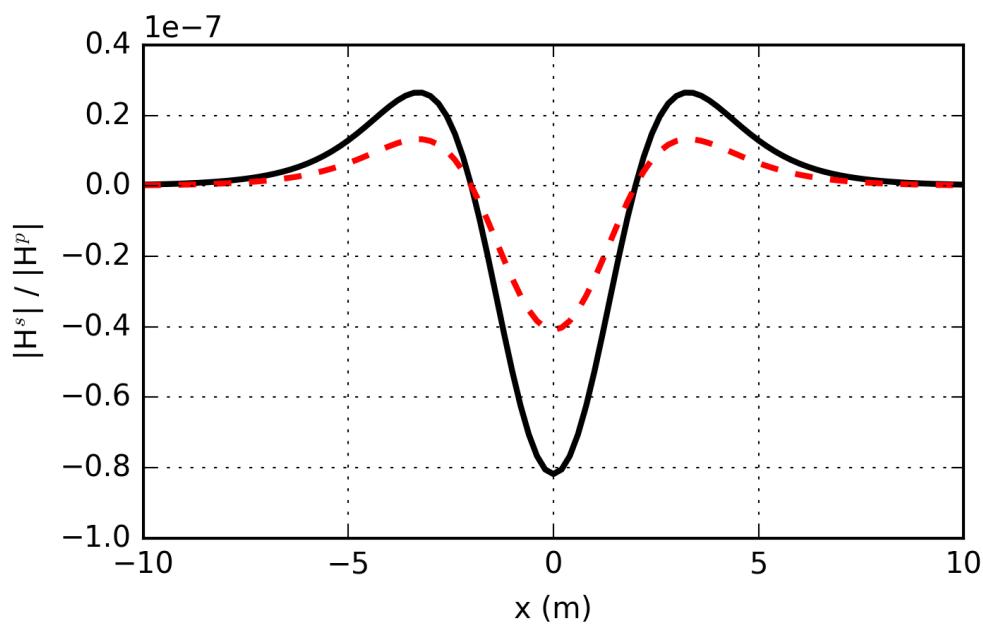
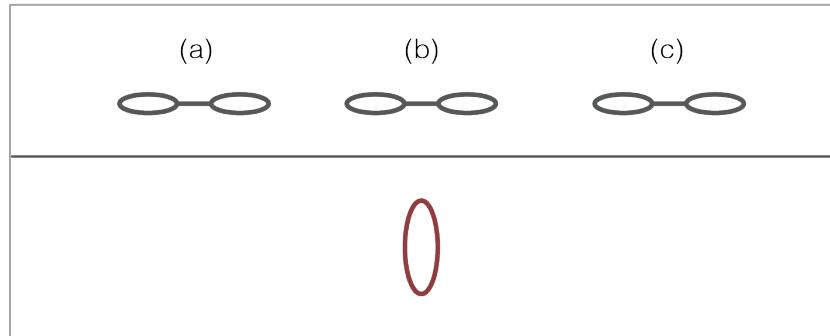


- Induction number
$$\alpha = \frac{\omega L}{R}$$
- When $\alpha < 1$
 - Real < Imag

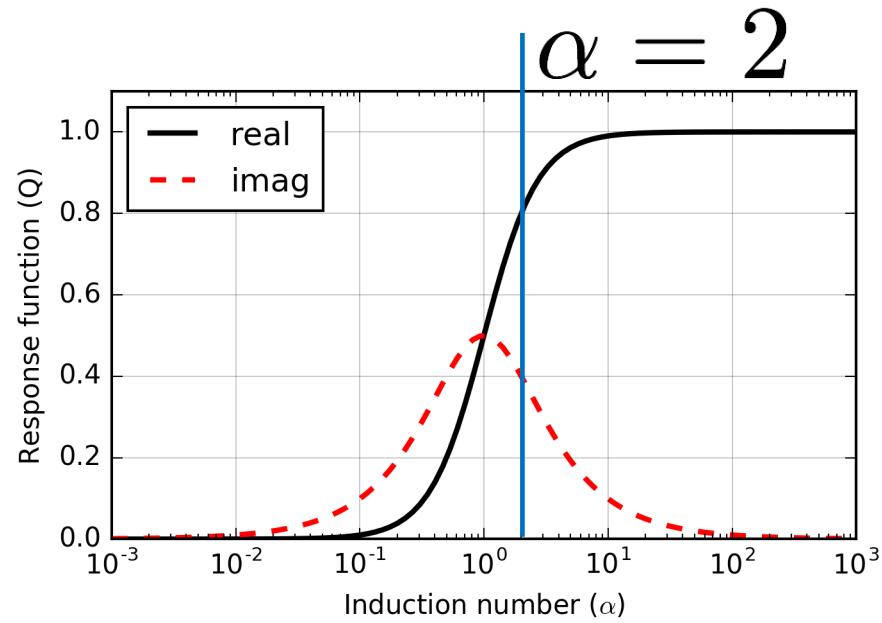


Conductor in a resistive earth: Frequency

Profile over the loop

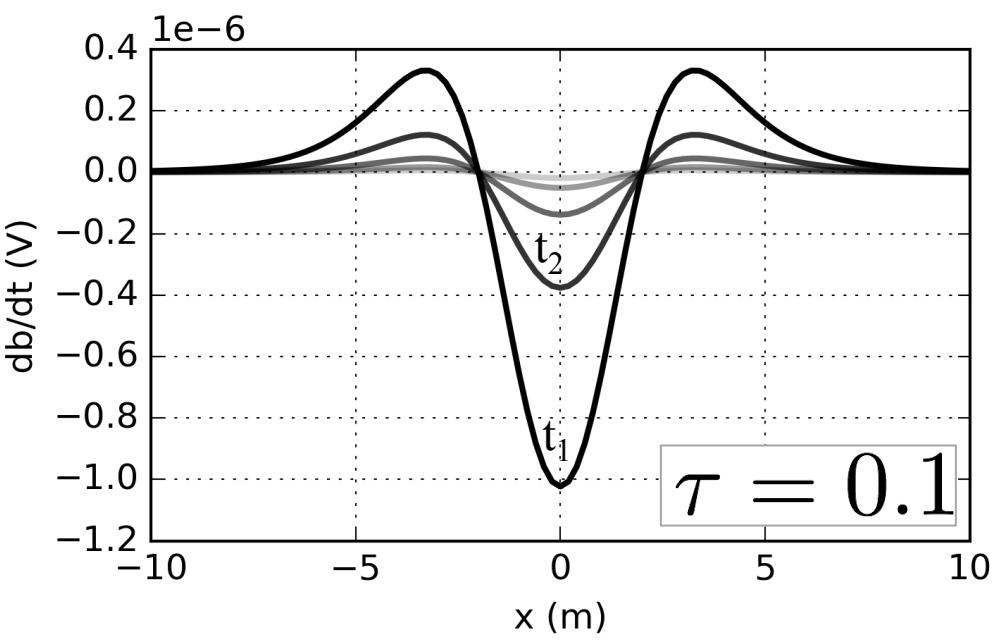
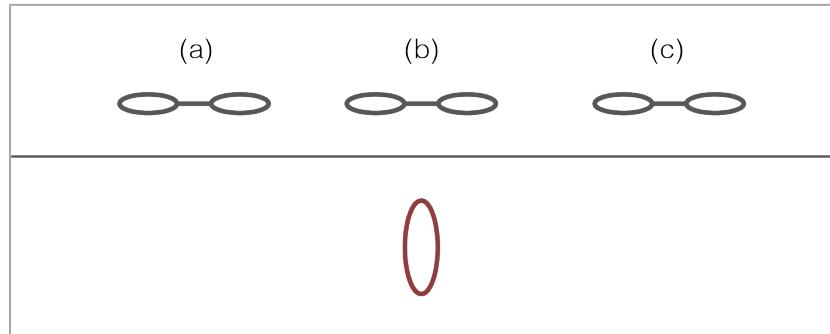


- Induction number
$$\alpha = \frac{\omega L}{R}$$
- When $\alpha > 1$
 - Real > Imag

