# OSCILLATIONS AND INSTABILITIES IN A PROPULSIVE MAGNETIC NOZZLE

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### **A**ZIMUTHAL OSCILLATIONS IN A MAGNETIC NOZZLE

- Helicon plasma thruster consists of
  - Cylindrical dielectric ionization chamber
  - EM inductor ("antenna")
  - Converging-diverging applied B field  $\rightarrow$  Magnetic nozzle





- Losses to lateral walls are large and suggest anomalous transport
- Gradient-driven drift instabilities are a candidate mechanism for that enhanced transport
- Previous works identify oscillations in the magnetic nozzle plasma plume



### EXPERIMENTAL SETUP

- HPT prototype running at 5 sccm Xe, 450 W EM power
- Max. B field (magnetic throat): 750 G
- Vacuum chamber: 1.5 m diameter, 3.5 m length,  $10^{-5}$  mbar during operation
- 3 floating, cylindrical tungsten probe tips, 1 cm apart (along  $\theta$  and d directions)
- Probe system displaced with radial-polar arm







 Averaging over 30 realizations to estimate the cross correlation spectrum between each 2 probes :

 $C_{12}(\omega) = X_1(\omega)X_2^*(\omega)$ 

• Mean and deviation of log power  $c = \log C$ :

$$\bar{\mu}_c = \frac{1}{n} \sum_{k}^{n} c_k \qquad \bar{\sigma}_c = \sqrt{\frac{1}{n-1} \sum_{k}^{n} (c_k - \bar{\mu}_c)^2}$$

 Mean and deviation phase difference (circular statistics):

$$ilde{z} = rac{1}{n} \sum_{k}^{n} \exp(i\phi_k) \quad ilde{\mu}_{\phi} = \arg ilde{z} \quad d_1 = rac{1}{n} \sum_{k}^{n} |\phi_k - \mu_{ ilde{\phi}}|$$

• Coherence (normalized cross correlation spectrum):

$$\hat{C}_{12}(\omega) = \frac{|C_{12}(\omega)|}{\sqrt{|X_1(\omega)|^2 |X_2(\omega)|^2}}$$





- At low  $\alpha$  angles ( $\alpha = 20$  deg), good coherence up to 100 kHz at 50 mm; not so good at 100 mm
- Dispersion relation in that range has  $m=0\div 1$ ,  $k_{\parallel}\simeq 0$ , corresponding to essentially global oscillations



![](_page_4_Picture_4.jpeg)

- At larger angles (α = 30 ÷ 35) and d = 100 ÷ 150 mm, a dominantly-azimuthal mode is found at f < 100 kHz with</li>
  - $m=1\div4$ ,  $k_{\parallel}<10$  rad/m
  - Peak in CSD power at ~60 kHz; harmonic at ~120 kHz

![](_page_5_Figure_4.jpeg)

![](_page_5_Picture_5.jpeg)

#### Similar features but flatter spectrum found at d = 150 mm

![](_page_6_Figure_2.jpeg)

![](_page_6_Picture_3.jpeg)

- Behavior at large  $\alpha$  is more complex.
  - Good coherence at low frequencies far downstream
  - $m = 1 \div 2$ ,  $k_{\parallel}$  has a dispersion relation

![](_page_7_Figure_4.jpeg)

![](_page_7_Picture_5.jpeg)

### SUMMARY OF FINDINGS - AND ISSUES

- 10 kHz Oscillations found everywhere, with  $m\simeq 0$  and  $k_{\parallel}\simeq 0$
- 40-60 kHz Oscillations at intermediate  $\alpha$  angles downstream, with m < 4 and  $k_{\parallel}$  likely nonzero
- 40-60 kHz Oscillations at large  $\alpha \simeq 50$  deg (plume periphery), with m < 2 and  $k_{\parallel} \neq 0$
- Phase velocities (< 10<sup>5</sup> m/s) comparable to estimated density-gradient and ExB drift velocities
- Coherence is not large ( $\sim 1$ ) in some of these
- Dispersion relation plots with bends could suggest multiple waves coexist
  - If two or more oscillations coexist at same ω,
     2-probe method is unable to resolve them
- Jumps exist at some, but not all locations, around  $\omega_{lh}$

![](_page_8_Figure_9.jpeg)

![](_page_8_Picture_10.jpeg)