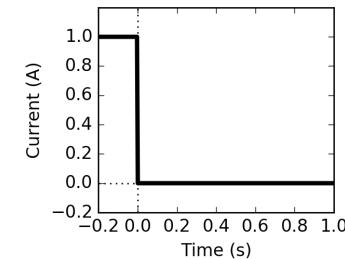


Conductor in a resistive earth: Transient

Profile over the loop

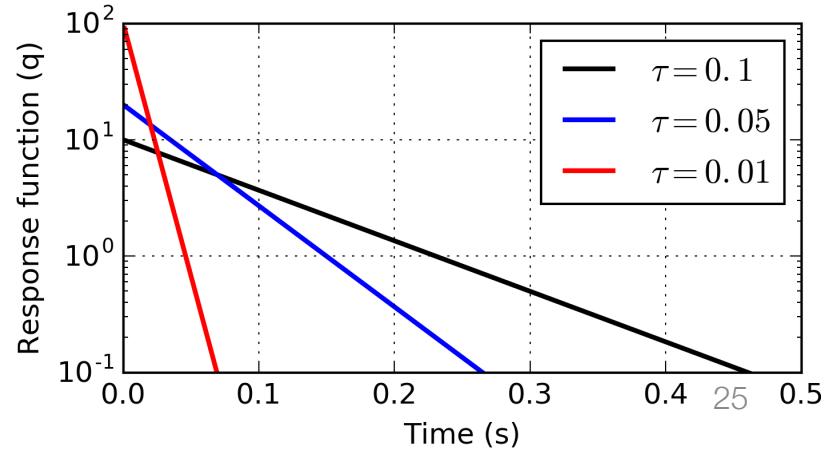


- Time constant
$$\tau = L/R$$
- Step-off current in Tx



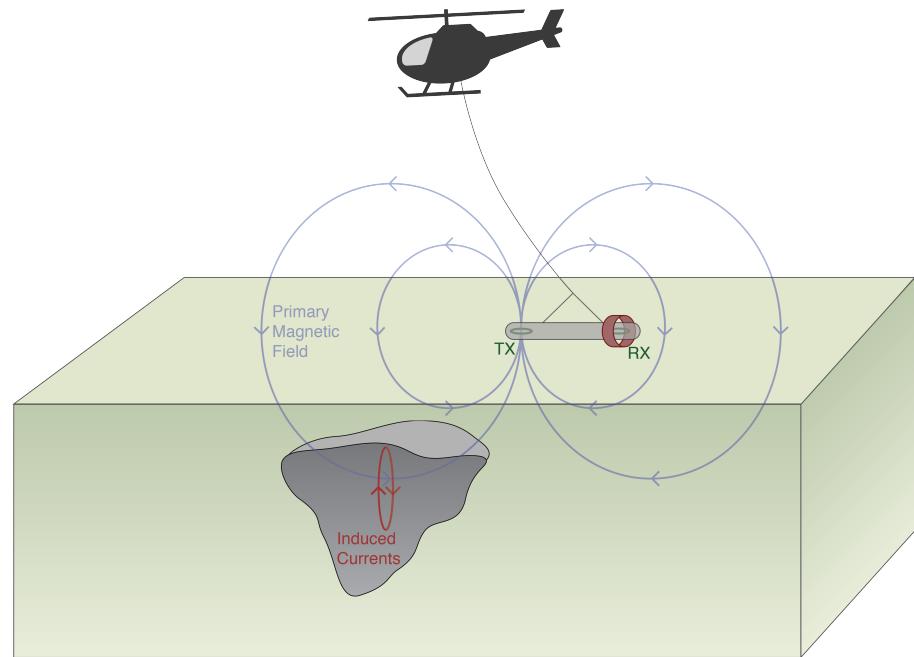
- Response function depends on time, τ

$$q(t) = e^{-t/\tau}$$



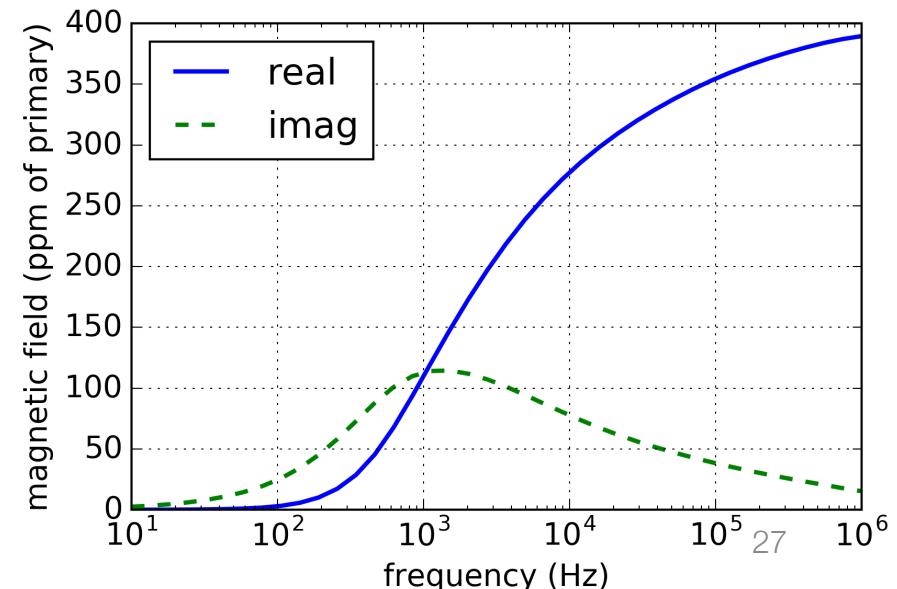
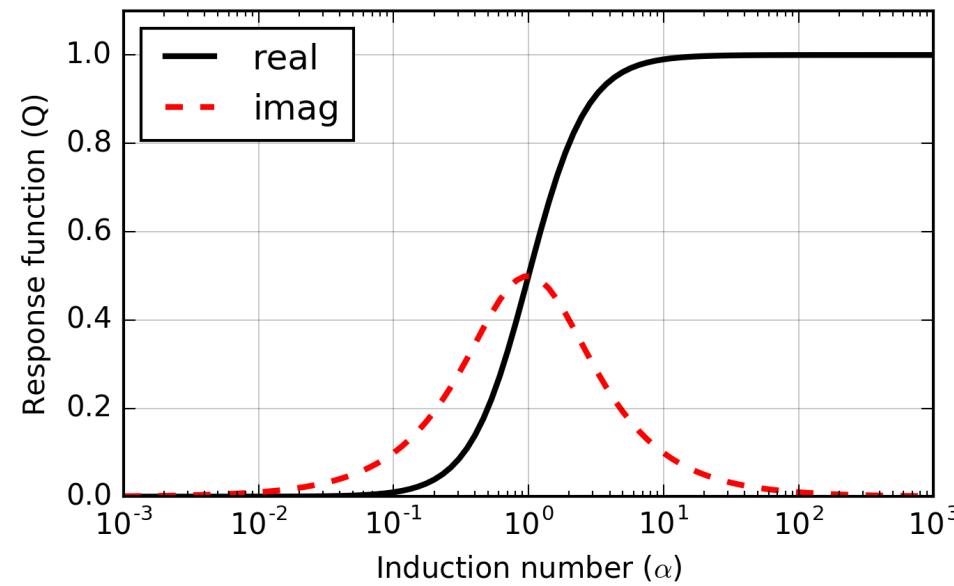
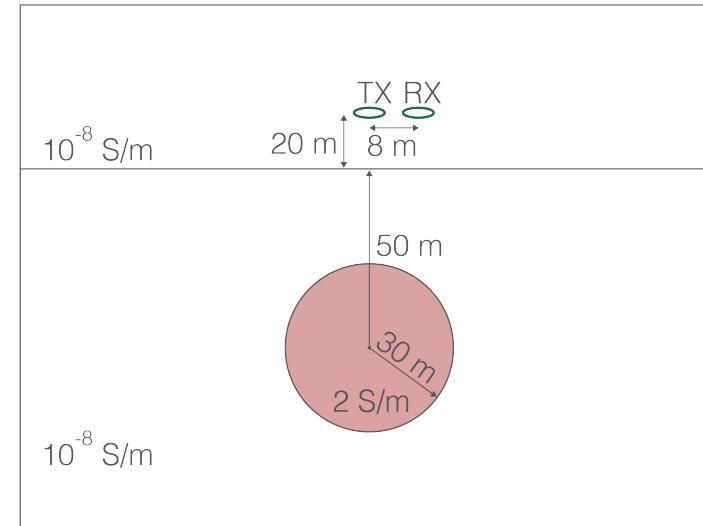
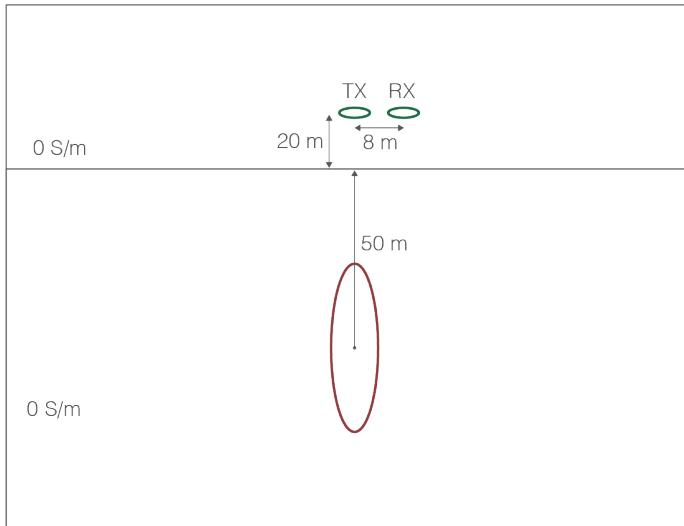
Recap: what have we learned?

- Basics of EM induction
- Response functions
- Mutual coupling
- Data for frequency or time domain systems
- Circuit model provides representative results
 - Applicable to geologic targets?



Sphere in a resistive background

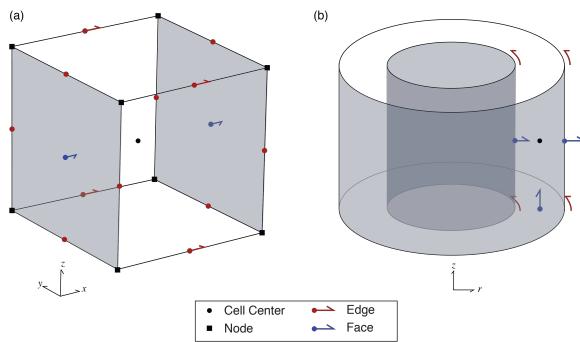
How representative is a circuit model?



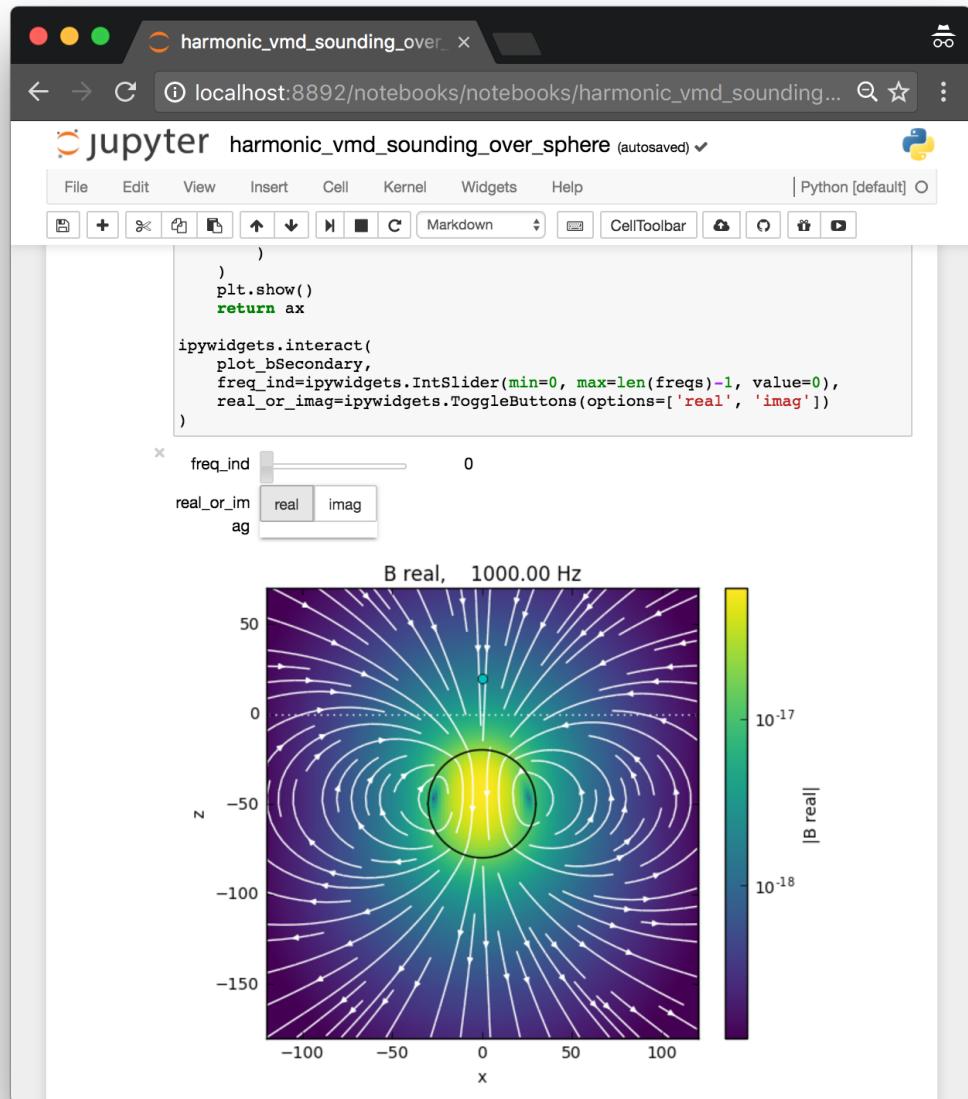
Cyl Code



- Finite Volume EM
 - Frequency and Time



- Built on SimPEG
- Open source, available at:
<http://em.geosci.xyz/apps.html>

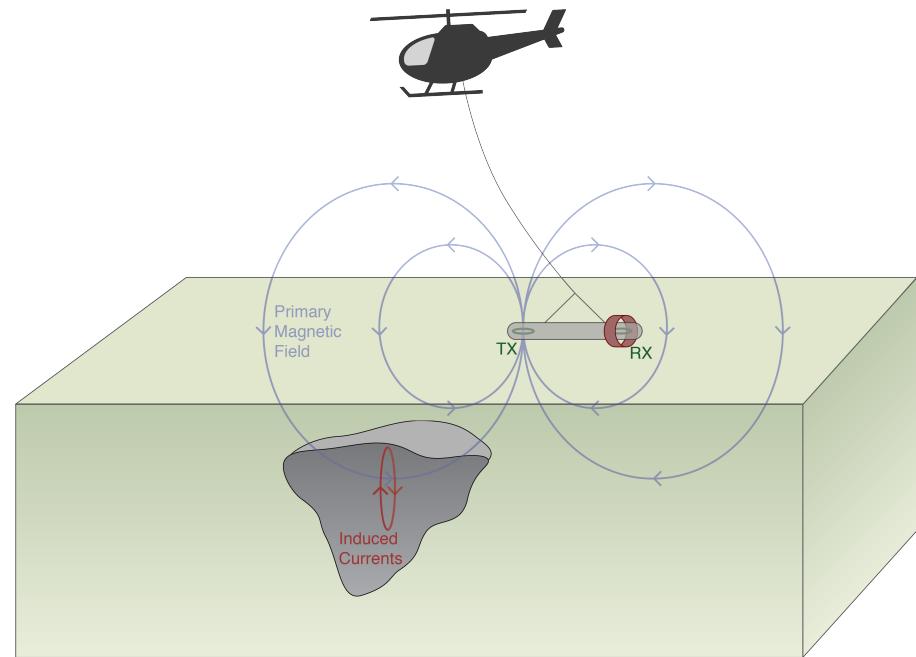


Recap: what have we learned?

- Basics of EM induction
- Response functions
- Mutual coupling
- Data for frequency or time domain systems
- Circuit model is a good proxy

Major item not yet accounted for...

- Propagation of energy from
 - Transmitter to target
 - Target to receiver



How do EM fields and fluxes behave in a conductive background?

Revisit Maxwell's equations

First order equations

$$\nabla \times \mathbf{e} = -\frac{\partial \mathbf{b}}{\partial t} \quad \mathbf{j} = \sigma \mathbf{e}$$
$$\nabla \times \mathbf{h} = \mathbf{j} + \frac{\partial \mathbf{d}}{\partial t} \quad \mathbf{b} = \mu \mathbf{h}$$
$$-\nabla \cdot \mathbf{d} = \epsilon \mathbf{e} \quad \mathbf{d} = \epsilon \mathbf{e}$$

Second order equations

$$\nabla^2 \mathbf{h} - \underbrace{\mu\sigma \frac{\partial \mathbf{h}}{\partial t}}_{\text{diffusion}} - \underbrace{\mu\epsilon \frac{\partial^2 \mathbf{h}}{\partial t^2}}_{\text{wave propagation}} = 0$$

In frequency

$$\nabla^2 \mathbf{H} + k^2 \mathbf{H} = 0$$

$$k^2 = \omega^2 \mu \epsilon - i \omega \mu \sigma$$

* Same equation holds[†] for E

Plane waves in a homogeneous media

In frequency

$$\nabla^2 \mathbf{H} + k^2 \mathbf{H} = 0$$

$$k^2 = \omega^2 \mu \epsilon - i \omega \mu \sigma$$

Quasi-static

$$\frac{\omega \epsilon}{\sigma} \ll 1$$

even if...

$$\sigma = 10^{-4} S/m$$

$$f = 10^4 Hz$$

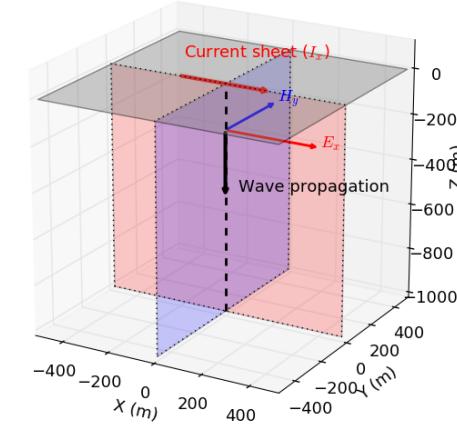
then

$$\frac{\omega \epsilon}{\sigma} \sim 0.005$$

$$k = \sqrt{-i \omega \mu \sigma} = (1 - i) \sqrt{\frac{\omega \mu \sigma}{2}}$$

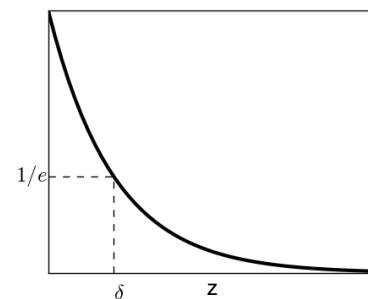
$$\equiv \alpha - i\beta$$

Plane wave solution



$$\mathbf{H} = \mathbf{H}_0 e^{\underbrace{-\alpha z}_{\text{attenuation}}} e^{\underbrace{-i(\beta z - \omega t)}_{\text{phase}}}$$

Skin depth



δ : skin depth

$$\delta = \sqrt{\frac{2}{\omega \mu \sigma}} = 503 \sqrt{\frac{1}{\sigma f}}$$

Plane waves in a homogeneous media

In time

$$\nabla^2 \mathbf{h} - \mu\epsilon \frac{\partial^2 \mathbf{h}}{\partial t^2} - \mu\sigma \frac{\partial \mathbf{h}}{\partial t} = 0$$

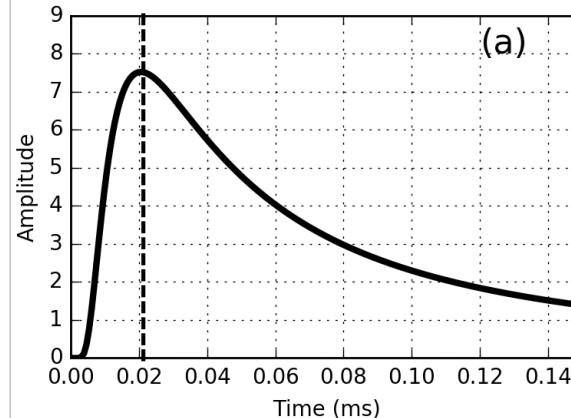
$$\mathbf{h}(t = 0) = \mathbf{h}_0 \delta(t)$$

Solution for quasi-static

$$\mathbf{h}(t) = -\frac{(\mu\sigma)^{1/2} z}{2\pi^{1/2} t^{3/2}} e^{-\mu\sigma z^2/(4t)}$$

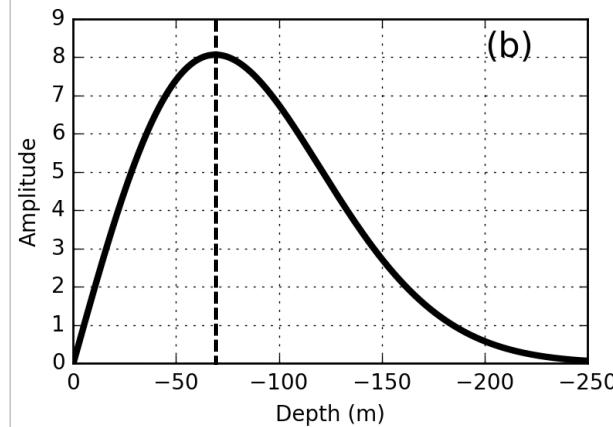
z : depth (m)

Peak time:



$$t_{max} = \frac{\mu\sigma z^2}{6}$$

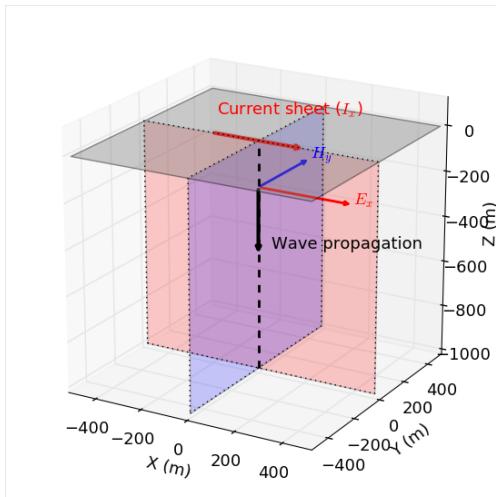
Diffusion distance



$$d = \sqrt{\frac{2t}{\mu\sigma}}$$
$$\approx 1260 \sqrt{\frac{t}{\sigma}}$$

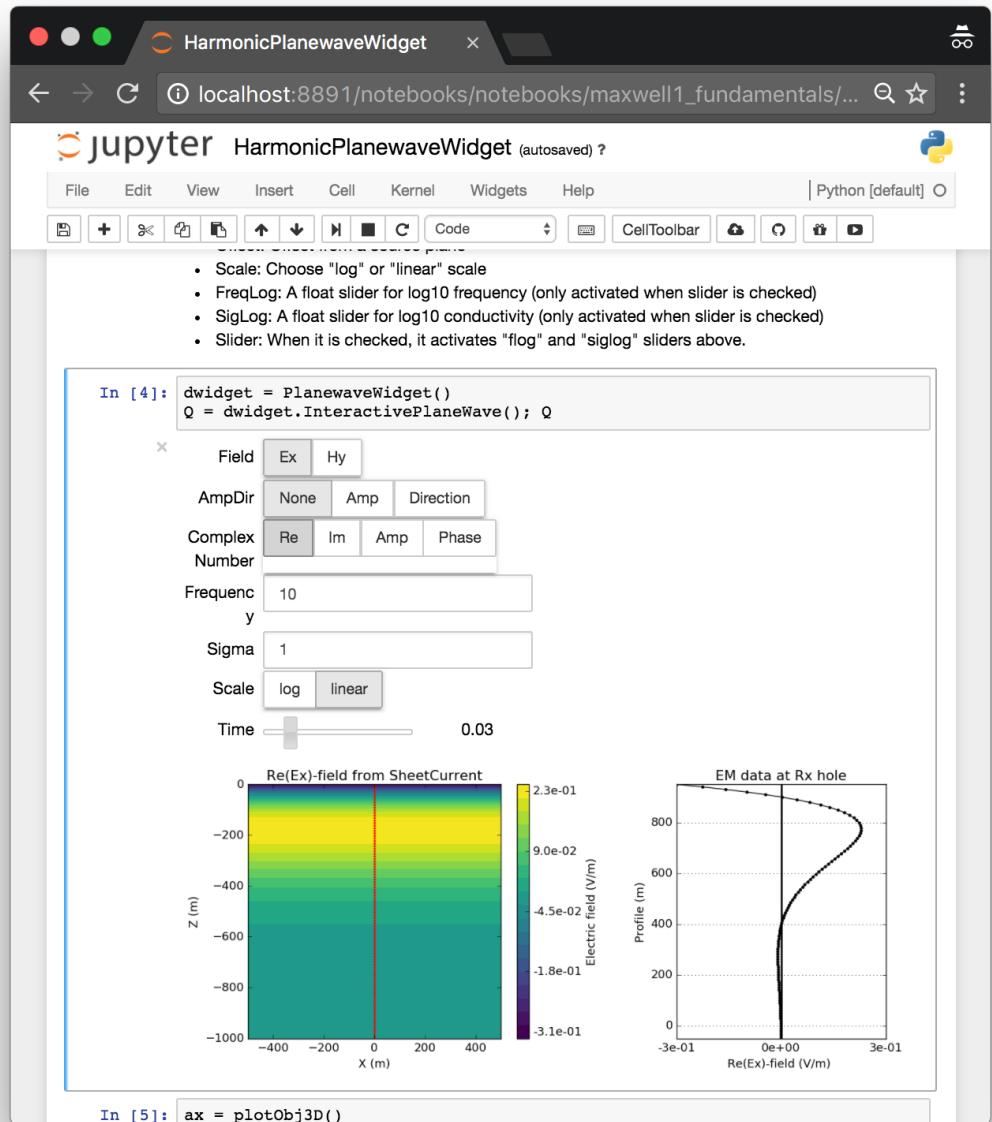
Frequency Domain App: Plane waves

- Plane wave

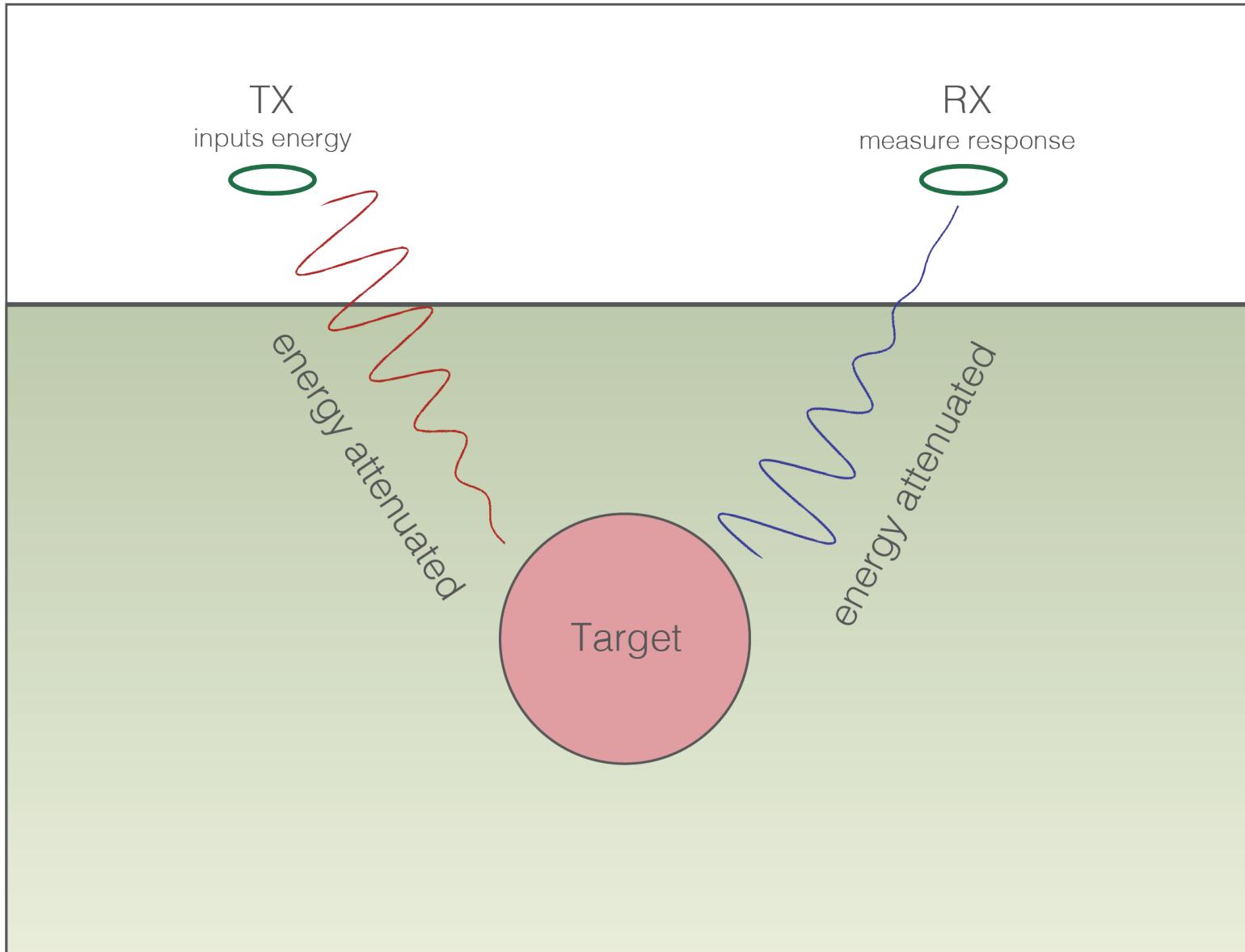


$$\mathbf{H} = \mathbf{H}_0 e^{-\alpha z} e^{-i(\beta z - \omega t)}$$

attenuation phase

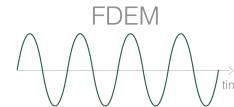


Effects of background resistivity

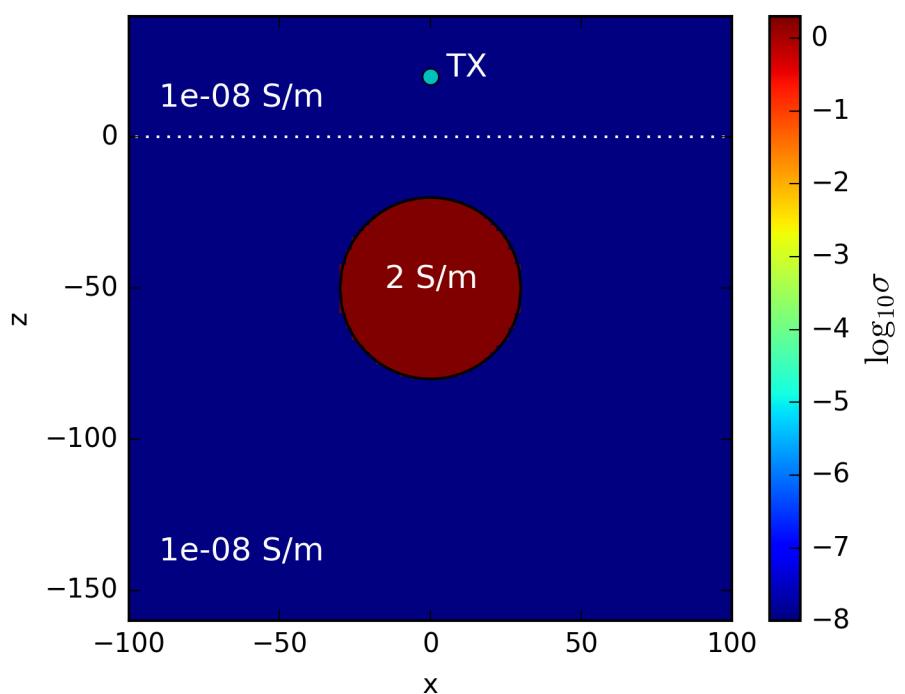


Effects of background resistivity: Frequency

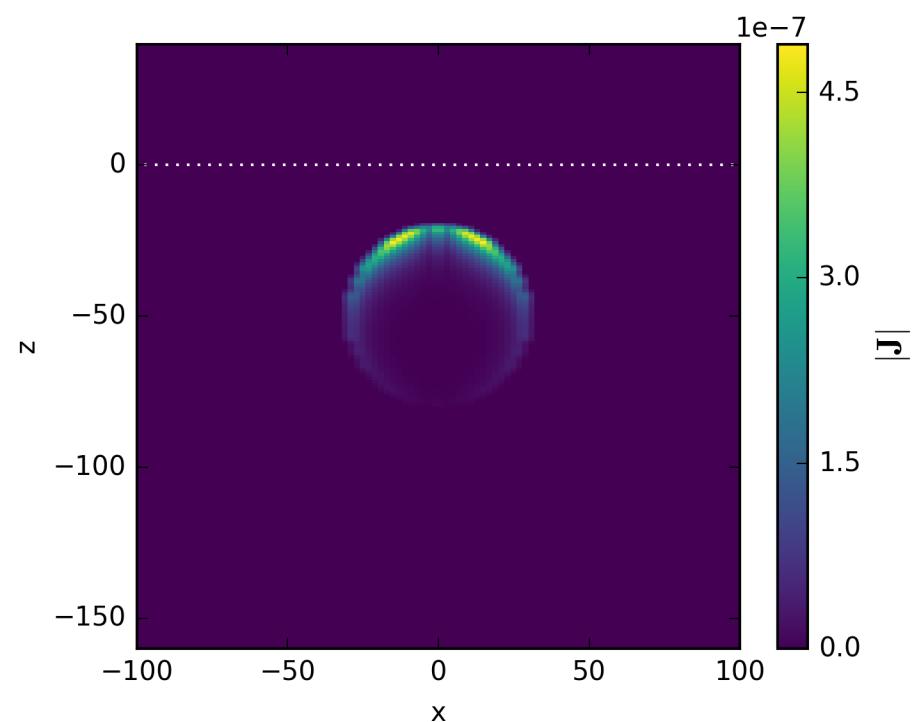
- Buried, conductive sphere
- Vary background conductivity
- Frequency: 10^4 Hz



10^{-8} S/m background

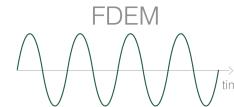


Current Density

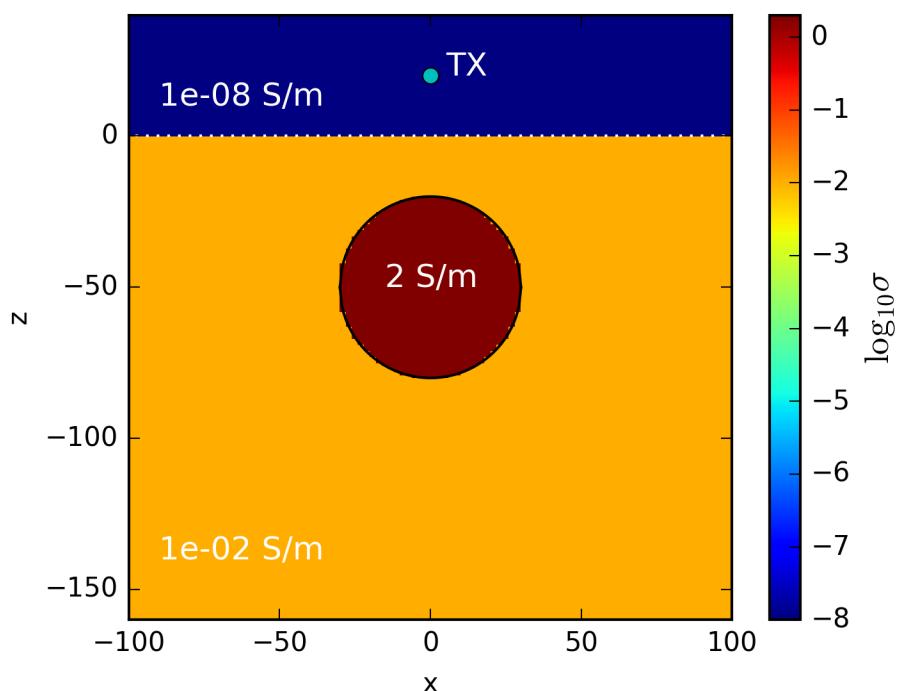


Effects of background resistivity: Frequency

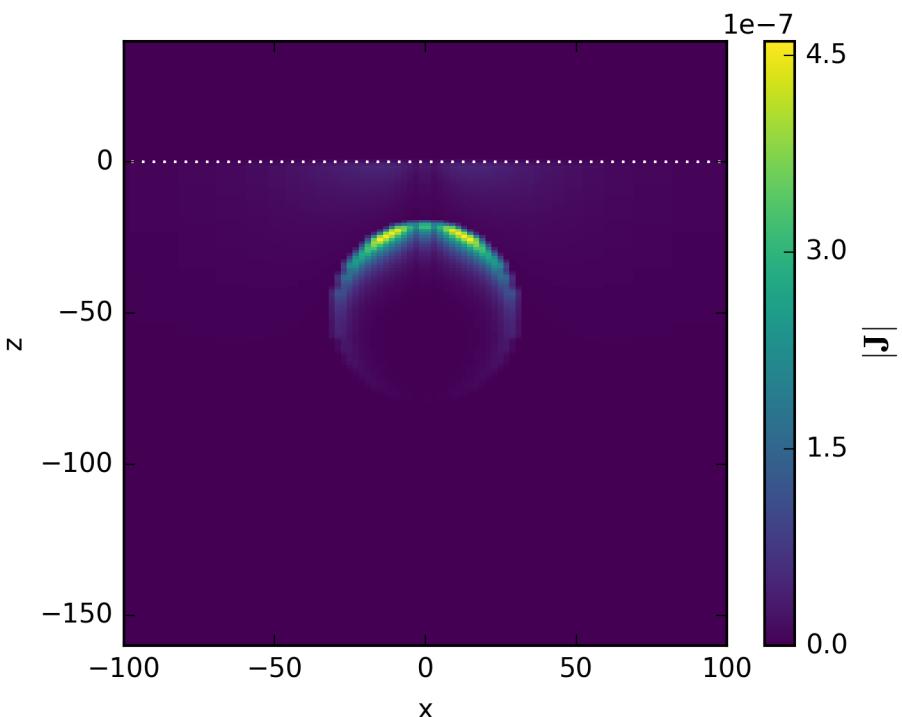
- Buried, conductive sphere
- Vary background conductivity
- Frequency: 10^4 Hz



10^{-2} S/m background

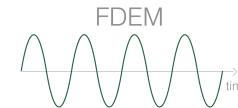


Current Density

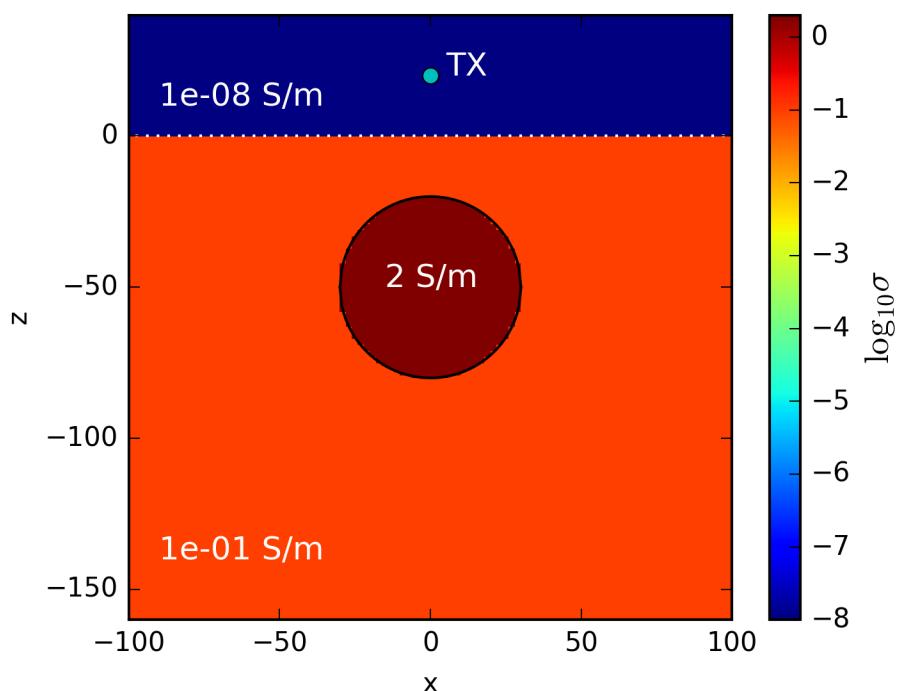


Effects of background resistivity: Frequency

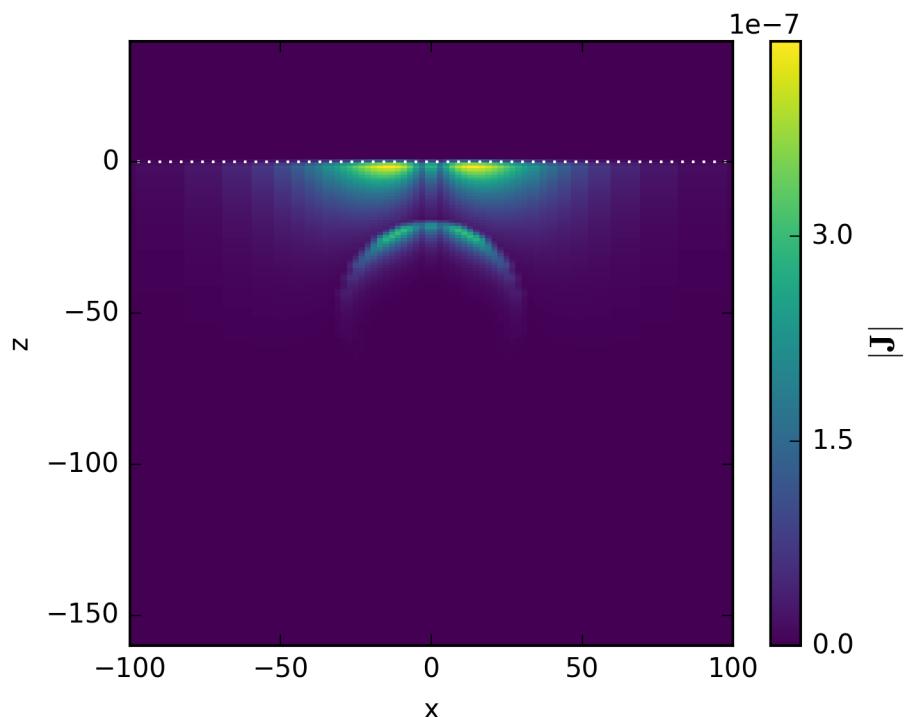
- Buried, conductive sphere
- Vary background conductivity
- Frequency: 10^4 Hz