### DISPERSION RELATION DISCUSSION

- Other works (below) also found azimuthal oscillations in similar setups
  - Hepner et al: suggest ECDI
  - Takahashi: suggests oscillations are a magnetosonic wave

![](_page_9_Figure_4.jpeg)

Hepner et al ECRT Appl. Phys. Lett. 116, 263502 (2020); Proposes outward electron transport

![](_page_9_Figure_6.jpeg)

Takahashi's HPT Scientific Reports (2022) 12:20137; Proposes inward electron transport

![](_page_9_Picture_8.jpeg)

- 3D wave dispersion relation for inhomogeneous plasma
  - Obtained from fluid approach
    - Locally Cartesian set of coordinates  $(r, \theta, z) \rightarrow (x, y, z)$

$$-i\omega_{s}n_{s} + u_{sx}\frac{\partial n_{0}}{\partial x} + u_{sx0}\frac{\partial n_{s}}{\partial x} + n_{0}\nabla\cdot\boldsymbol{u}_{s} + n_{s}\frac{\partial u_{sx0}}{\partial x} = \nu_{p}n_{s}$$

$$-i\omega_{s}\boldsymbol{u}_{s} + u_{sx}\frac{\partial \boldsymbol{u}_{s0}}{\partial x} + u_{sx0}\frac{\partial \boldsymbol{u}_{s}}{\partial x} = -\left(\left(\frac{\nabla\cdot\Pi}{nm}\right)_{s}^{(1)}\right) - \frac{q}{m}\nabla\phi_{1} + \frac{q}{|q|}\omega_{cs}\boldsymbol{u}_{s} \times \boldsymbol{b} - N_{s}\boldsymbol{u}_{s}$$

• **B** || **z** 

- Inertial, collisional and FLR effects considered for the electron flow
  - Stress tensor (II) comprised of both gyrotropic and gyroviscous parts

$$\Pi_{e} = p_{\perp} \boldsymbol{I} + \left( p_{\parallel} - p_{\perp} \right) \boldsymbol{b} \boldsymbol{b} + \hat{\Pi}_{e}$$

![](_page_10_Picture_9.jpeg)

- 3D wave dispersion relation for inhomogeneous plasma
  - Obtained from fluid approach

$$\frac{k^2 c_s^2}{\omega_{Pi}^2} \left( \frac{\omega_{Pi}^2}{\omega_i^2} - 1 \right) = \underbrace{\begin{array}{c} \omega_D - \omega_B + D_\perp + D_\parallel} \\ \omega_e + \omega_D - \omega_B + D_\perp + D_\parallel \end{array}$$

$$\omega_e \equiv \omega - \omega_E - \omega_D$$

 $\partial \ln B_0$ 

• Drift frequencies:

$$\omega_D \equiv -k_\perp \frac{c_e^2}{\omega_{ce}} \frac{\partial \ln n_0}{\partial x} \qquad \qquad \omega_E \equiv -k_\perp \frac{E_{x0}}{B_0} \qquad \qquad \omega_B \equiv -k_\perp \frac{c_e^2}{\omega_c}$$

- Inertial, collisional and FLR effects are accounted for
  - For perpendicular propagation:

$$D_{\perp} = D_{\perp} \left( \rho_e k_{\perp}, \nu_{\perp}, \omega, \omega_E, \omega_D, \omega_B \right)$$

• For parallel propagation:

$$D_{\parallel} = D_{\parallel} \left( \rho_e^2 k_{\parallel}^2, \rho_e^2 k_{\perp}^2, \nu_{\parallel}, \omega, \omega_E, \omega_D, \omega_B \right)$$

![](_page_11_Picture_12.jpeg)

- Delimited parametric regions for instabilities to occur
- For perpendicular propagation:
  - Blue region for collisionless instability
  - Dashed lines delimit region for collisional instability

![](_page_12_Figure_5.jpeg)

![](_page_12_Picture_6.jpeg)

- Delimited parametric regions for instabilities to occur
- For perpendicular propagation:
  - Blue region for collisionless instability
  - Dashed lines delimit region for collisional instability

![](_page_13_Figure_5.jpeg)

![](_page_13_Picture_6.jpeg)