10^{-1} S/m background



Current Density

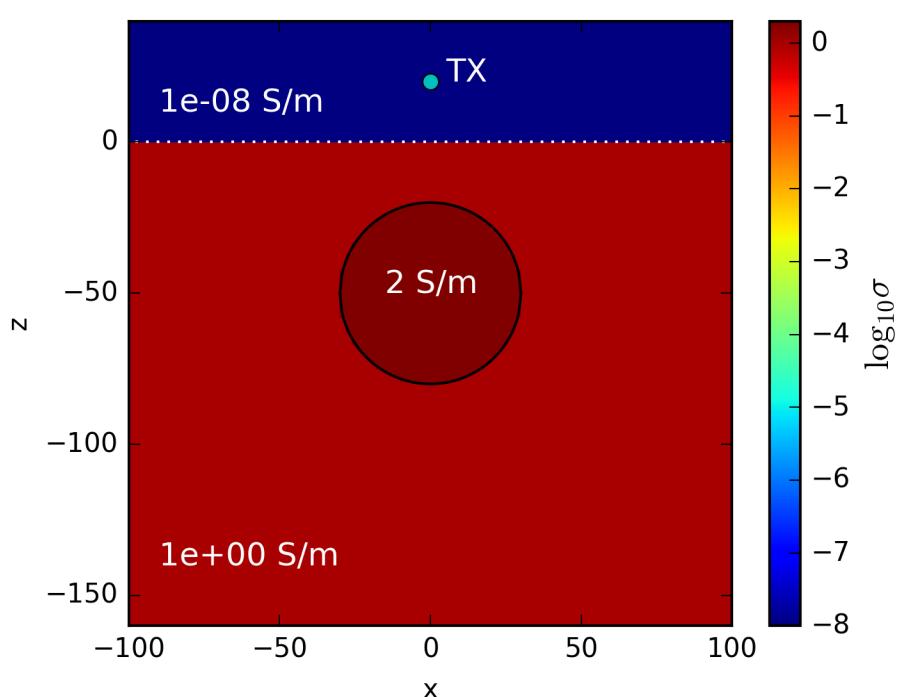


Effects of background resistivity: Frequency

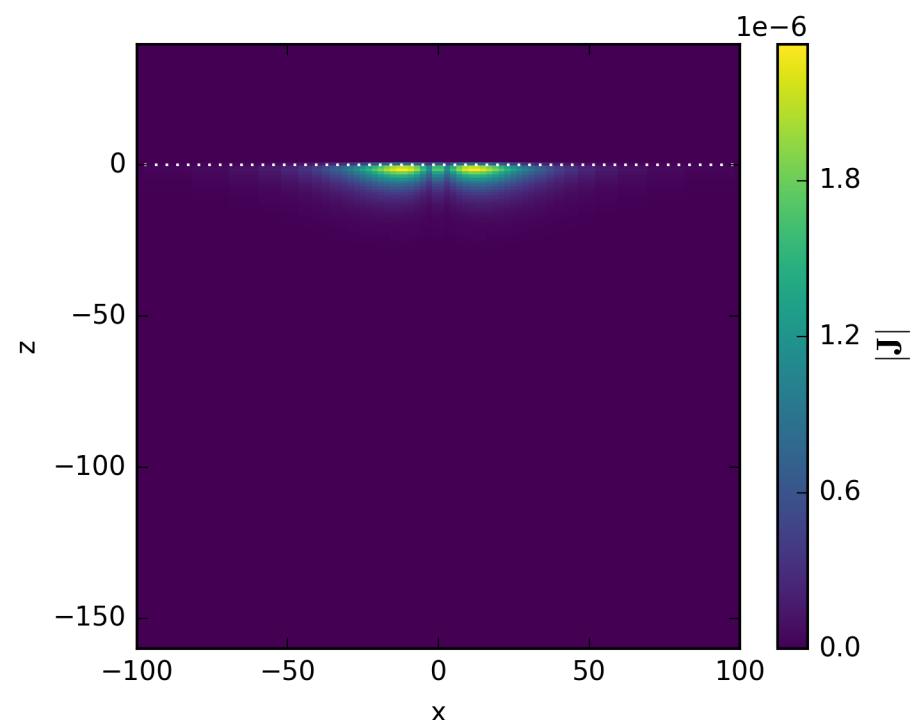
- Buried, conductive sphere
- Vary background conductivity
- Frequency: 10^4 Hz



1 S/m background

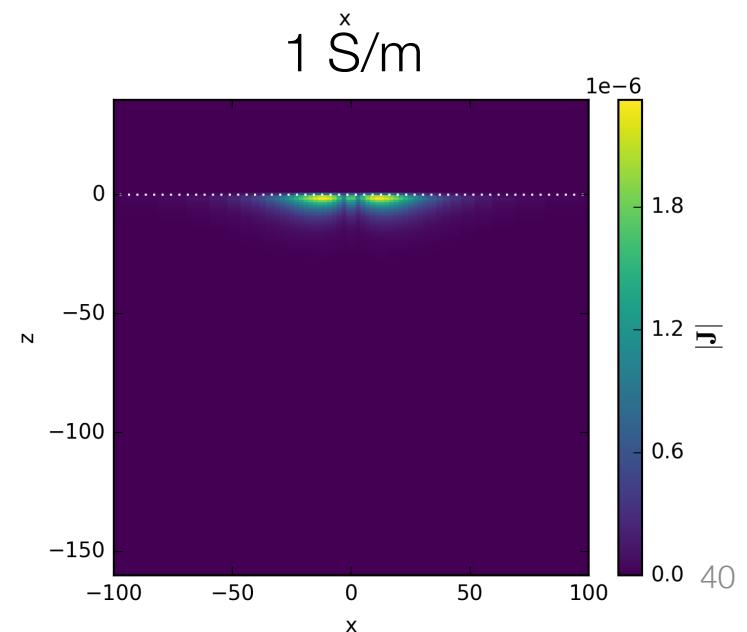
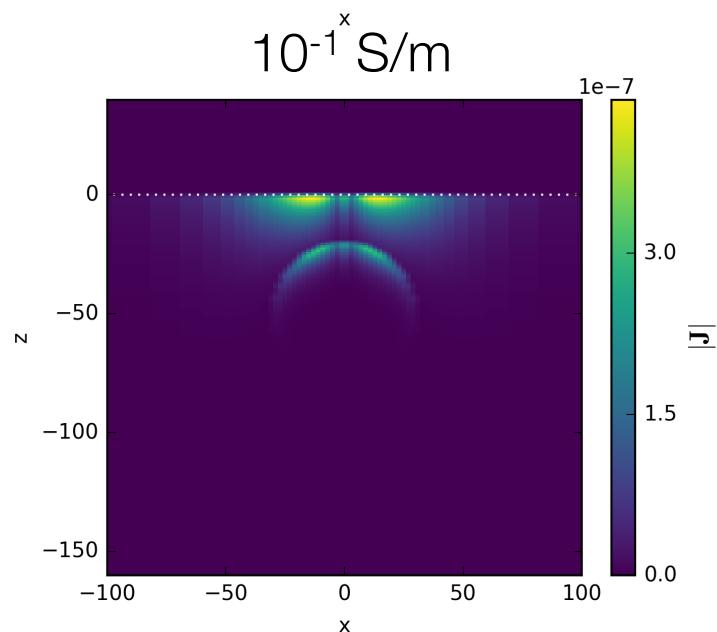
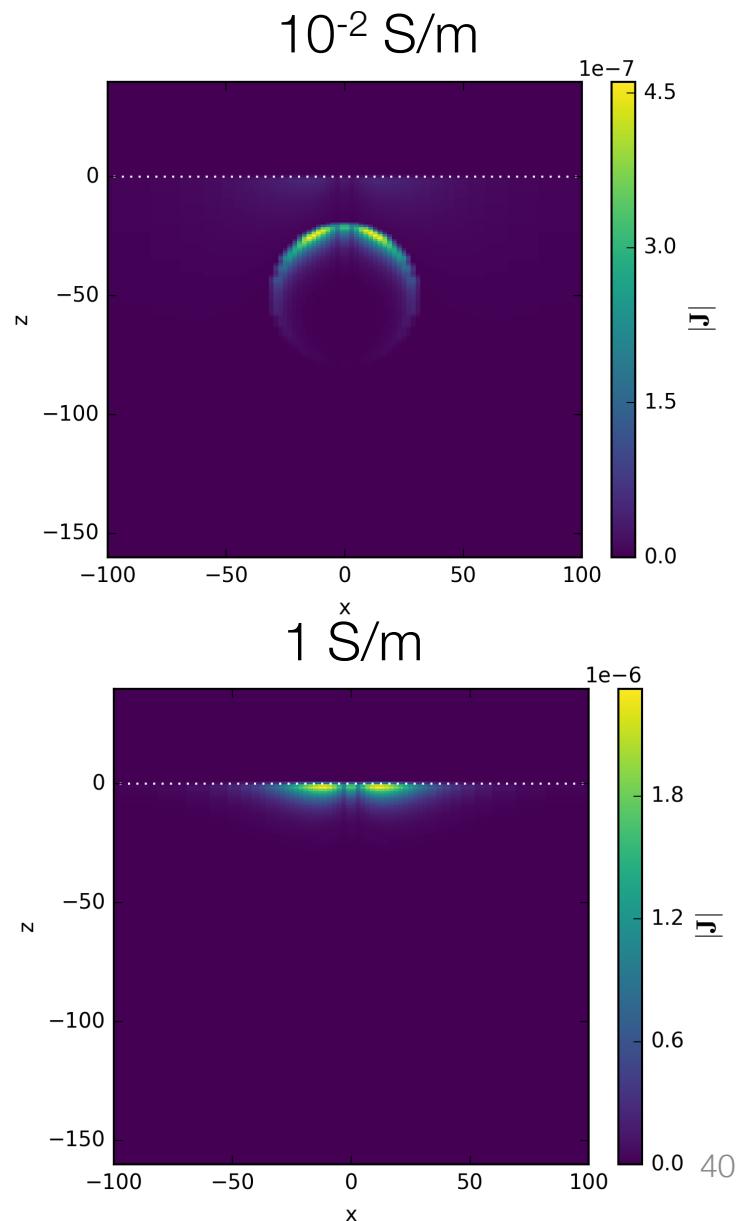
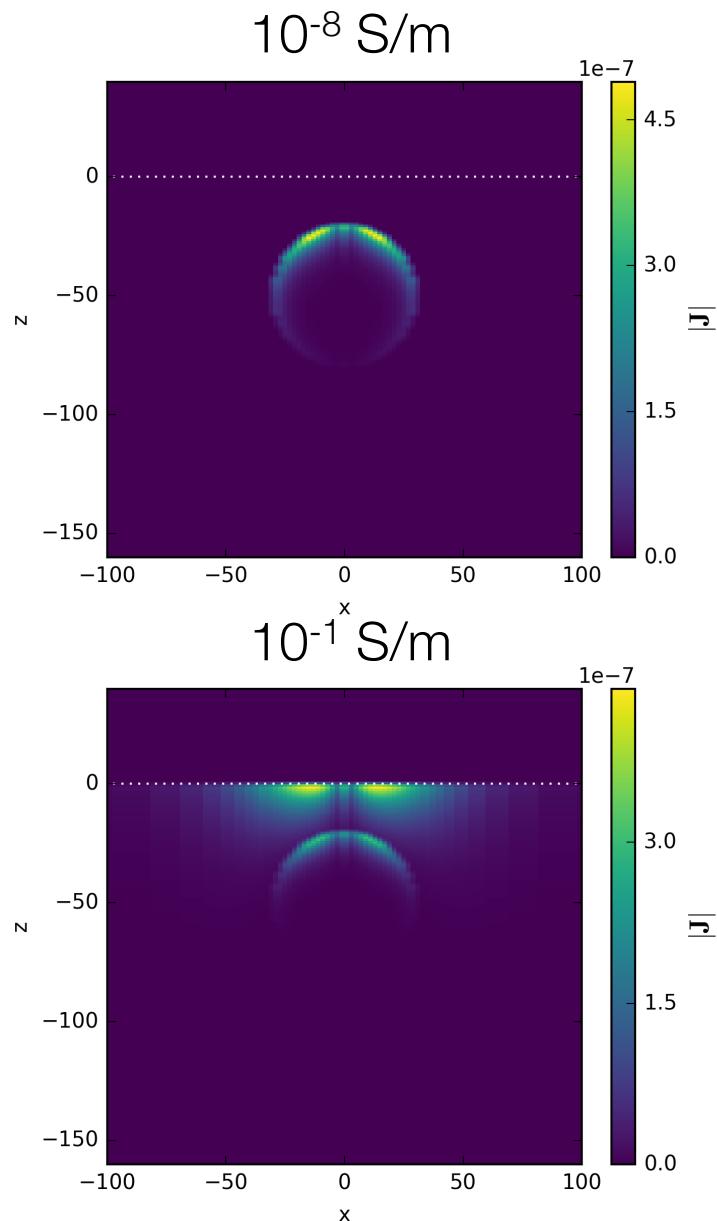


Current Density



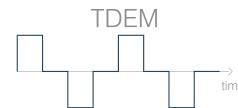
10^4 Hz

Effects of background resistivity: Frequency

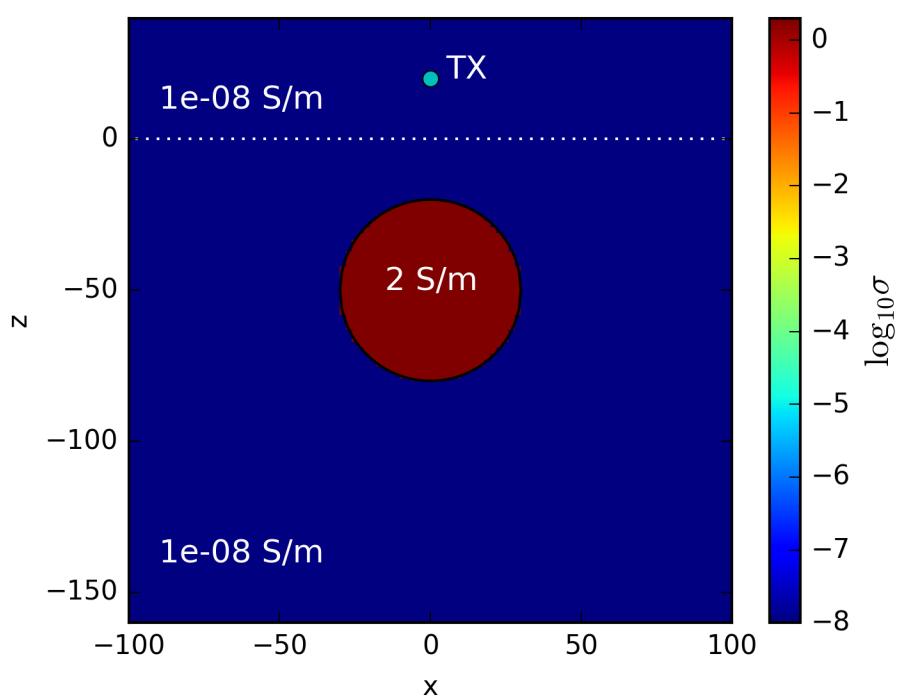


Effects of background resistivity: Time

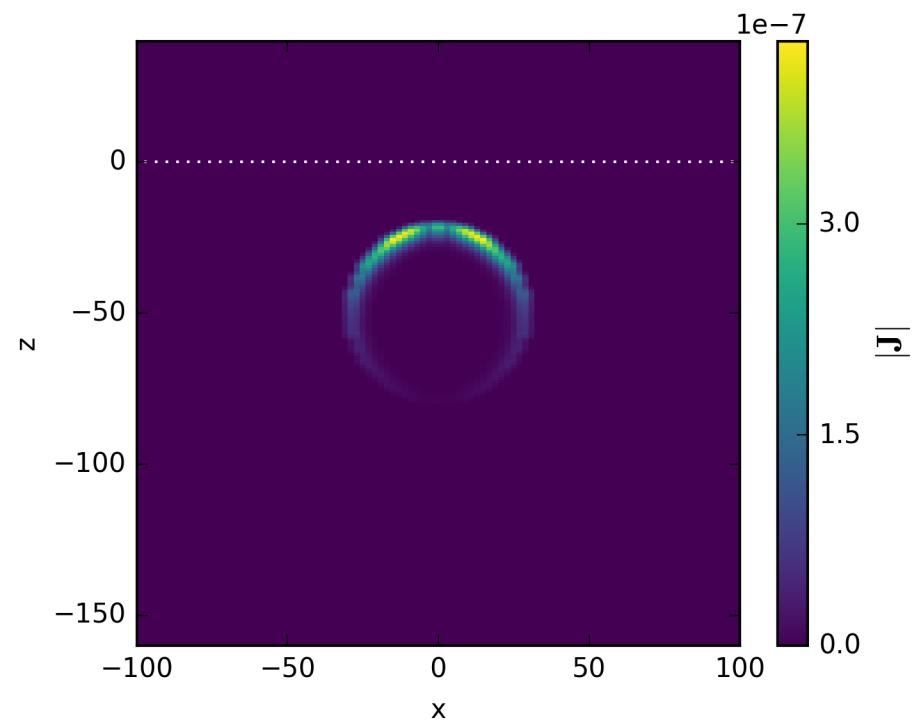
- Buried, conductive sphere
- Vary background conductivity
- Time: 10^{-5} s



10^{-8} S/m background

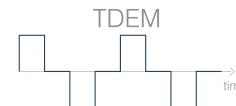


Current Density

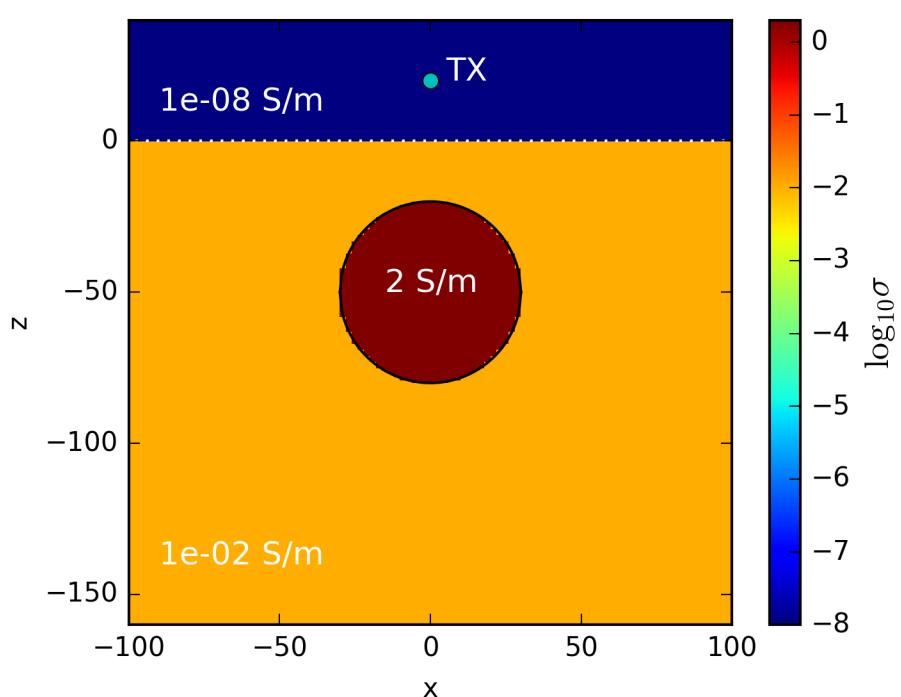


Effects of background resistivity: Time

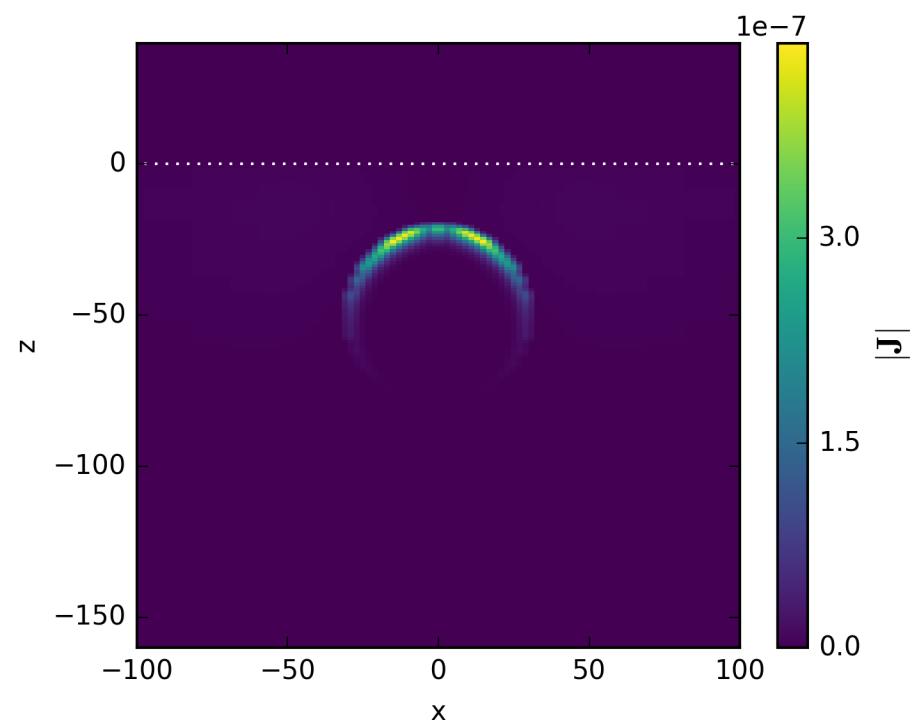
- Buried, conductive sphere
- Vary background conductivity
- Time: 10^{-5} s



10^{-2} S/m background



Current Density

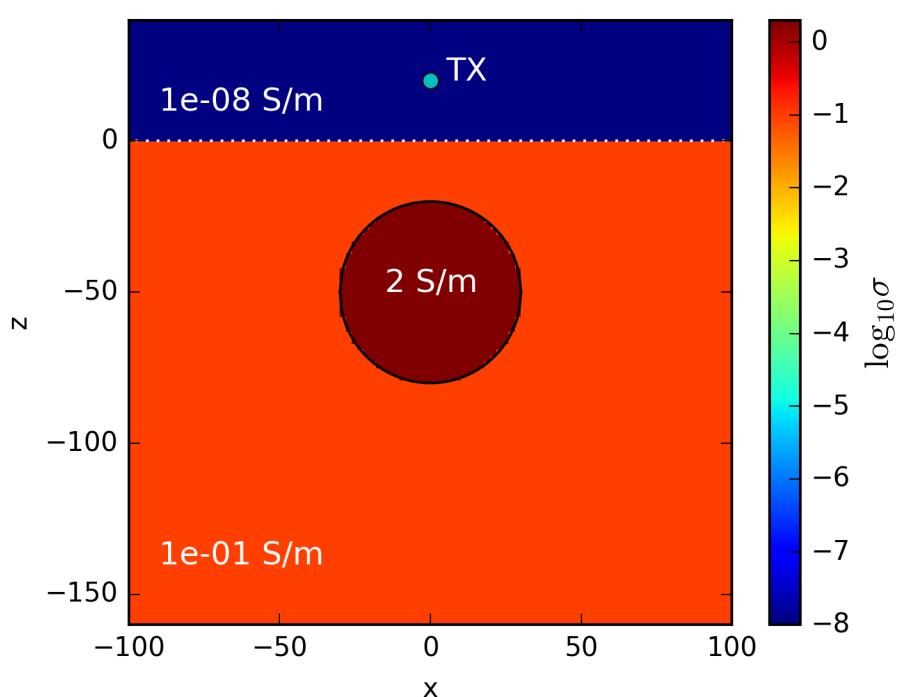


Effects of background resistivity: Time

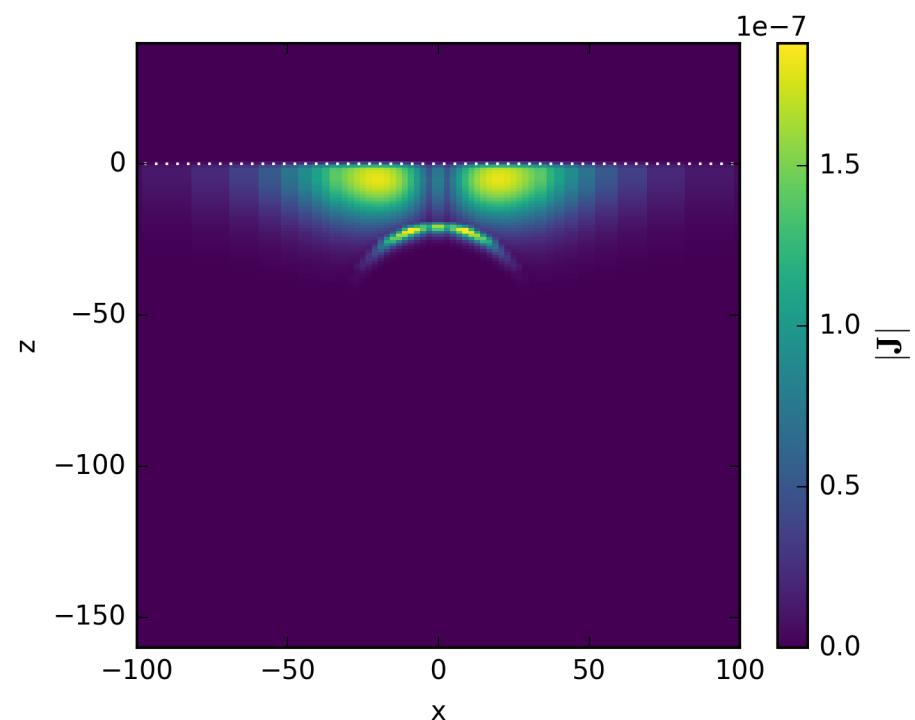
- Buried, conductive sphere
- Vary background conductivity
- Time: 10^{-5} s