- Delimited parametric regions for instabilities to occur
- For perpendicular propagation:
  - Blue region for collisionless instability
  - Dashed lines delimit region for collisional instability
- For both propagation directions:
  - Two different regions for collisonless instability
  - Collisions widen instability regions

![](_page_14_Figure_8.jpeg)

![](_page_14_Picture_9.jpeg)

- Delimited parametric regions for instabilities to occur
- For perpendicular propagation:
  - Blue region for collisionless instability
  - Dashed lines delimit region for collisional instability
- For both propagation directions:
  - Two different regions for collisonless instability
  - Collisions widen instability regions

![](_page_15_Figure_8.jpeg)

![](_page_15_Picture_9.jpeg)

### SUMMARY AND WAY FORWARD

- At low frequency ranges (< 100 kHz), floating probe measurements show existence of mainly azimuthal oscillations with low m number (and likely nonzero k<sub>||</sub>). Oscillations are more apparent downstream, at moderate and high angles α from axis. Likely, various types of oscillations are present at the same time.
- Present measurements are partially inconclusive: low coherence, "dirty" dispersion relation plots
- Adding a 3<sup>rd</sup> probe could help discriminate the coexistence of multiple waves (which would be affecting current results)
- Phase velocity is compatible with drift velocities
- Local linear dispersion analysis in slab geometry including collisionality,  $k_{\parallel}$ , and gyroviscous terms shows:
  - Collisionless regions of instability are affected by  $k_{\parallel}$
  - Collisions add a weak instability almost everywhere in the parametric plane
  - Our operating point corresponds with one such weak instability condition, which could justify experimental observations
- Effect of oscillations/instabilities on  $\perp$  transport still unclear

![](_page_16_Picture_10.jpeg)

# ACKNOWLEDGMENTS

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![](_page_17_Picture_2.jpeg)

# THANK YOU!

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![](_page_18_Picture_2.jpeg)

![](_page_18_Picture_3.jpeg)

### STEADY-STATE RESULTS

B[G]160 140 120 100  $r \; [mm]$ 80 60 40  $10^{1}$ 20 0 0 50 100 150 200  $z \,[\mathrm{mm}]$  $n \, [m^{-3}]$  $\phi$  [V]  $10^{17}$ 100 80 90 60  $r \, [\mathrm{mm}]$  $10^{16}$ 40 70 20 100 120 140 40 60 100 40 60 80 80 120 140  $z \, [\mathrm{mm}]$  $z \, [\mathrm{mm}]$ 

#### • $B \simeq 10 \div 100 \text{ G}$

- $\omega_{ce} \simeq 2.8 \cdot 10^7 \div 2.8 \cdot 10^8 \text{ Hz}$
- $\omega_{ci} \simeq 1.2 \cdot 10^2 \div 1.2 \cdot 10^3$  Hz (Xenon)
- $\omega_{lh} \simeq 5.7 \cdot 10^4 \div 5.7 \cdot 10^5 \text{ Hz}$
- $n \simeq 10^{15} \div 10^{17} \text{ m}^{-3}$ 
  - $\omega_{pe} = 9.0 \cdot 10^8 \text{ Hz}$
- $T_e \simeq 12.5 \text{ eV}$ 
  - $c_s = 3 \cdot 10^3 \, \text{m/s}$
- Estimated gradient lenghts
  - $\ln n \rightarrow 1 \text{ cm}$
  - $e\phi/T_e \rightarrow 1 \text{ cm}$
  - $\ln B \rightarrow 10 \text{ cm}$
- Estimated drift velocities
  - $u_D, u_E \simeq 10^4 \div 10^5 \text{ m/s}$

 Collisionless instabilities arising from interaction between branches

#### • No parallel propagation:

![](_page_20_Figure_3.jpeg)

![](_page_20_Picture_4.jpeg)

- Collisionless instabilities arising from interaction between branches
  - Branch 1: non-trivial solution of parallel component of momentum ( $\omega_{\parallel} = 0$ )
    - Does not interact with other branches

• No parallel propagation:

![](_page_21_Figure_5.jpeg)

![](_page_21_Picture_6.jpeg)

- Collisionless instabilities arising from interaction between branches
  - Branch 1: non-trivial solution of parallel component of momentum ( $\omega_{\parallel}=0$ )
    - Does not interact with other branches
  - Branches 2 and 3: destabilization of «anti-drift wave» (modified SHI, with inertial and FLR effects)