10^{-1} S/m background

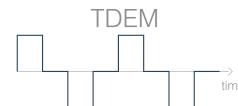


Current Density

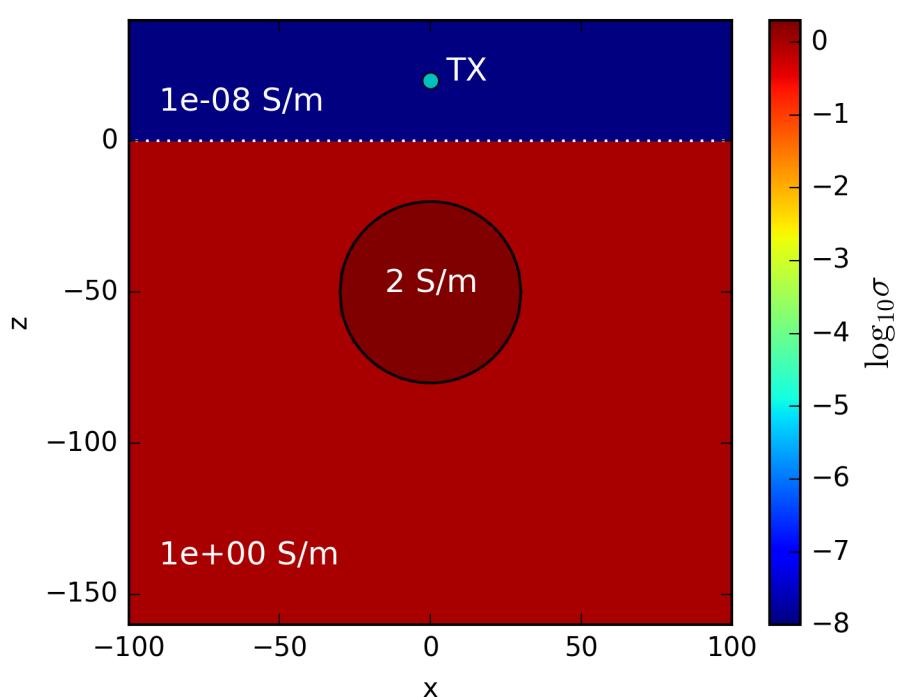


Effects of background resistivity: Time

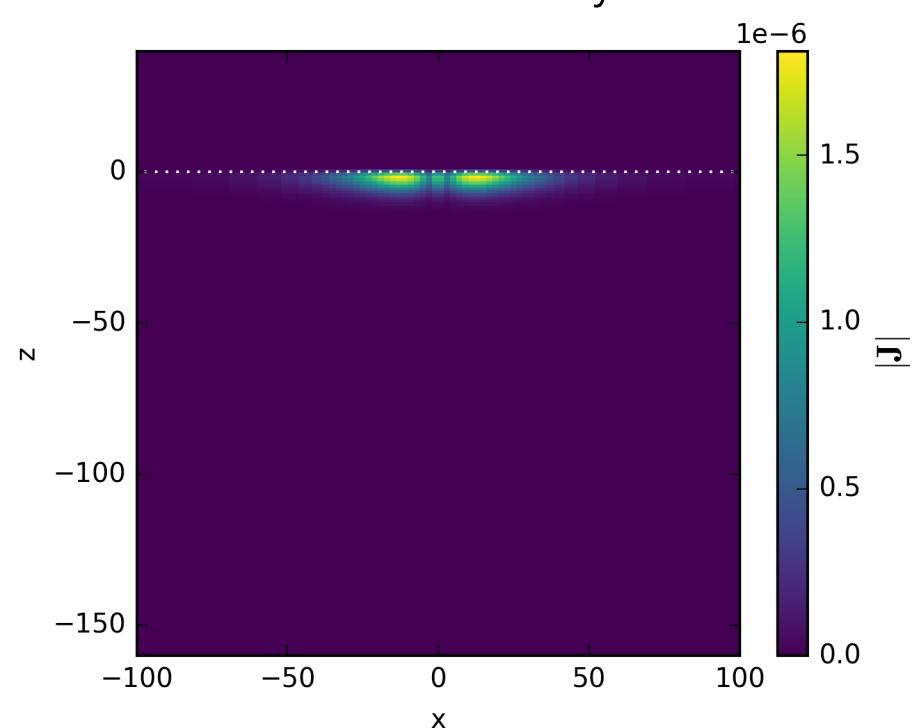
- Buried, conductive sphere
- Vary background conductivity
- Time: 10^{-5} s



1 S/m background

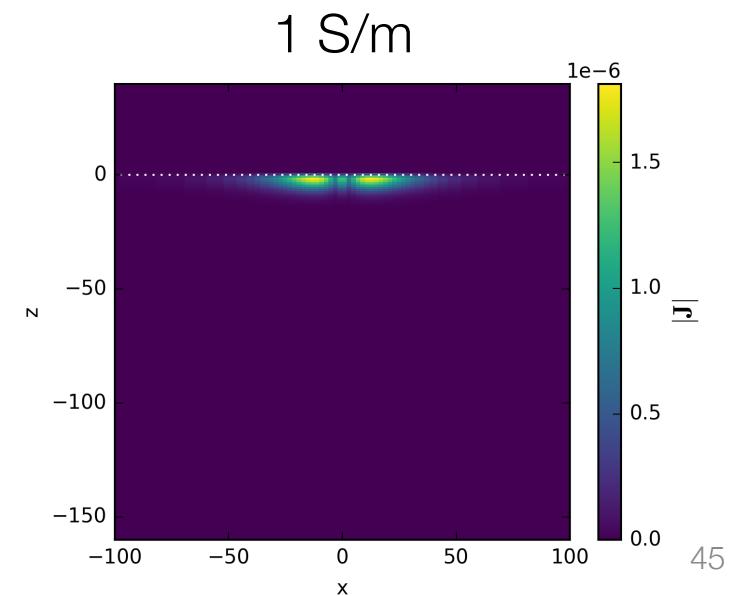
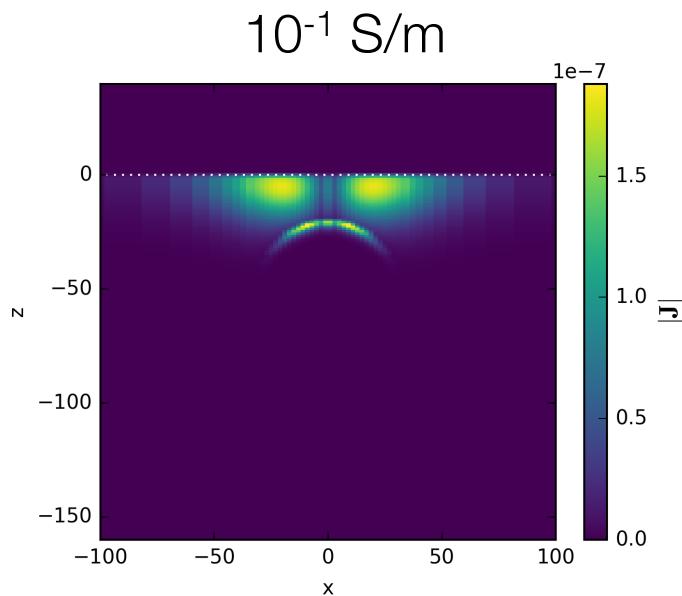
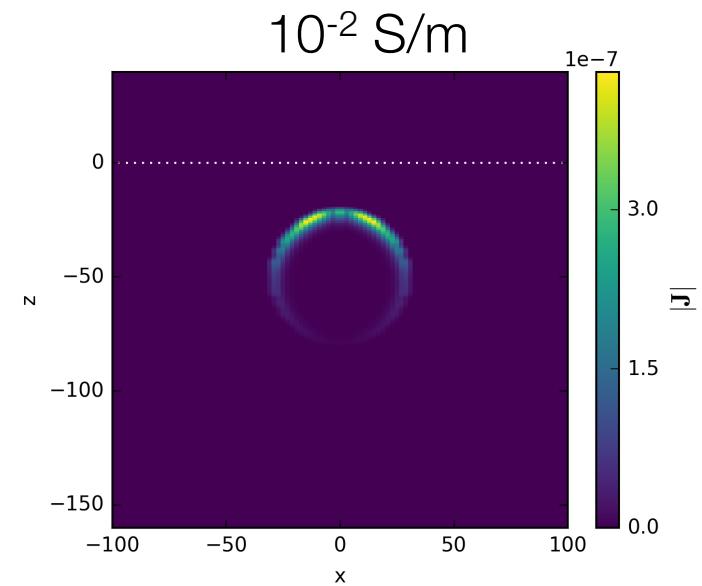
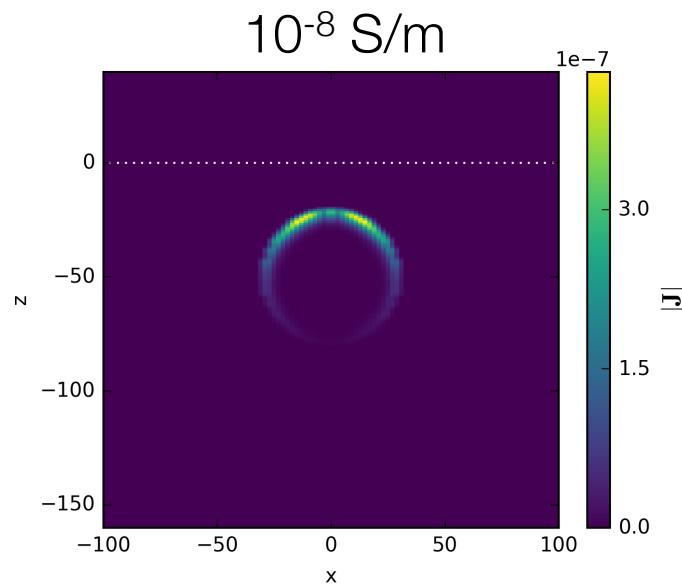