No parallel propagation: Imaginary part of branch 2  $\rho_e/L_n = 0.008, \rho_e/L_E = 0.02, \sqrt{m_e/m_i} = 0.0001, \rho_e \kappa_{\parallel} = 0.0, \nu_{\perp} = 0.0, \nu_{\parallel} = 0.0$ 1.00 0.6 0.75 0.4 0.50 0.2 0.25 W<sub>L</sub>H 0.0 × HI 0.00 Real parts of branches 2 and 3 -0.25-0.50-0.4-0.75-0.6 -1.000.1 0.2 0.4 0.5 0.0 0 Imaginary part of branch 3

![](_page_22_Picture_6.jpeg)

- Collisionless instabilities arising from interaction between branches
- Including parallel propagation introduces new branches (and new ways for destabilizations to take place)

![](_page_23_Figure_4.jpeg)

![](_page_23_Picture_5.jpeg)

- Collisionless instabilities arising from interaction between branches
- Including parallel propagation introduces new branches (and new ways for destabilizations to take place)
  - Branch 1 now «sees» other branches
    - Its solution is no longer decoupled from the rest of the problem

![](_page_24_Figure_6.jpeg)

![](_page_24_Picture_7.jpeg)

- Collisionless instabilities arising from interaction between branches
- Including parallel propagation introduces new branches (and new ways for destabilizations to take place)
  - Branch 1 now «sees» other branches
  - Branches 2 and 3 might still interact and destabilize under simil-SHI conditions

![](_page_25_Figure_6.jpeg)

![](_page_25_Picture_7.jpeg)

- Collisionless instabilities arising from interaction between branches
- Including parallel propagation introduces new branches (and new ways for destabilizations to take place)
  - Branch 1 now «sees» other branches
  - Branches 2 and 3 might still interact and destabilize under simil-SHI conditions
    - Stable close to the origin

![](_page_26_Figure_7.jpeg)

![](_page_26_Picture_8.jpeg)

- Collisionless instabilities arising from interaction between branches
- Including parallel propagation introduces new branches (and new ways for destabilizations to take place)

![](_page_27_Figure_4.jpeg)

![](_page_27_Picture_5.jpeg)

- Collisionless instabilities arising from interaction between branches
- Including parallel propagation introduces new branches (and new ways for destabilizations to take place)
  - Depending on the choice of parameters, branch 1 might interact with one of the two sonic branches

![](_page_28_Figure_5.jpeg)

![](_page_28_Picture_6.jpeg)

- Collisionless instabilities arising from interaction between branches
- Including parallel propagation introduces new branches (and new ways for destabilizations to take place)
  - Depending on the choice of parameters, branch 1 might interact with one of the two sonic branches
    - New conditions for instability

![](_page_29_Figure_6.jpeg)

![](_page_29_Picture_7.jpeg)

- Collisionless instabilities arising from interaction between branches
- Including parallel propagation introduces new branches (and new ways for destabilizations to take place)
  - Depending on the choice of parameters, branch 1 might interact with one of the two sonic branches
    - New conditions for instability

![](_page_30_Figure_6.jpeg)

![](_page_30_Picture_7.jpeg)

• Collisions widen parametric range favourable for instabilities to arise

![](_page_31_Picture_2.jpeg)

- Collisions widen parametric range favourable for instabilities to arise
- Consider a point where no collisionless instability takes place:

![](_page_32_Figure_3.jpeg)

![](_page_32_Picture_4.jpeg)

- Collisions widen parametric range favourable for instabilities to arise
- Consider a point where no collisionless instability takes place:

![](_page_33_Figure_3.jpeg)

#### • Without collisions:

![](_page_33_Figure_5.jpeg)

![](_page_33_Picture_6.jpeg)

- Collisions widen parametric range favourable for instabilities to arise
- Consider a point where no collisionless instability takes place:

![](_page_34_Figure_3.jpeg)

#### • With collisions:

![](_page_34_Figure_5.jpeg)

![](_page_34_Picture_6.jpeg)

- Collisions widen parametric range favourable for instabilities to arise
- Consider a point where no collisionless instability takes place:

![](_page_35_Figure_3.jpeg)

#### • With collisions:

![](_page_35_Figure_5.jpeg)

![](_page_35_Picture_6.jpeg)