Current Density



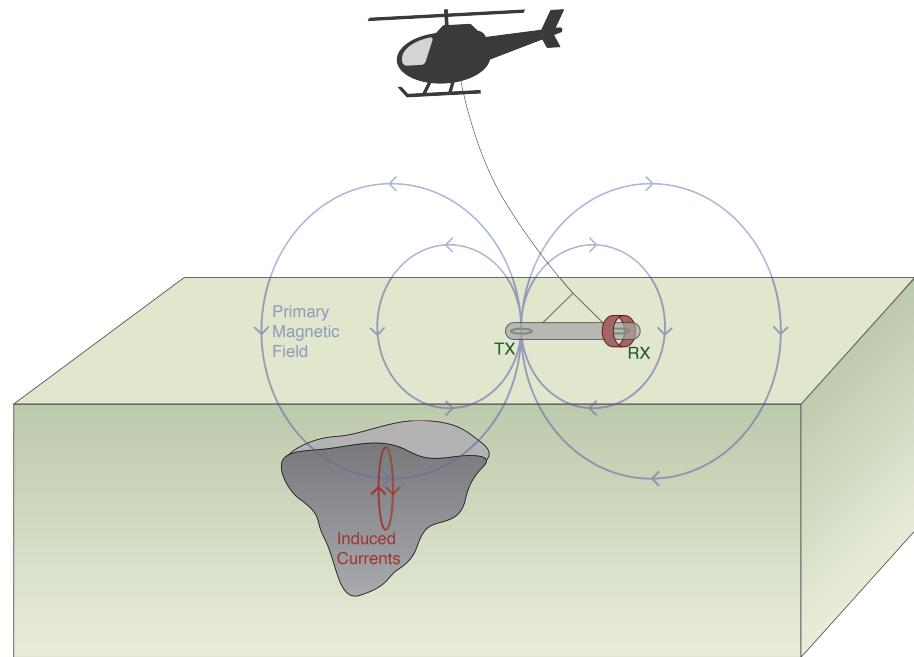
10^{-5} s

Effects of background resistivity: Time

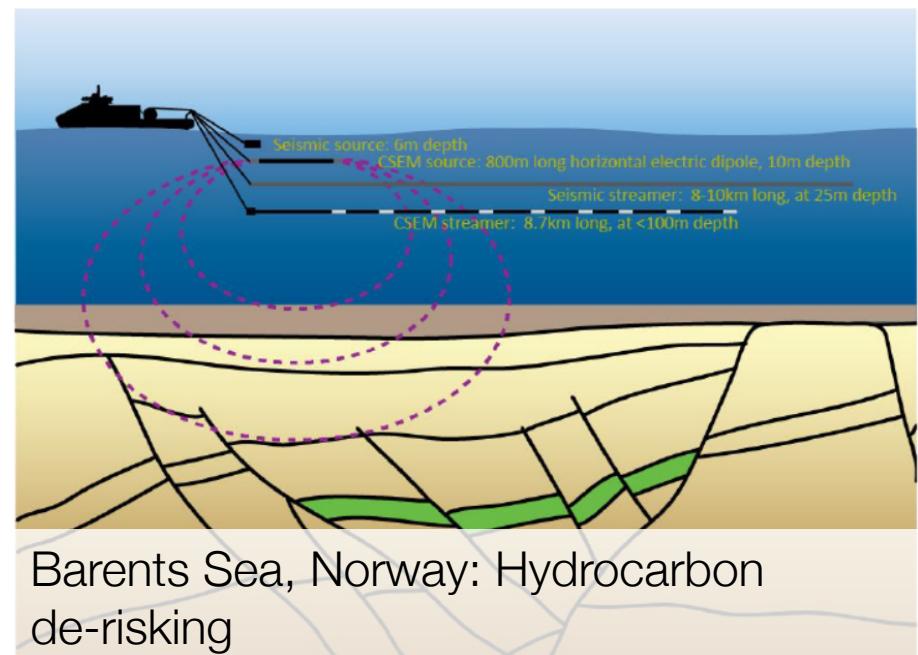
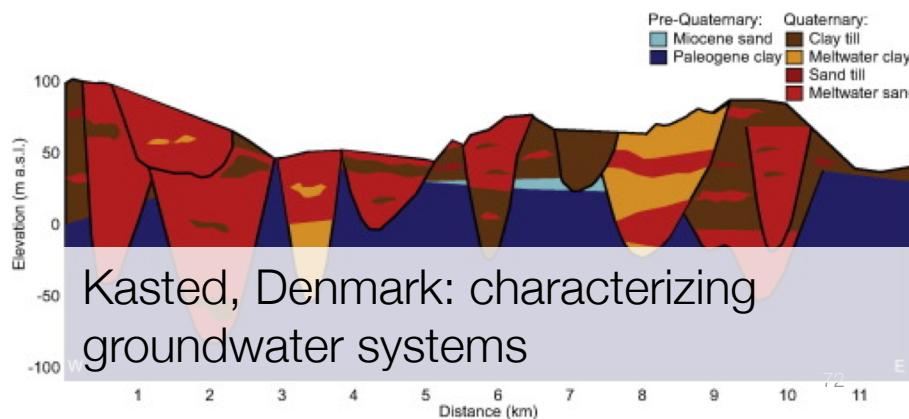
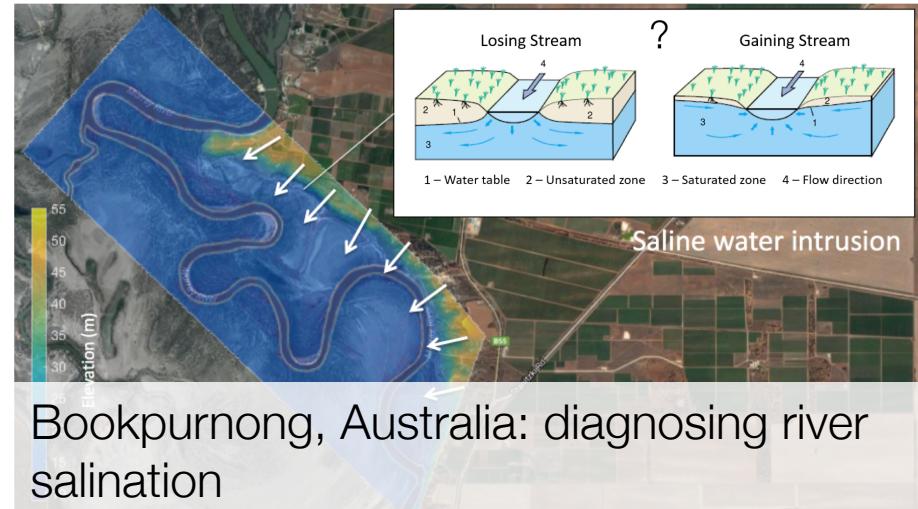
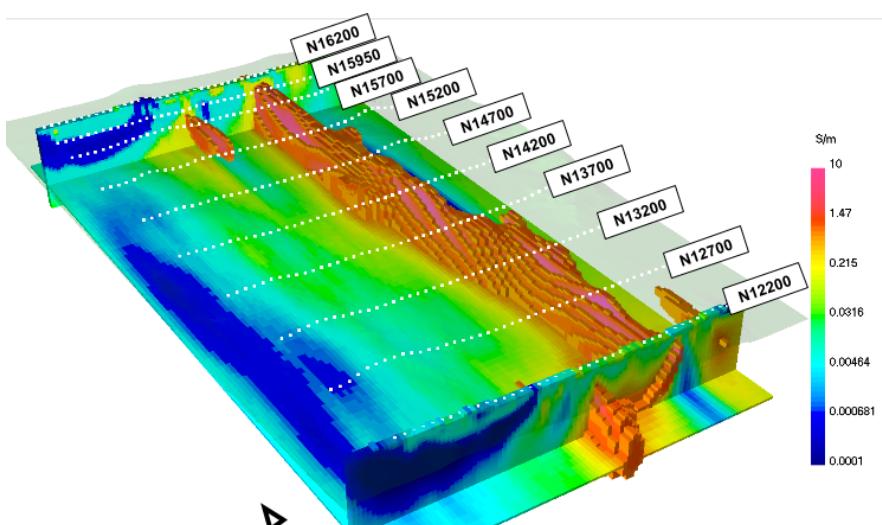


Recap: what have we learned?

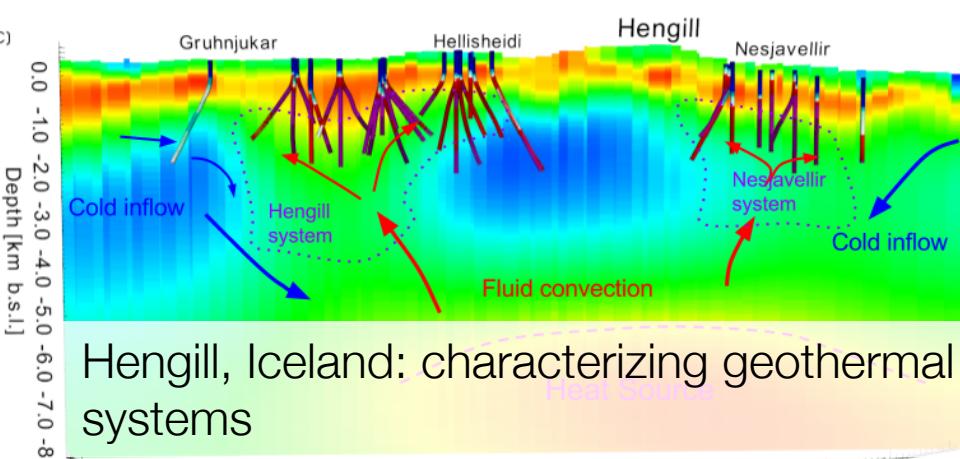
- Basics of EM induction
- Response functions
- Mutual coupling
- Data for frequency or time domain systems
- Circuit model is a good proxy
- Need to account for energy losses
- Ready to look at some field examples



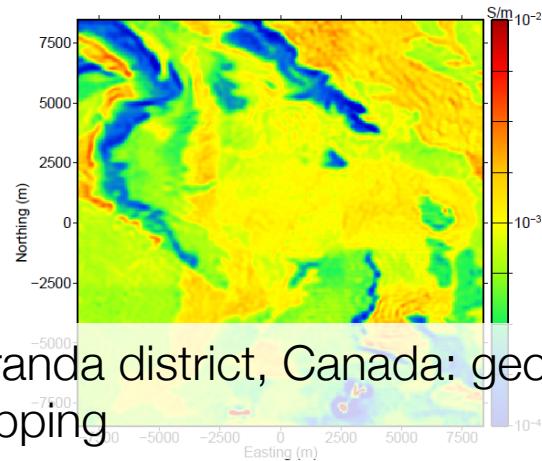
Today's Case Histories



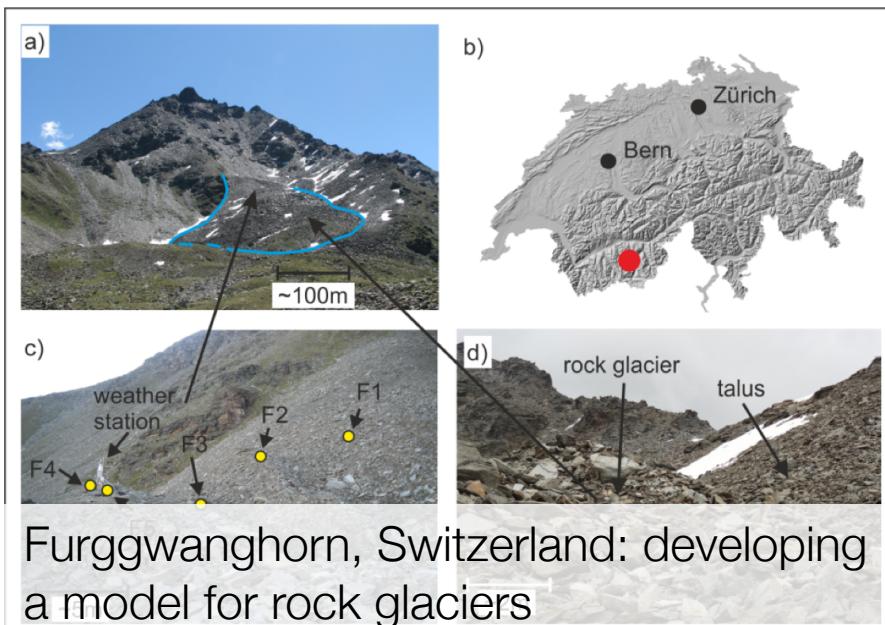
Today's Case Histories



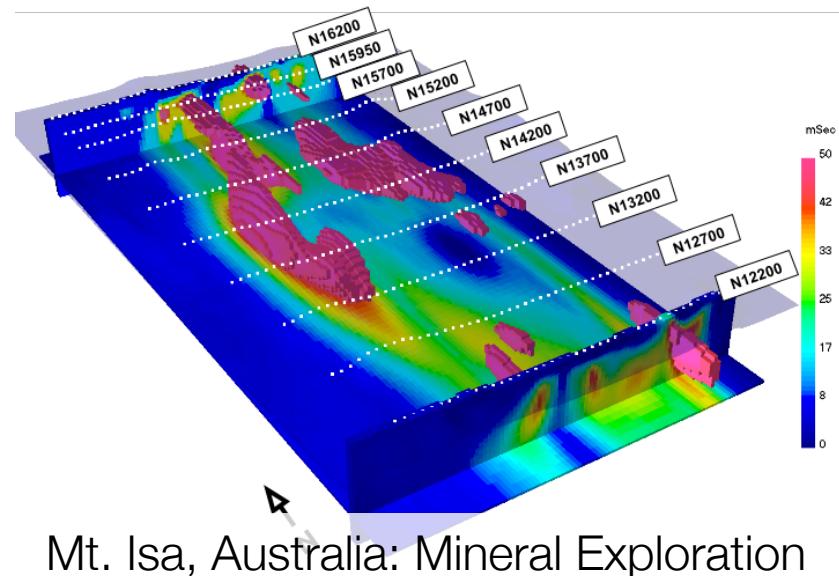
Hengill, Iceland: characterizing geothermal systems



Noranda district, Canada: geologic mapping



Furggwanghorn, Switzerland: developing a model for rock glaciers



Mt. Isa, Australia: Mineral Exploration

End of EM Fundamentals

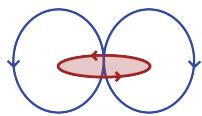
Next up



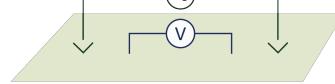
DC Resistivity



EM
Fundamentals



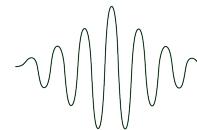
Inductive
Sources



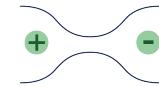
Grounded
Sources



Natural
Sources



GPR



Induced
Polarization



The
Future



Lunch: Play with apps